

NECESSARY AND SUFFICIENT CONDITION TO OBTAIN N PHASES IN A ONE-DIMENSIONAL MEDIUM WITH A FLUX CONDITION ON THE FIXED FACE

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ABSTRACT

We consider an n -phase Stefan Problem for a semi-infinite material $x \geq 0$ with constant initial condition and a heat flux of the type $q(t) = -q_0/\sqrt{t}$, imposed on the fixed face $x = 0$. We determine necessary and/or sufficient conditions on the parameter q_0 , in order to obtain the existence of the solution. We also show the equivalence between this problem and that in which a temperature on the fixed face is imposed.

Key words : Multi-phase Stefan problem, Similarity variable, Neumann solution, Exact solutions, Free boundary problems.

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I. INTRODUCTION

We consider a semi-infinite material represented by $(0, +\infty)$, with a constant initial temperature u_1 . We apply a heat flux of the type $q(t) = -q_0/\sqrt{t}$ with $q_0 > 0$ on the fixed face $x = 0$. If the temperature's values are between u_i and u_{i+1} ($u_i < u_{i+1}$), we will consider that the material is in the phase i . Each of these phases is separated by a free boundary.

In this work, based on [4,7], we determine which conditions must verify the parameter q_0 , for the n -phases to appear. We show that the necessary and sufficient condition is:

$$q_0 > \frac{\beta_{n-1}}{\operatorname{erf} \frac{G_{n-2}(x_{n-2})}{a_{n-1}}} > \frac{\beta_{n-2}}{\operatorname{erf} \frac{G_{n-3}(x_{n-3})}{a_{n-2}}} > \dots > \frac{\beta_2}{\operatorname{erf} \frac{x_1}{a_2}} > \beta_1$$

where β_i , G_i and x_i are defined by (13), (25) and (26) respectively.

II. THE PROBLEM

We wish to find conditions on q_0 for the n -phases to appear, in other words, conditions such that there exist the free boundaries $S_{n-1}(t) < S_{n-2}(t) < \dots < S_1(t)$ and the temperature $\Theta = \Theta(x, t)$, given by:

$$(1) \quad \Theta(x, t) = \Theta_i(x, t) \quad \text{if} \quad S_i(t) < x < S_{i-1}(t), \quad t > 0, \quad i = 1, \dots, n,$$

and such that the following conditions be satisfied:

$$(2) \quad \frac{\partial \Theta_i}{\partial t} = \alpha_i \frac{\partial \Theta_i}{\partial x^2}, \quad x \in (S_i(t), S_{i-1}(t)), \quad t > 0, \quad i = 1, \dots, n,$$

$$(3) \quad \Theta_1(x, 0) = u_1, \quad x > 0,$$

$$(4) \quad \Theta_1(S_0(t), t) = u_1, \quad t > 0,$$

$$(5) \quad \Theta_i(S_i^+(t), t) = \Theta_{i+1}(S_{i+1}^-(t), t) = u_{i+1}, \quad t > 0, \quad i = 1, \dots, n-1,$$

$$(6) \quad k_n \frac{\partial \Theta_n}{\partial x}(0, t) = -\frac{q_0}{\sqrt{t}}, \quad t > 0,$$

$$(7) \quad S_i(0) = 0, \quad i = 1, \dots, n-1,$$

$$(8) \quad k_i \frac{\partial \Theta_i}{\partial x}(S_i^+(t), t) - k_{i+1} \frac{\partial \Theta_{i+1}}{\partial x}(S_i^-(t), t) = \delta_i \dot{S}_i(t), \quad t > 0, \quad i = 1, \dots, n-1,$$

where k_i and $\alpha_i = \frac{k_i}{\rho c_i}$ are respectively the thermal conductivity and the diffusivity for the phase i , c_i is the specific heat of the phase i , ρ is the common mass density, $\delta_i > 0$ represents the latent heat used for passing from phase i to phase $i + 1$, the temperatures u_i verify $u_{i+1} > u_i$ for $i = 1, \dots, n-1$ and we consider $S_0(t) = +\infty$ and $S_n(t) = 0$, $\forall t > 0$. For more physical considerations see [6,7].

