# DETERMINATION OF UNKNOWN COEFFICENTS OF A SEMI-INFINITE MATERIAL THROUGH A SIMPLE MUSHY ZONE MODEL FOR THE TWO-PHASE STEFAN PROBLEM 

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#### Abstract

We use a simple mushy zone model in a two-phase solidification problem (Stefan problem) for the simultaneous determination of unknown coefficients of a semi-infinite material with an overspecified condition on the fixed face. We find the necessary and sufficient conditions for the existence of a solution and the corresponding formulae for the unknown coefficients. Copyright (C) 1996 Elsevier Science Ltd


## 1. INTRODUCTION

We consider a semi-infinite material with mass densities $\rho>0$ equal in both solid and liquid phases and we can assume, without loss of generality, that the phase-change temperature is $0^{\circ} \mathrm{C}$.

If the material is initially assumed to be liquid at the constant temperature $E>0$ and a constant temperature $-D<0$ is imposed on the fixed face $x=0$, then three distinct regions can be distinguished (for a mathematical and properties description of this simple model see [1]; for the one-phase model see [2]):
$\left(\mathrm{H}_{1}\right)$ The liquid phase, at temperature $\theta_{2}=\theta_{2}(x, t)>0$, occupying the region $x>r(t)$, $t>0$.
$\left(\mathrm{H}_{2}\right)$ The solid phase, at temperature $\theta_{1}=\theta_{1}(x, t)<0$, occupying the region $0<x<s(t)$, $t>0$.
$\left(\mathrm{H}_{3}\right)$ The mushy zone, at temperature 0 , occupying the region $s(t) \leq x \leq r(t), t>0$. We make two assumptions on its structure:
(a) The material in the mushy zone contains a fixed fraction $\epsilon \lambda$ (with $0<\epsilon<1$ ) of the total latent heat $\lambda>0$, i.e.

$$
\begin{equation*}
k_{1} \theta_{1_{r}}(s(t), t)-k_{2} \theta_{2_{s}}(r(t), t)=\lambda \rho(\epsilon \dot{s}(t)+(1-\epsilon) \dot{r}(t)), \quad t>0 . \tag{1.1}
\end{equation*}
$$

(b) The width of the mushy zone is inversely proportional (with constant $\gamma>0$ ) to the temperature gradient at the point $(s(t)$, $t$, i.e.

$$
\begin{equation*}
\theta_{1_{r}}(s(t), t)(r(t)-s(t))=\gamma, \quad t>0 . \tag{1.2}
\end{equation*}
$$

We suppose that the temperature $\theta=\theta(x, t)$ of the material is defined by

$$
\theta(x, t)= \begin{cases}\theta_{1}(x, t)<0 & \text { if } 0<x<s(t), t>0  \tag{1.3}\\ 0 & \text { if } s(t) \leq x \leq r(t), t>0 \\ \theta_{2}(x, t)>0 & \text { if } x>r(t), t>0\end{cases}
$$

The governing differential equations take the following forms for the solid and liquid phases:

$$
\begin{array}{ccc}
\alpha_{1} \theta_{1_{x t}}(x, t)=\theta_{1_{1}}(x, t), & 0<x<s(t), & t>0 \\
\alpha_{2} \theta_{2_{x t}}(x, t)=\theta_{2, t}(x, t), & x>r(t), & t>0 \tag{1.5}
\end{array}
$$

where $c_{i}>0, k_{i}>0$ and $\alpha_{i}=a_{i}^{2}=k_{i} / \rho c_{i}>0$ are the specific heat, the thermal conductivity and the diffusion coefficient for the phase $i(i=1$ denotes solid phase; $i=2$ denotes liquid phase) respectively.

The conditions at the solid-mushy interface $x=s(t)$ and the mushy-liquid interface $x=r(t)$ are given by (1.1), (1.2) and the requirement of the continuity of the temperature, i.e.

$$
\begin{equation*}
\theta_{1}(s(t), t)=\theta_{2}(r(t), t)=0, \quad t>0 \tag{1.6}
\end{equation*}
$$

The initial and boundary conditions are given by

$$
\begin{gather*}
\theta_{1}(0, t)=-D<0, \quad t>0  \tag{1.7}\\
\theta_{2}(x, 0)=\theta_{2}(+\infty, t)=E>0, \quad x>0, \quad t>0  \tag{1.8}\\
s(0)=r(0)=0 \tag{1.9}
\end{gather*}
$$

We consider an overspecified heat flux condition $[3,4]$ on the fixed face $x=0$ which is given by $[1,4-7]$

$$
\begin{equation*}
k_{1} \theta_{1_{x}}(0, t)=\frac{h_{\mathrm{o}}}{\sqrt{t}}, \quad t>0, \quad \text { with } \quad h_{\mathrm{o}}>0 \tag{1.10}
\end{equation*}
$$

If by means of a phase-change experiment we are able to measure certain quantities, then we shall find formulae for the simultaneous determination of the unknown coefficients ( $\epsilon, \gamma$ denote parameters of the mushy zone; $\lambda, \rho, c_{1}, c_{2}, k_{1}, k_{2}$ denote thermal coefficients of the material).

We shall also prove that the different problems for determining several unknown coefficients, posed in the next sections, do not always have an explicit solution. Moreover, it does exist iff some complementary conditions for the corresponding data are verified. In this paper, we generalize the results obtained in [5] for the particular case $\epsilon=1$ and $\gamma=0$ (i.e. without mushy region) and those obtained in [7] for the one-phase case. In [4] several references on free-moving boundary problems and determination of physical coefficients are given.

In Section 2 we shall consider the simple mushy zone model for the two-phase Stefan problem for determining one unknown thermal coefficient of a semi-infinite material with an overspecified condition on the fixed face, supposing the free boundaries $x=s(t)$ and $x=r(t)$ are unknown. The results obtained for the eight possible cases are considered in Appendix C (Table 1) which shows both the necessary and sufficient conditions to be verified by the data for the existence and uniqueness of the solution and the expression of the corresponding unknown coefficient. Moreover, we shall also prove the respective properties for the determination of $\epsilon$ (case 3) and the determination of $k_{2}$ (case 7).

