

Exact solution for a Stefan problem with convective boundary condition and density jump

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We consider the solidification of a semi-infinite material which is initially at its liquid phase at a uniform temperature T_i . Suddenly at time $t > 0$ the fixed face $x = 0$ is submitted to a convective cooling condition with a time-dependent heat transfer coefficient of the type $H(t) = ht^{-1/2}$ ($h > 0$). The bulk temperature of the liquid at a large distance from the solid-liquid interface is T_∞ , a constant temperature such that $T_\infty < T_f < T_i$ where T_f is the freezing temperature. We also consider the density jump between the two phases. We obtain that the corresponding phase-change (solidification) process has an explicit solution of a similarity type for the temperature of both phases and the solid-liquid interface, if and only if the coefficient h is large enough, that is $h > h_0 = \frac{k_l}{\sqrt{\pi\alpha_l}} \frac{T_i - T_f}{T_i - T_\infty}$ where k_l and α_l are the conductivity and diffusion coefficients of the initial liquid phase. Moreover, when $h \leq h_0$ we only have a heat conduction problem for the initial liquid phase and the corresponding change of phase does not occur. We do the mathematical analysis of a problem which was presented in S.M. Zubair - M.A. Chaudhry, Wärme-und-Stoffübertragung, 30 (1994), 77-81 and we generalized the results obtained in D.A. Tarzia, MAT-Serie A, 8 (2004), 21-27 when the density jump between the two phases was neglected.

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1 Introduction

Heat transfer problems involving a change of phase due to melting or freezing processes are very important in science and technology [3], [5]. This kind of problems are generally referred as moving-free boundary problems which have been the subject of numerous theoretical, numerical and experimental investigations, e.g. we can see the large bibliography on the subject given in [7]. We consider the solidification of a semi-infinite material which is initially at its liquid phase at a uniform temperature T_i . Suddenly at time $t > 0$ the fixed face $x = 0$ is submitted to a convective cooling condition due to a sudden drop in the ambient temperature. The bulk temperature of the liquid at a large distance from the solid-liquid interface is T_∞ , a constant such that $T_\infty < T_f < T_i$ where T_f is the freezing temperature. The density jump between the two phases are also considered. In order to solve the phase-change process with a convective condition at the fixed face $x = 0$, approximate methods were used, for example in [1], [2], [4]. In [9] the problem was analyzed and a closed-form expression for the solid-liquid interface and both temperatures were found when the heat transfer coefficient h is time-dependent and proportional to $t^{-1/2}$. The solution is obtained graphically. The goal of this paper is to give the mathematical analysis of this problem, that is the solidification of a semi-infinite material which is initially at the constant temperature T_i and a convective cooling condition is imposed at the fixed boundary $x = 0$ for a time-dependent heat transfer coefficient of the type

$$H(t) = h/\sqrt{t}, \quad h > 0, \quad t > 0, \quad (1)$$

and the density jump between the two solid and liquid phases are also considered. We prove that there exists an instantaneous phase-change process if and only if the coefficient h is large enough, that is

$$h > \frac{k_l}{\sqrt{\pi\alpha_l}} \frac{T_i - T_f}{T_i - T_\infty} \quad (2)$$

where k_i and $\alpha_i = k_i/\rho_i c_i$ are the conductivity and diffusion coefficients of the phase i ; c_i and ρ_i are the specific heat and the mass density of the phase i ($i = s$: solid; $i = l$: liquid). Moreover, the explicit expression for the solid-liquid interface $s(t)$ and the temperatures of the solid $T_s(x, t)$ and liquid $T_l(x, t)$ phases respectively can be obtained if and only if the inequality (2) is verified for the coefficient h which characterizes the dependent-time heat transfer coefficient $H(t)$ given by (1).