Remark 1: The problem (2)–(8) can be stated in a similar way for dependent variable concentration or enthalpy.

III. SOLUTION OF THE PROBLEM (2)–(8)

Following the Neumann's solution idea, for the two-phase Stefan problem [1,2,5] (for the multi-phase Stefan problem see [6,7]) we propose:

$$(9) \quad \begin{aligned} \Theta_i(x, t) &= A_i + B_i \operatorname{erf}\left(\frac{x}{2a_i\sqrt{t}}\right), & i &= 1, \dots, n, \\ S_i(t) &= 2\omega_i \sqrt{t}, & \omega_i &> 0, \quad i = 1, \dots, n-1, \end{aligned}$$

where for convenience we consider $a_i = \sqrt{\alpha_i}$ and

$$\operatorname{erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x \exp(-t^2) dt \quad \text{and} \quad \operatorname{erfc}(x) = 1 - \operatorname{erf}(x)$$

are the error and the complementary error functions respectively.

From conditions (3)–(6) we obtain:

$$(10) \quad A_i = \frac{u_{i+1} \operatorname{erf}\left(\frac{\omega_{i+1}}{a_i}\right) - u_i \operatorname{erf}\left(\frac{\omega_i}{a_i}\right)}{\operatorname{erf}\left(\frac{\omega_{i+1}}{a_i}\right) - \operatorname{erf}\left(\frac{\omega_i}{a_i}\right)}, \quad B_i = \frac{u_i - u_{i+1}}{\operatorname{erf}\left(\frac{\omega_{i+1}}{a_i}\right) - \operatorname{erf}\left(\frac{\omega_i}{a_i}\right)}, \quad i=1, \dots, n-1,$$

$$(11) \quad A_n = u_n + \operatorname{erf}\left(\frac{\omega_{n-1}}{a_n}\right) \frac{q_0 a_n \sqrt{\pi}}{k_n}, \quad B_n = -\frac{q_0 a_n \sqrt{\pi}}{k_n},$$

here we have inserted a dummy parameter $\omega_0 = +\infty$.

Conditions (8) ($i = 1, \dots, n-1$) are verified if and only if ω_i satisfy the following system:

$$(12) \quad \begin{aligned} \text{a)} \quad \delta_i \omega_i &= \beta_{i+1} \frac{\eta(\omega_i, a_{i+1})}{\phi(\omega_i, \omega_{i+1}, a_{i+1})} - \beta_i \frac{\eta(\omega_i, a_i)}{\phi(\omega_{i-1}, \omega_i, a_i)}, \quad i = 1, \dots, n-2, \\ \text{b)} \quad \delta_{n-1} \omega_{n-1} &= q_0 \eta(\omega_{n-1}, a_n) - \beta_{n-1} \frac{\eta(\omega_{n-1}, a_{n-1})}{\phi(\omega_{n-2}, \omega_{n-1}, a_{n-1})}, \end{aligned}$$

where:

$$(13) \quad \beta_i = -\frac{k_i}{a_i \sqrt{\pi}} (u_i - u_{i+1}) > 0, \quad i = 1, \dots, n-1,$$

$$(14) \quad \eta(\omega, a) = \exp\left(-\frac{\omega^2}{a^2}\right),$$

$$(15) \quad \phi(\alpha, \beta, \gamma) = \operatorname{erf}\left(\frac{\alpha}{\beta}\right) - \operatorname{erf}\left(\frac{\beta}{\gamma}\right).$$

In order to show the existence of a solution of the system (12) we will define the sequences of functions $\{h_i(x)\}_{i=1}^{n-1}$, $\{H_i(x)\}_{i=1}^{n-2}$ y $\{G_i(x)\}_{i=1}^{n-1}$.