In Section 3 we shall consider the same model for determining two unknown thermal coefficients of a semi-infinite material with an overspecified condition on the fixed face, supposing known the expression for the moving boundary $x=s(t)$. The results obtained for the 28 possible cases are considered in Appendix D (Table 2) which shows both the necessary and sufficient conditions to be verified by the data for the existence of the solution and the expression of the corresponding unknown coefficients. There are several cases where the moving boundary problem has a unique solution iff some conditions are verified. Moreover, we
shall also prove the respective properties for the determination of $k_{1}$ and $k_{2}$ (case 9), the determination of $\epsilon$ and $k_{2}$ (case 17), the determination of $c_{2}$ and $k_{2}$ (case 19) and, the determination of $\gamma$ and $k_{2}$ (case 23).

The functions and the restrictions used in the text and, Appendices C and D are summarized in Appendix A and Appendix B respectively.

## 2. DETERMINATION OF ONE UNKNOWN THERMAL COEFFICIENT

Taking into account the hypotheses $\left(\mathrm{H}_{1}\right)-\left(\mathrm{H}_{3}\right)$ we can formulate the following:
Problem ( $\mathrm{P}_{1}$ ). Find the free boundaries $x=s(t)$ and $x=r(t)$, defined for $t>0$ with $0<s(t)<$ $r(t)$ and $s(0)=r(0)=0$, the temperature $\theta=\theta(x, t)$, defined by (1.3) for $x>0$ and $t>0$, and one of the eight unknown thermal coefficients $\epsilon, \gamma, \lambda, \rho, c_{1}, c_{2}, k_{1}, k_{2}$ such that they satisfy the conditions (1.1), (1.2), (1.4)-(1.10) where $D>0, E>0$ and $h_{0}>0$ are data and they must be known or determined by an experience of phase-change [8].

The solution of this problem is given $[1,6,9-11]$ by

$$
\begin{gather*}
\theta_{1}(x, t)=-D+\frac{D}{f\left(\frac{\sigma}{a_{1}}\right)} f\left(\frac{x}{2 a_{1} \sqrt{t}}\right)  \tag{2.1}\\
\theta_{2}(x, t)=\frac{-E f\left(\frac{\omega}{a_{2}}\right)}{1-f\left(\frac{\omega}{a_{2}}\right)}+\frac{E}{1-f\left(\frac{\omega}{a_{2}}\right)} f\left(\frac{x}{2 a_{2} \sqrt{t}}\right)  \tag{2.2}\\
s(t)=2 \sigma \sqrt{t}, \quad \sigma>0  \tag{2.3}\\
r(t)=2 \omega \sqrt{t}, \quad \omega>\sigma \tag{2.4}
\end{gather*}
$$

where $f$ is the error function, the coefficient $\omega$ is given by

$$
\begin{equation*}
\omega=\omega(\sigma)=a_{1} W\left(\frac{\sigma}{a_{1}}\right) \tag{2.5}
\end{equation*}
$$

and the coefficient $\sigma$ and the unknown thermal coefficient are obtained by solving the following system of equations:

$$
\begin{gather*}
\frac{h_{\mathrm{o}}}{\lambda \rho a_{1}} \exp \left(-\frac{\sigma^{2}}{a_{1}^{2}}\right)-\frac{E k_{2}}{\lambda \rho a_{1} a_{2} \sqrt{\pi}} F_{1}\left(\frac{\omega(\sigma)}{a_{1}}\right)=G\left(\frac{\sigma}{a_{1}}\right)  \tag{2.6a}\\
\frac{a_{1}}{k_{1}} f\left(\frac{\sigma}{a_{1}}\right)=\frac{D}{h_{o} \sqrt{\pi}} . \tag{2.6~b}
\end{gather*}
$$

The eight possible cases for Problem ( $\mathrm{P}_{1}$ ) are considered in Appendix C (Table 1) which shows both the necessary and sufficient conditions to be verified by the data for the existence and uniqueness of the solution of the problem and the expression of the coefficient $\sigma$ together with the corresponding unknown coefficient. We remark here that the coefficient $\omega$ is always given by the expression (5) as a function of $\sigma$ and $a_{1}$.

Now, we shall prove the following properties only for the determination of $\epsilon$ (case 3 ) and the determination of $k_{2}$ (case 7), which gave us different difficulties in all cases.

Theorem 1 (Case 3). The necessary and sufficient condition for Problem $\left(\mathrm{P}_{1}\right)$, with $\sigma$ and $\epsilon$ unknown, to have a unique solution is that data $D>0, E>0, h_{\mathrm{o}}>0$, mushy zone coefficient $\gamma>0$ and thermal coefficients of the phase-change material $\lambda, \rho, c_{1}, c_{2}, k_{1}, k_{2}>0$ do verify the conditions

$$
\begin{equation*}
h_{\mathrm{o}}>\frac{E k_{2}}{a_{2} \sqrt{\pi}}, \quad f\left(x_{5}\right)<\frac{D k_{1}}{h_{\mathrm{o}} a_{1} \sqrt{\pi}}<f\left(x_{4}\right) \tag{2.7}
\end{equation*}
$$

where $x_{4}$ and $x_{5}$ are unique positive zeros of functions $H_{4}$ and $H_{5}$ respectively. In such a case, the solution is given by (2.1)-(2.4) with

$$
\begin{equation*}
\epsilon=\frac{2 D}{\gamma \sqrt{\pi}} F_{2}\left(\xi_{1}\right) H_{5}\left(\xi_{1}\right), \quad \sigma=a_{1} \xi_{1}, \quad \omega=a_{1} W\left(\xi_{1}\right) \tag{2.8}
\end{equation*}
$$

where $\xi_{1}$ is the unique solution of the equation

$$
\begin{equation*}
f(x)=\frac{D k_{1}}{h_{0} a_{1} \sqrt{\pi}}, \quad x>0 \tag{2.9}
\end{equation*}
$$

Proof. We define

$$
\begin{equation*}
\xi_{1}=\frac{\sigma}{a_{1}}, \quad \text { with } \quad a_{1}=\sqrt{\frac{k_{1}}{\rho c_{1}}} \tag{2.10}
\end{equation*}
$$

The coefficient $\sigma$ is obtained from (2.10) and the element $\xi_{1}$ is given from (2.6b) as the solution of (2.9) iff the data verify the condition

$$
\begin{equation*}
\frac{D k_{1}}{h_{\mathrm{o}} a_{1} \sqrt{\pi}}<1 \tag{2.11}
\end{equation*}
$$