2 Explicit solution for a convective boundary condition and density jump.

We consider the following free boundary problem: find the solid-liquid interface $x = s(t)$ and the temperature $T(x, t)$ defined by $T_s(x, t)$ if $0 < x < s(t)$, $t > 0$, T_f if $x = s(t)$, $t > 0$ and $T_l(x, t)$ if $x > s(t)$, $t > 0$ which satisfy the following equations and boundary conditions

$$T_{s,t} = \alpha_s T_{s,xx}, \quad 0 < x < s(t), \quad t > 0 \quad (3)$$

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$$T_{l,t} \left(\frac{\rho_s - \rho_l}{\rho_l} \dot{s}(t) T_{l,x} = \alpha_l T_{l,xx}, x > s(t), t > 0 \right) \quad (4)$$

$$T_s(s(t), t) = T_l(s(t), t) = T_f, x = s(t), t > 0 \quad (5)$$

$$T_l(x, 0) = T_l(+\infty, t) = T_i, x > 0, t > 0 \quad (6)$$

$$k_s T_{s,x}(0, t) = \frac{h}{\sqrt{t}} (T_s(0, t) - T_\infty), t > 0 \quad (7)$$

$$k_s T_{s,x}(s(t), t) - k_l T_{l,x}(s(t), t) = \rho_s l \dot{s}(t), t > 0 \quad (8)$$

$$s(0) = 0 \quad (9)$$

where l is the latent heat of fusion, and $T_\infty < T_f < T_i$.

We obtain the following results:

Theorem 2.1 (i) If the coefficient h verifies the inequality (2) then the free boundary problem (3)–(9) has the explicit solution of a similarity type given by

$$s(t) = 2\xi\sqrt{\alpha_s t} \quad (10)$$

$$T_s(x, t) = T_\infty + \frac{(T_f - T_\infty) \left[1 + \frac{h\sqrt{\pi\alpha_s}}{k_s} \operatorname{erf}\left(\frac{x}{2\sqrt{\alpha_s t}}\right) \right]}{1 + \frac{h\sqrt{\pi\alpha_s}}{k_s} \operatorname{erf}(\xi)}, 0 < x < s(t), t > 0 \quad (11)$$

$$T_l(x, t) = T_f + (T_i - T_f) \left[\operatorname{erf}\left(\delta + \frac{x}{2\sqrt{\alpha_l t}}\right) - \operatorname{erf}\left((1+\varepsilon)\sqrt{\frac{\alpha_s}{\alpha_l}}\xi\right) \right], x > s(t), t > 0 \quad (12)$$

where erf is the error function and erfc is the complementary error function defined by $\operatorname{erfc}(z) = 1 - \operatorname{erf}(z)$, and the dimensionless parameter $\xi > 0$ satisfies the following equation

$$G(x) = x, x > 0 \quad (13)$$

where function G and the coefficients δ and ε are given by

$$G(x) = b_1 \frac{\exp(-x^2)}{1 + b_2 \operatorname{erf}(x)} - b_3 \frac{\exp(-b^2 x^2)}{\operatorname{erfc}(bx)} \quad (14)$$

$$\delta = \varepsilon \xi \sqrt{\frac{\alpha_s}{\alpha_l}}, \quad \varepsilon = \frac{|\rho_s - \rho_l|}{\rho_l}, \quad \alpha_0 = \frac{\alpha_l}{(1+\varepsilon)^2}, \quad b = \sqrt{\frac{\alpha_s}{\alpha_0}}, \quad (15)$$

$$b_1 = \frac{h(T_f - T_\infty)}{\rho_s l \sqrt{\alpha_l}}, \quad b_2 = \frac{h}{k_s} \sqrt{\pi \alpha_s}, \quad b_3 = \frac{k_l(T_i - T_f)}{\rho_s l \sqrt{\alpha_s \alpha_l \pi}}. \quad (16)$$

(ii) The Eq. (13) has a unique solution ξ if and only if the coefficient h satisfies the inequality (2). In this case, there exists an instantaneous solidification process.

Remark 2.2 (i) We note that the temperature at the fixed face $x = 0$ is given by $T_s(0, t) = \text{Const.} < T_f$.

(ii) We note that the inequality (2) for the coefficient h is of the type that it was obtained in [6] when a time-dependent heat flux condition on the fixed face is imposed.

(iii) When the jump density between the two phases is neglected, i.e. $\varepsilon = 0$, then we find the result obtained in [8].

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