We define:

$$(16) \quad G_1(x) = x, \quad x > 0,$$

$$(17) \quad h_1(x) = \delta_1 x + \beta_1 \frac{\eta(x, a_1)}{\operatorname{erfc}\left(\frac{x}{a_1}\right)}, \quad x > 0.$$

Taking into account that $F_1(x) = \frac{\exp(-x^2)}{\operatorname{erfc}(x)}$ ($x > 0$) is an increasing function, it results that h_1 verifies:

$$h_1(0) = \beta_1 > 0, \quad h_1(+\infty) = +\infty, \quad h_1'(x) > 0, \quad \forall x > 0,$$

then the equation (12a) for $i = 1$ is equivalent to:

$$(18) \quad h_1(\omega_1) = \beta_2 \frac{\eta(\omega_1, a_2)}{\phi(\omega_1, \omega_2, a_2)},$$

and hence $\phi(\omega_1, \omega_2, a_2) > 0$ and $\omega_1 > \omega_2$.

Now we define:

$$(19) \quad H_1(x) = \operatorname{erf} \frac{x}{a_2} - \beta_2 \frac{\eta(x, a_2)}{h_1(x)}, \quad x > 0.$$

From properties of the function h_1 it results that H_1 satisfies:

$$H_1(0) = -\frac{\beta_2}{\beta_1} < 0, \quad H_1(+\infty) = 1, \quad H_1'(x) > 0, \quad \forall x > 0.$$

This implies the existence of $x_1 > 0$ such that $H_1(x_1) = 0$ and then we can define the function

$$(20) \quad G_2(x) = a_2 \operatorname{erf}^{-1}[H_1(x)] \quad , \quad x \in (x_1, +\infty),$$

which is an increasing function and verifies $G_2(x_1) = 0$ and $G_2(+\infty) = +\infty$.

We can write equation (18) as:

$$(21) \quad \omega_2 = G_2(\omega_1) \quad ,$$

or equivalently by

$$(21\text{bis}) \quad \frac{\beta_2}{\phi(\omega_1, \omega_2, a_2)} = \frac{h_1(\omega_1)}{\eta(\omega_1, a_2)}.$$

Hence $G_2(x) < G_1(x)$, $\forall x \in (x_1, +\infty)$. Following [7] and proceeding inductively we can define now the functions:

$$(22) \quad h_i(x) = \delta_i G_i(x) + \beta_i \frac{\eta(G_i(x), a_i)}{\phi(G_{i-1}(x), G_i(x), a_i)} \quad , \quad x \in (x_{i-1}, +\infty) \quad , i = 2, \dots, n-2,$$

$$(23) \quad h_{n-1}(x) = \delta_{n-1} G_{n-1}(x) + h_{n-2}(x) \frac{\eta(G_{n-1}(x), a_{n-1})}{\eta(G_{n-2}(x), a_{n-1})} \quad , \quad x \in (x_{n-2}, +\infty) \quad ,$$

$$(24) \quad H_i(x) = \operatorname{erf} \frac{G_i(x)}{a_{i+1}} - \beta_{i+1} \frac{\eta(G_i(x), a_{i+1})}{h_i(x)} \quad , \quad x \in (x_{i-1}, +\infty) \quad , \quad i = 2, \dots, n-2,$$

$$(25) \quad G_i(x) = a_i \operatorname{erf}^{-1}[H_{i-1}(x)] \quad , \quad x \in (x_{i-1}, +\infty) \quad , \quad i = 3, \dots, n-1 \quad ,$$

where

$$(26) \quad x_i > x_{i-1} \quad / \quad H_i(x_i) = 0 \quad , \quad i = 2, \dots, n.$$

Having in mind that:

$$\frac{\beta_i}{\phi(G_{i-1}(x), G_i(x), a_i)} = \frac{h_{i-1}(x)}{\eta(G_{i-1}(x), a_i)} \quad , \quad i = 3, \dots, n-1,$$

it results that $G_{i-1}(x) > G_i(x)$, $\forall x \in (x_{i-1}, +\infty)$ and then the functions defined below are increasing functions and satisfy these conditions:

$$\begin{aligned} h_i(x_{i-1}) &> 0 \quad , & h_i(+\infty) &= +\infty, \\ H_i(x_{i-1}) &< 0 \quad , & H_i(x_i) &= 0 \quad , & H_i(+\infty) &= 1 \quad , \\ G_i(x_{i-1}) &= 0 \quad , & G_i(+\infty) &= +\infty. \end{aligned}$$