From (2.6a) it follows that $\xi_{1}$ should verify

$$
\begin{equation*}
\frac{h_{\mathrm{o}}}{\lambda \rho a_{1}} \exp \left(-\xi_{1}^{2}\right)-\frac{E k_{2}}{\lambda \rho a_{1} a_{2} \sqrt{\pi}} F_{1}\left(\frac{a_{1}}{a_{2}} W\left(\xi_{1}\right)\right)=\xi_{1}+\frac{(1-\epsilon) \gamma \sqrt{\pi}}{2 D} f\left(\xi_{1}\right) \exp \left(\xi_{1}^{2}\right) \tag{2.12}
\end{equation*}
$$

then we obtain the expression for $\epsilon$ in (2.8).
Therefore we have the following properties:

$$
\begin{align*}
& \epsilon<1 \text { iff } H_{4}\left(\xi_{1}\right)>0 \\
&  \tag{2.13}\\
& \\
& \text { iff } \quad H_{4}\left(0^{+}\right)>0\left(\text { i.e. } h_{\mathrm{o}}>\frac{E k_{2}}{a_{2} \sqrt{\pi}}\right) \text { and } \xi_{1}<x_{4}\left(\text { i.e. } f\left(\xi_{1}\right)<f\left(x_{4}\right)\right),
\end{align*}
$$

where $x_{4}$ is the only positive root of $H_{4}$ (because $H_{4}$ is a decreasing function for $x>0$ ), and

$$
\begin{equation*}
\left.\epsilon>0 \quad \text { iff } \quad H_{5}\left(\xi_{1}\right)>0 \quad \text { iff } \quad \xi_{1}>x_{5} \text { (i.e. } f\left(\xi_{1}\right)>f\left(x_{5}\right)\right) \tag{2.14}
\end{equation*}
$$

where $x_{5}$ is the only positive root of $H_{5}$ (because $H_{5}$ is an increasing function for $x>0$, $H_{5}\left(0^{+}\right)=-\frac{h_{\mathrm{o}}}{\lambda \rho a_{1}} \alpha_{20}<0$ and $\left.H_{5}(+\infty)=+\infty\right)$. We can deduce (2.7) from (2.13) and (2.14) because $x_{5}<x_{4}$.

Theorem 2. (Case 7). The necessary and sufficient condition for Problem $\left(\mathrm{P}_{1}\right)$, with $\sigma$ and $k_{2}$ unknown, to have a unique solution is that data $D>0, E>0, h_{\mathrm{o}}>0$, mushy zone coefficients $0<\epsilon<1$ and $\gamma>0$, and thermal coefficients of the phase-change material $\lambda, \rho, c_{1}, c_{2}, k_{1}>0$ do verify the condition

$$
\begin{equation*}
\frac{D k_{1}}{h_{\mathrm{o}} a_{1} \sqrt{\pi}}<f\left(x_{17}\right) \tag{2.15}
\end{equation*}
$$

where $x_{17}$ is the unique positive zero of function $H_{17}$. In such a case, the solution is given by (2.1)-(2.4) with

$$
\begin{equation*}
k_{2}=\frac{k_{1} c_{2}}{c_{1}} \frac{W^{2}\left(\xi_{1}\right)}{B^{2}}, \quad \sigma=a_{1} \xi_{1}, \quad \omega=a_{1} W\left(\xi_{1}\right) \tag{2.16}
\end{equation*}
$$

where $\xi_{1}$ is the unique solution of (2.9) and $B$ is the only solution of the equation

$$
\begin{equation*}
\frac{1}{H_{16}(x)}=\frac{\lambda}{E c_{2}} \frac{H_{2}\left(\xi_{1}\right)}{W\left(\xi_{1}\right)}, \quad x>0 \tag{2.17}
\end{equation*}
$$

Proof. We obtain the coefficient $\sigma$ as in Theorem 1. From (2.6a) it follows that $\xi_{1}$ should verify

$$
\begin{equation*}
\frac{E k_{2}}{\lambda \rho a_{1} a_{2} \sqrt{\pi}} F_{1}\left(\frac{a_{1}}{a_{2}} W\left(\xi_{1}\right)\right)=H_{2}\left(\xi_{1}\right) \tag{2.18}
\end{equation*}
$$

If we define

$$
\begin{equation*}
B=\frac{a_{1}}{a_{2}} W\left(\xi_{1}\right), \quad \text { with } \quad a_{2}=\frac{\sqrt{k_{2}}}{\sqrt{\rho c_{2}}}, \tag{2.19}
\end{equation*}
$$

then equation (2.18) is equivalent to

$$
\begin{equation*}
\frac{F_{1}(B)}{B}=\frac{\lambda \sqrt{\pi}}{E c_{2}} \frac{H_{2}\left(\xi_{1}\right)}{W\left(\xi_{1}\right)}, \quad B>0 \tag{2.20}
\end{equation*}
$$

that is, $B$ is the solution of (2.17). Taking into account the properties of the function $H_{16}$ we can deduce that there exists a unique solution of (2.17) if and only if

$$
\begin{equation*}
H_{17}\left(\xi_{1}\right)>0 \quad \text { iff } \quad \xi_{1}<x_{17}(\text { i.e. }(2.15)) \tag{2.21}
\end{equation*}
$$

where $x_{17}$ is the only positive root of $H_{17}$ (because $H_{17}$ is a decreasing function for $x>0$, $H_{17}\left(0^{+}\right)>0$ and $\left.H_{17}(+\infty)=-\infty\right)$. From (2.19) we obtain the coefficient $k_{2}$.

## 3. DETERMINATION OF TWO UNKNOWN THERMAL COEFFICIENTS

Taking into account the hypotheses $\left(\mathrm{H}_{1}\right)-\left(\mathrm{H}_{3}\right)$ we can formulate the following:
Problem $\left(\mathrm{P}_{2}\right)$. Find the free boundary $x=r(t)$, defined for $t>0$ with $r(0)=0$, the temperature $\theta=\theta(x, t)$, defined by (1.3) for $x>0$ and $t>0$, and two of the eight unknown thermal coefficients $\epsilon, \gamma, \lambda, \rho, c_{1}, c_{2}, k_{1}, k_{2}$ such that they satisfy the conditions (1.1), (1.2), (1.4)-(1.10) where the moving boundary $x=s(t)$, defined for $t>0$ with $s(0)=0$, is given by (2.3) with $a$ known coefficient $\sigma>0$ and $D, E, h_{0}>0$ are data and they must be known or determined by an experience of phase-change [8].

The solution of that problem is given by (2.1), (2.2) and (2.4) where the coefficient $\omega$ and the unknown thermal coefficients are obtained by solving the system of equations (2.6).

The 28 cases for Problem ( $\mathrm{P}_{2}$ ) (cases 9 to 36) are considered in Appendix D (Table 2) which shows both the necessary and sufficient conditions to be verified by the data for the existence of the solution of the problem and the expression of the coefficient $\omega$ together with the corresponding unknown coefficients. There are several cases where the moving boundary problem has a unique solution iff some conditions are verified.

Now, we shall prove the properties corresponding only for the determination of $k_{1}$ and $k_{2}$ (case 9), the determination of $\epsilon$ and $k_{2}$ (case 17), the determination of $c_{2}$ and $k_{2}$ (case 19) and the determination of $\gamma$ and $k_{2}$ (case 23), which give us different difficulties in all cases.

Theorem 3 (Case 9). The necessary and sufficient condition for Problem $\left(\mathrm{P}_{2}\right)$, with $\omega, k_{1}$ and $k_{2}$ unknown, to have a unique solution is that data $\sigma>0, D>0, E>0, h_{\mathrm{o}}>0$, mushy zone coefficients $0<\epsilon<1$ and $\gamma>0$, and thermal coefficients of the phase-change material $\lambda, \rho, c_{1}, c_{2}>0$ do verify the conditions

$$
\begin{equation*}
\frac{h_{\mathrm{o}}}{E \rho \sigma c_{2}}>1+\frac{\gamma}{D}+\frac{\lambda}{E c_{2}}\left(1+\frac{(1-\varepsilon) \gamma}{D}\right), \quad \frac{D \rho \sigma c_{1}}{h_{\mathrm{o}} \sqrt{\pi}}<H_{20}\left(x_{23}\right) \tag{3.1}
\end{equation*}
$$

where $x_{23}$ is the unique positive zero of function $H_{23}$. In such a case, the solution is given by (2.1), (2.2) and (2.4) with

$$
\begin{equation*}
\omega=\sigma H_{25}\left(\xi_{1}\right), \quad k_{1}=\rho \sigma^{2} c_{1} \frac{1}{\xi_{1}^{2}}, \quad k_{2}=\rho \sigma^{2} c_{2} \frac{H_{25}^{2}\left(\xi_{1}\right)}{B^{2}} \tag{3.2}
\end{equation*}
$$

where $\xi_{1}$ is the unique solution of the equation

$$
\begin{equation*}
H_{20}(x)=\frac{D \rho \sigma c_{1}}{h_{\mathrm{o}} \sqrt{\pi}}, \quad x>0 \tag{3.3}
\end{equation*}
$$

and $B$ is the only solution of the equation

$$
\begin{equation*}
\frac{1}{H_{16}(x)}=\frac{\lambda}{E c_{2}} \frac{H_{21}\left(\xi_{1}\right)}{W\left(\xi_{1}\right)}, \quad x>0 \tag{3.4}
\end{equation*}
$$

Proof. We define

$$
\begin{equation*}
\xi_{1}=\frac{\sigma}{a_{1}}, \quad \text { with } \quad a_{1}=\frac{\sqrt{k_{1}}}{\sqrt{\rho c_{1}}} \tag{3.5}
\end{equation*}
$$

The coefficients $\omega$ and $k_{1}$ are obtained using (3.5) and the element $\xi_{1}$ is given from (2.6b), as the solution of (3.3). From (2.6a) it follows that $\xi_{1}$ and $k_{2}$ should verify

$$
\begin{equation*}
\frac{E k_{2}}{\lambda \rho a_{1} a_{2} \sqrt{\pi}} F_{1}\left(\frac{\sigma}{a_{2}} \frac{W\left(\xi_{1}\right)}{\xi_{1}}\right)=H_{21}\left(\xi_{1}\right) . \tag{3.6}
\end{equation*}
$$

If we define

$$
\begin{equation*}
B=\frac{\sigma}{a_{2}} \frac{W\left(\xi_{1}\right)}{\xi_{1}}=\frac{\sigma}{a_{2}} H_{25}\left(\xi_{1}\right), \quad \text { with } \quad a_{2}=\sqrt{\frac{k_{2}}{\rho c_{2}}} \tag{3.7}
\end{equation*}
$$

then (3.6) is equivalent to

$$
\begin{equation*}
\frac{F_{1}(B)}{B}=\frac{\lambda \sqrt{\pi}}{E c_{2}} \frac{H_{21}\left(\xi_{1}\right)}{W\left(\xi_{1}\right)}, \quad B>0 \tag{3.8}
\end{equation*}
$$

that is, $B$ is the solution of (3.4). Taking into account the properties of the function $H_{16}$ we can deduce that there exists a unique solution of (3.4) if and only if

$$
H_{23}\left(\xi_{1}\right)>0 \quad \text { iff } \quad H_{23}\left(0^{+}\right)>0 \quad \text { and } \quad \xi_{1}<x_{23} \text { (i.e. (3.1)), }
$$

where $x_{23}$ is the only positive root of $H_{23}$ (because $H_{23}$ is a decreasing function for $x>0$ and $\left.H_{23}(+\infty)=-\infty\right)$. From (3.7) we obtain the coefficient $k_{2}$.

Theorem 4 (Case 17). The necessary and sufficient condition for Problem $\left(\mathrm{P}_{2}\right)$, with $\omega, \epsilon$ and $\boldsymbol{k}_{2}$ unknown, to have at least one solution is that data $\sigma>0, D>0, E>0, h_{0}>0$, mushy zone coefficient $\gamma>0$, and thermal coefficients of the phase-change material $\lambda, \rho, c_{1}, c_{2}, k_{1}>0$ do verify the conditions

$$
\begin{equation*}
h_{\mathrm{o}}=\frac{D k_{1}}{a_{1} f\left(\frac{\sigma}{a_{1}}\right) \sqrt{\pi}}, \quad \xi_{1}=\frac{\sigma}{a_{1}}<x_{36} \tag{3.9}
\end{equation*}
$$

where $x_{36}$ is the unique positive zero of function $H_{36}$. In such a case, there exist infinite solutions which have the form (2.1), (2.2) and (2.4) where

$$
\begin{gather*}
\omega=a_{1} W\left(\xi_{1}\right), \quad k_{2}=\frac{k_{1} c_{2}}{c_{1}} \frac{W^{2}\left(\xi_{1}\right)}{B^{2}}, \\
\epsilon=\frac{2 D}{\gamma \sqrt{\pi}} F_{2}\left(\xi_{1}\right)\left(W\left(\xi_{1}\right)-\frac{D c_{1}}{\lambda \sqrt{\pi}} F_{2}\left(\xi_{1}\right)+\frac{E c_{2}}{\lambda} \frac{W\left(\xi_{1}\right)}{H_{16}(B)}\right), \tag{3.10}
\end{gather*}
$$