We can also define the function:

$$(27) \quad Q(x) = q_0 \eta(G_{n-1}(x), a_n) \quad , \quad x > x_{n-2},$$

which is a decreasing function and satisfies:

$$Q(x_{n-2}) = q_0, \quad Q(+\infty) = 0.$$

From (23), (25) and (27), the system ((12)a-b) can be written as:

$$(28) \quad \begin{aligned} \omega_i &= G_i(\omega_1), \quad i = 2, \dots, n-1, \\ h_{n-1}(\omega_1) &= Q(\omega_1), \end{aligned}$$

and consequently there will be a solution if and only if there exists an x^* such that $h_{n-1}(x^*) = Q(x^*)$. Note that because of the properties of functions h_{n-1} and Q this will happen if and only if $h_{n-1}(x_{n-2}) < Q(x_{n-2})$ or equivalently (on account of the properties of each G_i):

$$(29) \quad q_0 > \frac{\beta_{n-1}}{\operatorname{erf}\left(\frac{G_{n-2}(x_{n-2})}{\delta_{n-1}}\right)} > \frac{\beta_{n-2}}{\operatorname{erf}\left(\frac{G_{n-3}(x_{n-3})}{\delta_{n-2}}\right)} > \dots > \frac{\beta_2}{\operatorname{erf}\left(\frac{x_1}{\delta_2}\right)} > \beta_1.$$

Remark 2 : For the particular case $n = 2$ we find the result obtained in [3,4] for the existence of two phases.

IV. RELATIONSHIP WITH THE PROBLEM WITH CONSTANT TEMPERATURE ON THE FIXED FACE

The temperature (9) on the fixed face $x = 0$ is given by:

$$(30) \quad u^* = \theta_n(0, t) = A_n(\omega_{n-1}) = u_n + \operatorname{erf}\left(\frac{\omega_{n-1}}{\delta_n}\right) \frac{q_0 \delta_n \sqrt{\pi}}{k_n}.$$

As $u^* > u_n$ we can consider the multi-phase Stefan problem which consists in finding the functions $S_{n-1}(t) < S_{n-2}(t) < \dots < S_1(t)$ and the temperature $\Theta = \Theta(x, t)$, defined by:

$$\Theta(x, t) = \Theta_i(x, t) \quad \text{if} \quad S_i(t) < x < S_{i-1}(t), \quad t > 0, \quad i = 1, \dots, n,$$

solutions of (2)-(5), (6bis), (7) and (8) where

$$(6bis) \quad \Theta_n(0, t) = u_{n+1}, \quad \text{with } u_{n+1} > u_n.$$

The solution of this problem, following the above method, is given by:

$$(31) \quad \Theta_i(x,t) = M_i + N_i \operatorname{erf}\left(\frac{x}{2a_i\sqrt{t}}\right), \quad i = 1, \dots, n,$$

$$S_i(t) = 2\sigma_i\sqrt{t}, \quad \sigma_i > 0, \quad i = 1, \dots, n-1,$$

with

$$(32) \quad M_i = \frac{u_{i+1} \operatorname{erf}\left(\frac{\sigma_{i-1}}{a_i}\right) - u_i \operatorname{erf}\left(\frac{\sigma_i}{a_i}\right)}{\operatorname{erf}\left(\frac{\sigma_{i-1}}{a_i}\right) - \operatorname{erf}\left(\frac{\sigma_i}{a_i}\right)}, \quad N_i = \frac{u_i - u_{i+1}}{\operatorname{erf}\left(\frac{\sigma_{i-1}}{a_i}\right) - \operatorname{erf}\left(\frac{\sigma_i}{a_i}\right)}, \quad i=1, \dots, n,$$

where we have inserted two dummy parameters $\sigma_0 = +\infty$ y $\sigma_n = 0$.