with $B$ an arbitrary parameter which is defined by

$$
\begin{equation*}
B>H_{16}^{-1}\left(\frac{1}{A}\right) \quad \text { if } \quad x_{37} \leq \xi_{1}<x_{36}\left(A=\frac{\lambda}{E c_{2}} \frac{1}{W\left(\xi_{1}\right)}\left(\frac{D c_{1}}{\lambda \sqrt{\pi}} F_{2}\left(\xi_{1}\right)-\xi_{1}\right)\right), \tag{3.11}
\end{equation*}
$$

and

$$
\begin{equation*}
H_{16}^{-1}\left(\frac{1}{A}\right)<B<H_{16}^{-1}\left(\frac{1}{C}\right) \quad \text { if } \quad \xi_{1}<x_{37}\left(C=\frac{\lambda}{E c_{2}} \frac{1}{W\left(\xi_{1}\right)}\left(\frac{D c_{1}}{\lambda \sqrt{\pi}} F_{2}\left(\xi_{1}\right)-W\left(\xi_{1}\right)\right)\right) \tag{3.12}
\end{equation*}
$$

where $x_{36}$ and $x_{37}$ are the unique positive zeros of functions $H_{36}$ and $H_{37}$ respectively.
Proof. From (2.6b) it follows that $h_{\mathrm{o}}=\frac{D k_{1}}{a_{1} f\left(\xi_{1}\right) \sqrt{\pi}}$. From (2.6a), $\epsilon$ and $k_{2}$ should verify

$$
\begin{equation*}
\frac{D c_{1}}{\lambda \sqrt{\pi}} F_{2}\left(\xi_{1}\right)-\frac{E}{\lambda a_{1}} \sqrt{\frac{c_{2} k_{2}}{\rho \pi}} F_{1}\left(a_{1} \sqrt{\frac{\rho c_{2}}{k_{2}}} W\left(\xi_{1}\right)\right)=\xi_{1}+\frac{(1-\varepsilon) \gamma \sqrt{\pi}}{2 D} f\left(\xi_{1}\right) \exp \left(\xi_{1}\right) \tag{3.13}
\end{equation*}
$$

i.e. the expression for $\epsilon$ in (3.10), when we define $B=a_{1} \sqrt{\frac{\rho c_{2}}{k_{2}}} W\left(\xi_{1}\right)$. Then we have

$$
\begin{equation*}
\epsilon<1 \quad \text { iff } \frac{1}{H_{16}(B)}<A \text { and } \epsilon<0 \text { iff } \frac{1}{H_{16}(B)}>C . \tag{3.14}
\end{equation*}
$$

Taking into account the properties of the function $H_{16}$ we can deduce

$$
\begin{align*}
& A>1 \Leftrightarrow H_{36}\left(\xi_{1}\right)>0 \Leftrightarrow \xi_{1}<x_{36}  \tag{3.15}\\
& C>1 \Leftrightarrow H_{37}\left(\xi_{1}\right)>0 \Leftrightarrow \xi_{1}<x_{37} .
\end{align*}
$$

Then we obtain (3.11) and (3.12) from (3.14) and (3.15).

Theorem 5 (Case 19). The necessary and sufficient condition for Problem ( $\mathrm{P}_{2}$ ), with $\omega, c_{2}$ and $k_{2}$ unknown, to have at least one solution is that data $\sigma>0, D>0, E>0, h_{0}>0$, mushy zone coefficients $0<\varepsilon<1$ and $\gamma>0$, and thermal coefficients of the phase-change material $\lambda, \rho, c_{1}, k_{1}>0$ do verify the conditions

$$
\begin{equation*}
h_{\mathrm{o}}=\frac{D k_{1}}{a_{1} f\left(\frac{\sigma}{a_{1}}\right) \sqrt{\pi}}, \quad H_{57}\left(\frac{\sigma}{a_{1}}\right)>0 \quad\left(\text { or } \xi_{1}=\frac{\sigma}{a_{1}}<x_{57}\right) \tag{3.16}
\end{equation*}
$$

where $x_{57}$ is the unique positive zero of function $H_{57}$. In such a case, there exist infinite solutions which have the form (2.1), (2.2) and (2.4) where

$$
\begin{equation*}
\omega=a_{1} W\left(\xi_{1}\right), \quad c_{2}=\frac{\lambda}{E} \frac{H_{16}(B) H_{57}\left(\xi_{1}\right)}{W\left(\xi_{1}\right)}, \quad k_{2}=\frac{\lambda k_{1} \sqrt{\pi}}{E c_{1}} \frac{W\left(\xi_{1}\right) H_{57}\left(\xi_{1}\right)}{H_{1}(B)} \tag{3.17}
\end{equation*}
$$

for any $B>0$.
Proof. From (2.6b) it follows that $h_{\mathrm{o}}=\frac{D k_{1}}{a_{1} f\left(\xi_{1}\right) \sqrt{\pi}}$. We define

$$
\begin{equation*}
B=\sqrt{\frac{k_{1} c_{2}}{c_{1} k_{2}}} W\left(\xi_{1}\right) . \tag{3.18}
\end{equation*}
$$

From (2.6a), $B$ and $k_{2}$ should verify

$$
\begin{equation*}
\frac{E c_{1} k_{2}}{\lambda k_{1} \sqrt{\pi}} \frac{H_{1}(B)}{W\left(\xi_{1}\right)}=H_{57}\left(\xi_{1}\right) \tag{3.19}
\end{equation*}
$$

i.e. we deduce the expression for $k_{2}$ in (3.17). Then we obtain the coefficient $c_{2}$ from (3.18) and (3.19).

Thus we have

$$
\begin{equation*}
k_{2}>0 \Leftrightarrow H_{57}\left(\xi_{1}\right)>0 \Leftrightarrow 0<\xi_{1}<x_{57} \tag{3.20}
\end{equation*}
$$

where $x_{57}$ is the only positive root of $H_{57}$ (because $H_{57}$ is a decreasing function for $x>0$, with $H_{57}\left(0^{+}\right)=+\infty$ and $\left.H_{57}(+\infty)=-\infty\right)$.

Theorem 6 (Case 23). The necessary and sufficient condition for Problem $\left(\mathrm{P}_{2}\right)$, with $\omega, \gamma$ and $k_{2}$ unknown, to have at least one solution is that data $\sigma>0, D>0, E>0, h_{0}>0$, mushy zone coefficient $0<\varepsilon<1$, and thermal coefficients of the phase-change material $\lambda, \rho, c_{1}, c_{2}, k_{1}>0$ do verify the conditions

$$
\begin{equation*}
h_{\mathrm{o}}=\frac{D k_{1}}{a_{1} f\left(\frac{\sigma}{a_{1}}\right) \sqrt{\pi}}, \quad H_{50}\left(\frac{\sigma}{a_{1}}\right)>0 \quad\left(\text { or } \xi_{1}=\frac{\sigma}{a_{1}}<x_{50}\right) \tag{3.21}
\end{equation*}
$$

where $x_{50}$ is the unique positive zero of function $H_{50}$. In such a case, there exist infinite solutions which have the form (2.1), (2.2) and (2.4) where

$$
\begin{equation*}
\omega=a_{1} W\left(\xi_{1}, \gamma\right), \quad k_{2}=\frac{k_{1} c_{2}}{c_{1}} \frac{W^{2}\left(\xi_{1}, \gamma\right)}{B^{2}} \tag{3.22}
\end{equation*}
$$

for any $0<\gamma<\gamma_{\mathrm{o}}$, with

$$
\begin{equation*}
\gamma_{\mathrm{o}}=\frac{2 D}{\left(1-\epsilon+\frac{E c_{2}}{\lambda}\right) \sqrt{\pi}} F_{2}\left(\xi_{1}\right) H_{50}\left(\xi_{1}\right) \tag{3.23}
\end{equation*}
$$

and $B=B(\gamma)$ is the only solution of the equation

$$
\begin{equation*}
\frac{1}{H_{16}(x)}=\frac{\lambda}{E c_{2}} \frac{H_{57}\left(\xi_{1}, \gamma\right)}{W\left(\xi_{1}, \gamma\right)}, \quad x>0 . \tag{3.24}
\end{equation*}
$$

Proof. From (2.6b) it follows that $h_{\mathrm{o}}=\frac{D k_{1}}{(\sigma)}$. We define

$$
a_{1} f\left(\frac{\sigma}{a_{1}}\right) \sqrt{\pi}
$$

$$
\begin{equation*}
\xi_{2}=\frac{\sigma}{a_{2}}, \quad \text { with } \quad a_{2}=\sqrt{\frac{k_{2}}{\rho c_{2}}} . \tag{3.25}
\end{equation*}
$$

From (2.6a), $\xi_{2}>0$ should verify

$$
\begin{equation*}
F_{1}\left(\frac{W\left(\xi_{1}, \gamma\right)}{\xi_{1}} \xi_{2}\right) \frac{1}{\xi_{2}}=\frac{\lambda \sqrt{\pi}}{E c_{2}} \frac{1}{\xi_{1}}\left(\frac{h_{\mathrm{o}}}{\lambda \rho \sigma} \xi_{1} \exp \left(-\xi_{1}^{2}\right)-G_{1}\left(\xi_{1}, \gamma\right)\right) . \tag{3.26}
\end{equation*}
$$

If we define

$$
\begin{equation*}
B=B(\gamma)=\frac{W\left(\xi_{1}, \gamma\right)}{\xi_{1}} \xi_{2}, \tag{3.27}
\end{equation*}
$$

equation (3.26) for $\xi_{2}$ is equivalent to (3.24) for $B$. We obtain the coefficients $\omega$ and $k_{2}$ from (3.25) and (3.27).

Taking into account the properties of the function $H_{16}$ we can deduce that there exists a unique solution of (3.24) if and only if

$$
\begin{equation*}
H_{50}\left(\xi_{1}\right)>0 \quad \text { (or equivalently } \xi_{1}<x_{50} \text { ), } \tag{3.28}
\end{equation*}
$$

for any $0<\gamma<\gamma_{0}$, where $x_{50}$ is the only positive root of $H_{50}$ (because $H_{50}$ is a decreasing function for $x>0$ with $H_{50}\left(0^{+}\right)=+\infty$ and $\left.H_{50}(+\infty)=-\infty\right)$.