The parameters σ_i are the solutions of the system of equations

$$(33) \quad \delta_i \sigma_i = \beta_{i+1} \frac{\eta(\sigma_i, a_{i+1})}{\phi(\sigma_i, \sigma_{i+1}, a_{i+1})} - \beta_i \frac{\eta(\sigma_i, a_i)}{\phi(\sigma_{i-1}, \sigma_i, a_i)}, \quad i = 1, \dots, n-1,$$

or equivalently:

$$\sigma_i = G_i(\sigma_1), \quad i = 2, \dots, n-1,$$

(34)

$$h_{n-1}(\sigma_1) = K(\sigma_1),$$

where the functions G_i and h_i ($i = 1, \dots, n-1$) are given by (16), (20), (22), (23) and (25) respectively and the function K is defined by:

$$(35) \quad K(x) = \beta_n \frac{\eta(G_{n-1}(x), a_n)}{\operatorname{erf}\left(\frac{G_{n-1}(x)}{a_n}\right)} \quad \text{with} \quad \beta_n = -\frac{k_n}{a_n\sqrt{\pi}}(u_n - u_{n+1}) > 0.$$

It satisfies $K(x_{n-2}) = +\infty$, $K(+\infty) = 0$ and $K'(x) < 0$, $\forall x \in (x_{n-2}, +\infty)$.

Because of

$$(36) \quad u_{n+1} = u_n + \operatorname{erf}\left(\frac{\omega_{n-1}}{a_n}\right) \frac{q_0 a_n \sqrt{\pi}}{k_n}$$

it results that $K(x) = Q(x)$ and from the uniqueness of the solution of (34) we deduce the equivalence of both problems, and then $M_i = A_i$, $N_i = B_i$, $\sigma_i = \omega_i$, $\forall i=1, \dots, n$.

Furthermore, if we consider $u^* = u_{n+1}$ it results

$$q_0 = -\frac{N_n k_n}{a_n \sqrt{\pi}} = \frac{\beta_n}{\operatorname{erf}\left(\frac{\sigma_{n-1}}{a_n}\right)},$$

hence the inequality (29), obtained for the first problem, is transformed in

$$(37) \quad \frac{\beta_n}{\operatorname{erf} \frac{\sigma_{n-1}}{a_n}} > \frac{\beta_{n-1}}{\operatorname{erf} \frac{G_{n-2}(x_{n-2})}{a_{n-1}}} > \frac{\beta_{n-2}}{\operatorname{erf} \frac{G_{n-3}(x_{n-3})}{a_{n-2}}} > \dots > \frac{\beta_2}{\operatorname{erf} \frac{x_1}{a_2}} > \beta_1.$$

for the second problem. Notice that this last inequality was obtained only with condition of temperature on the fixed face.

Remark 3 : The solution of problem (2)-(5), (6bis), (7) and (8) has been found before in [7] using a method similar to this.

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REFERENCES

- [1] H.S. CARSLAW, J.C. JAEGER, "Conduction of Heat in Solids", Clarendon Press, Oxford (1959).
- [2] L.I. RUBINSTEIN, "The Stefan Problem", Translations of Mathematical Monographs, Vol.27, American Mathematical Society, Providence (1971).
- [3] A.D. SOLOMON-D.G. WILSON-V. ALEXIADES, "Explicit solutions to change problems", Quart. Appl. Math., 41(1983), 237-243.
- [4] D.A. TARZIA, "An inequality for the coefficient σ of the free boundary $s(t)=2\sigma\sqrt{t}$ of the Neumann Solution for the two-phase Stefan Problem," Quart. Appl. Math., 39(1981/82), 491-497.
- [5] D.A. TARZIA "Soluciones exactas del problema de Stefan unidimensional", CUADERNOS del Instituto de Matemática "Beppo Levi", N°12, Rosario(1984), 5-36. See also "Analysis of a bibliography on moving and free boundary problems for the heat equation. Some results for the one-dimensional Stefan problem using the Lamé-Clapeyron and Neumann solutions", Research Notes in Math. N° 120, Pitman, London (1985), 84-102.
- [6] J.H. WEINER, "Transient heat conduction in multi-phase media", Brit. J. Appl. Physics, 6(1955), 361-363.
- [7] D.G. WILSON, "Existence and uniqueness for similarity solutions of one dimensional multi-phase Stefan problems", SIAM J. Appl. Math., 35(1978), 135-147.