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## APPENDIX A

The following real functions are defined, for $x>0$, by

$$
\begin{aligned}
& f(x)=\operatorname{erf}(x)=\frac{2}{\sqrt{\pi}} \int_{0}^{x} \exp \left(-t^{2}\right) \mathrm{d} t \\
& F_{2}(x)=\frac{\exp \left(-x^{2}\right)}{f(x)} \\
& G(x)=G(x, \gamma)=x+\frac{(1-\epsilon) \gamma \sqrt{\pi}}{2 D} f(x) \exp \left(x^{2}\right) \\
& H_{3}(x)=\exp \left(-x^{2}\right)-\frac{E k_{2}}{h_{\mathrm{o}} a_{2} \sqrt{\pi}} F_{1}\left(\frac{a_{1}}{a_{2}} W(x)\right) \\
& H_{4}(x)=\frac{h_{\mathrm{o}}}{\lambda \rho a_{1}} \exp \left(-x^{2}\right)-x-\frac{E k_{2}}{\lambda \rho a_{1} a_{2} \sqrt{\pi}} F_{1}\left(\frac{a_{1}}{a_{2}} W(x)\right) \\
& H_{5}(x)=\frac{\gamma^{\sqrt{\pi}}}{2 D} f(x) \exp \left(x^{2}\right)-H_{4}(x) \\
& H_{6}(x)=(1-\epsilon) \frac{a_{2}}{a_{1}} x+\frac{E k_{2}}{\lambda \rho a_{1} a_{2} \sqrt{\pi}} F_{1}(x) \\
& H_{\mathrm{s}}(x)=\frac{W(x)}{f(x)} \\
& H_{7}(x)=\frac{h_{0}}{\lambda \rho a_{1}} \exp \left(-x^{2}\right)-x-\frac{E k_{2}}{\lambda \rho a_{1} a_{2} \sqrt{\pi}} F_{1}\left(\frac{a_{1}}{a_{2}} x\right) \\
& H_{9}(x)=\frac{G(x)}{f(x)} \\
& H_{10}(x)=\frac{h_{\mathrm{o}}^{2} \sqrt{\pi}}{D \lambda \rho k_{1}} \exp \left(-x^{2}\right)-\frac{E h_{\mathrm{o}} k_{2}}{D \lambda \rho k_{1} a_{2}} F_{1}\left(\frac{D k_{1}}{h_{\mathrm{o}} a_{2} \sqrt{\pi}} H_{8}(x)\right) \\
& H_{11}(x)=\frac{\beta_{2}}{\beta_{1}} x+F_{1}(x) \\
& H_{12}(x)=\beta_{1} \beta_{3} \frac{1}{x} \\
& \beta_{1}=\frac{D k_{1}}{h_{\mathrm{o}} a_{2} \sqrt{\pi}}\left(\frac{2}{\sqrt{\pi}}+\frac{\gamma \sqrt{\pi}}{2 D}\right) \\
& H_{13}(x)=f(x) W(x) \\
& \beta_{2}=\frac{D \lambda \rho k_{1} a_{2}}{E h_{\mathrm{o}} k_{2}}\left(\frac{2}{\sqrt{\pi}}+\frac{(1-\epsilon) \gamma \sqrt{\pi}}{2 D}\right) \\
& \beta_{3}=\frac{h_{0} a_{2} \sqrt{\pi}}{E k_{2}} \\
& H_{14}(x)=f(x) G(x) \\
& H_{15}(x)=\frac{D c_{1}}{\lambda \sqrt{\pi}} \exp \left(-x^{2}\right)-\frac{D E c_{1} k_{2}}{\lambda \pi h_{\mathrm{o}} a_{2}} F_{1}\left(\frac{h_{0} \sqrt{\pi}}{D \rho c_{1} a_{2}} H_{13}(x)\right) \\
& H_{16}(x)=\sqrt{\pi} x \exp \left(x^{2}\right)(1-f(x)) \\
& H_{17}(x)=\frac{h_{\mathrm{o}}}{E \rho a_{1} c_{2}} \exp \left(-x^{2}\right)-\left(1+\frac{(1-\epsilon) \lambda}{E c_{2}}\right) \\
& \times \frac{y \sqrt{\pi}}{2 D} f(x) \exp \left(x^{2}\right)-\left(1+\frac{\lambda}{E c_{2}}\right) x \\
& H_{19}(x)=f(x) H_{\text {Is }}(x) \\
& H_{20}(x)=x f(x) \\
& H_{2!}(x)=H_{21}(x, \gamma)=\frac{h_{0}}{\lambda \rho \sigma} x \exp \left(-x^{2}\right)-G(x, \gamma) \\
& H_{22}(x)=\frac{h_{\mathrm{o}}}{E \rho \sigma c_{2}} x \exp \left(-x^{2}\right)-\left(1+\frac{(1-\epsilon) \lambda}{E c_{2}}\right) \\
& \times \frac{\gamma \sqrt{\pi}}{2 D} f(x) \exp \left(x^{2}\right)-\left(1+\frac{\lambda}{E c_{2}}\right) x \\
& H_{23}(x)=\frac{H_{22}(x)}{x} \\
& H_{24}(x)=\frac{f(x)}{x} \\
& H_{25}(x)=\frac{W(x)}{x} \\
& H_{27}(x)=\frac{H_{26}(x)}{x} \\
& H_{26}(x)=x \exp \left(-x^{2}\right)-\frac{E k_{2}}{h_{0} a_{2} \sqrt{\pi}} x F_{1}\left(\frac{\sigma}{a_{2}} H_{25}(x)\right) \\
& H_{28}(x)=\frac{H_{21}(x)}{x} \\
& H_{29}(x)=\frac{h_{0}}{\lambda \rho \sigma} \exp \left(-x^{2}\right)-1-\frac{E k_{2}}{\lambda \rho \sigma a_{2} \sqrt{\pi}} F_{1}\left(\frac{\sigma}{a_{2}} H_{25}(x)\right) \\
& H_{31}(x)=\frac{\gamma \sqrt{\pi}}{2 D} \frac{1}{x F_{2}(x)}-H_{29}(x) \\
& H_{31}(x)=(1-\epsilon) \frac{a_{2}}{\sigma} x+\frac{E k_{2}}{\lambda \rho \sigma a_{2} \sqrt{\pi}} F_{1}(x) \\
& H_{33}(x)=\frac{\sigma h_{1} c_{1}}{\lambda k_{1}} \exp \left(-x^{2}\right)-x G(x)-\frac{E c_{2}}{\lambda} x W(x) \\
& H_{32}(x)=\frac{h_{0}}{\lambda \rho \sigma} \exp \left(-x^{2}\right)-1-\frac{E k_{2}}{\lambda \rho \sigma a_{2} \sqrt{\pi}} F_{1}\left(\frac{\sigma}{a_{2}}\right) \\
& H_{34}(x)=\frac{E c_{2} \sqrt{\pi}}{D c_{1}} \frac{W(x)}{F_{2}(x)} \\
& H_{35}(x)=1-H_{34}(x) \\
& H_{36}(x)=\frac{D c_{1}}{\lambda \sqrt{\pi}} F_{2}(x)-x-\frac{E c_{2}}{\lambda} W(x) \\
& H_{37}(x)=H_{36}(x)-\frac{\gamma \sqrt{\pi}}{2 D} \frac{1}{F_{2}(x)} \\
& H_{3 k}(x)=x+\frac{E c_{2}}{\lambda \sqrt{\pi}} F_{1}(x)
\end{aligned}
$$

$$
\begin{array}{rlrl}
H_{39}(x)= & \sqrt{H_{21}(x)}\left(1+\frac{\gamma \sqrt{\pi}}{2 D} \frac{1}{x F_{2}(x)}\right) & H_{40}(x)= & G(x)+\frac{E}{\lambda} \frac{1}{\sqrt{H_{24}(x)}} \sqrt{\frac{D c_{1} c_{2} k_{2}}{\sigma h_{0}} \sqrt{\pi^{3}}} \\
& \times F_{1}\left(\sqrt{\frac{\sigma h_{0} c_{2} \sqrt{\pi}}{D c_{1} k_{2}}} H_{39}(x)\right) \\
H_{41}(x)=\frac{D c_{1}}{\lambda \sqrt{\pi}} F_{2}(x) & H_{42}(x)=\frac{G(x)}{x} \\
H_{43}(x)=\frac{h_{0}}{\lambda \rho \sigma} \exp \left(-x^{2}\right)-\left(1+\frac{E c_{2}}{\lambda}\right) & H_{44}(x)=\frac{E c_{2}}{\lambda \sqrt{\pi}} W\left(\xi_{1}\right) H_{1}(x)+G\left(\xi_{1}\right) x^{2}\left(\xi_{1}>0\right) \\
H_{45}(x)=\frac{\sigma h_{0} c_{1}}{\lambda k_{1}} \exp \left(-x^{2}\right)-x G(x) & H_{46}(x)=\frac{D c_{1}}{\lambda \sqrt{\pi}} F_{2}(x)-W(x) \\
H_{47}(x)= & H_{46}(x)+\frac{\gamma \sqrt{\pi}}{2 \cdot D} \frac{1}{F_{2}(x)} & H_{48}(x)=\frac{h_{0}}{\lambda \rho \sigma} \exp \left(-x^{2}\right)-1 \\
H_{49}(x)=\frac{\exp \left(-x^{2}\right)}{x}-\frac{E}{\sigma h_{0}} \sqrt{\frac{k_{1} c_{2} k_{2}}{\pi c_{1}}} F_{1}\left(\sqrt{\frac{k_{1} c_{2}}{c_{1} k_{2}}} W(x)\right) & H_{50}(x)=\frac{D c_{1}}{\lambda \sqrt{\pi}} F_{2}(x)-\left(1+\frac{E c_{2}}{\lambda}\right) x \\
H_{51}(x)=\frac{\sigma h_{0} c_{1}}{\lambda k_{1}} \exp \left(-x^{2}\right)-x^{2}-\frac{E}{\lambda} \sqrt{\frac{c_{1} c_{2} k_{2}}{\pi k_{1}}} & H_{52}(x)=H_{51}(x)-\frac{\gamma \sqrt{\pi}}{2 D} \frac{x}{F_{2}(x)} \\
& \times x F_{1}\left(\sqrt{\frac{k_{1} c_{2}}{c_{2}}} W(x)\right) & \\
H_{53}(x)=(1-\epsilon) \sqrt{\frac{c_{1}}{k_{1} k_{2}}} x+\frac{E}{\lambda} \sqrt{\frac{c_{1} c_{2} k_{2}}{\pi k_{1}}} F_{1}(x) & H_{54}(x)=\frac{\sigma h_{0} c_{1}}{\lambda k_{1}} \exp \left(-x^{2}\right)-\left(\epsilon+(1-\epsilon) \sqrt{\frac{c_{1} k_{2}}{k_{1} c_{2}}}\right) x^{2} \\
& -\frac{E}{\lambda} \sqrt{\frac{c_{1} c_{2} k_{2}}{\pi k_{1}}} x F_{1}\left(\sqrt{\left.\frac{k_{1} c_{2}}{c_{1} k_{2}} x\right)}\right. \\
H_{55}(x)=x f(x) \exp \left(x^{2}\right)=\frac{x}{F_{2}(x)} & H_{56}(x)=\frac{1}{H_{55}(x)}-\frac{E k_{2}}{D \rho \sigma c_{1} a_{2}} F_{1}\left(\frac{\sigma}{a_{2}} H_{25}(x)\right) \\
H_{57}(x)=H_{57}(x, \gamma)=\frac{D c_{1}}{\lambda \sqrt{\pi}} F_{2}(x)-G(x, \gamma) . & H_{58}(x)=x F_{2}(x)
\end{array}
$$

The principal properties of some of these functions, for $x>0$, are

| $f\left(0^{+}\right)=0$ | $f(+\infty)=1$ | $f^{\prime}(x)>0$ |
| ---: | :--- | ---: |
| $F_{1}\left(0^{+}\right)==1$ | $F_{1}(+\infty)=+\infty$ | $F_{1}^{\prime}(x)>0$ |
| $F_{2}\left(0^{+}\right)==+\infty$ | $F_{2}(+\infty)=0$ | $F_{2}^{\prime}(x)<0$ |
| $W\left(0^{+}\right)=0$ | $W(+\infty)=+\infty$ | $W^{\prime}(x)>0$ |
| $G\left(0^{+}\right)=0$ | $G(+\infty)=+\infty$ | $G^{\prime}(x)>0$ |
| $H_{1}\left(0^{+}\right)=0$ | $H_{1}(+\infty)=+\infty$ | $H_{1}^{\prime}(x)>0$ |
| $H_{2}\left(0^{+}\right)=\frac{h_{0}}{\lambda \rho a_{1}}$ | $H_{2}(+\infty)=-\infty$ | $H_{2}^{\prime}(x)<0$ |
| $H_{3}\left(0^{+}\right)=\alpha_{20}$ | $H_{3}(+\infty)=-\infty$ | $H_{3}^{\prime}(x)<0$ |
| $H_{4}\left(0^{+}\right)=\frac{h_{0}}{\lambda \rho a_{1}} \alpha_{20}$ | $H_{4}(+\infty)=-\infty$ | $H_{4}^{\prime}(x)<0$ |
| $H_{5}\left(0^{+}\right)=-\frac{h_{0}}{\lambda \rho a_{1}} \alpha_{21}$ | $H_{5}(+\infty)=+\infty$ | $H_{5}^{\prime}(x)>0$ |
| $H_{6}\left(0^{+}\right)=\frac{E k_{2}}{\lambda \rho a_{1} a_{2} \sqrt{\pi}}$ | $H_{6}(+\infty)=+\infty$ | $H_{6}^{\prime}(x)>0$ |
| $H_{7}\left(0^{+}\right)=\frac{h_{0}}{\lambda \rho a_{1}} \alpha_{20}$ | $H_{7}(+\infty)=-\infty$ | $H_{7}^{\prime}(x)<0$ |
| $H_{8}\left(0^{+}\right)=\alpha_{3}$ | $H_{8}(+\infty)=+\infty$ | $H_{8}^{\prime}(x)>0$ |
| $H_{y}\left(0^{+}\right)=\alpha_{4}$ | $H_{9}(+\infty)=+\infty$ | $H_{9}^{\prime}(x)>0$ |
| $H_{10}\left(0^{+}\right)=\alpha_{5}$ | $H_{10}(+\infty)=-\infty$ | $H_{10}^{\prime}(x)<0$ |
| $H_{11}\left(0^{+}\right)=1$ | $H_{11}(+\infty)=+\infty$ | $H_{11}^{\prime}(x)>0$ |
| $H_{12}\left(0^{+}\right)=+\infty$ | $H_{12}^{\prime}(x)<0$ |  |
| $H_{13}\left(0^{+}\right)=0$ | $H_{13}^{\prime}(x)>0$ |  |
| $H_{14}\left(0^{+}\right)=0$ | $H_{13}(+x)=+\infty$ | $H_{14}^{\prime}(x)>0$ |


| $H_{15}\left(0^{+}\right)=\alpha_{6}$ |  |
| :---: | :---: |
| $H_{16}\left(0^{+}\right)=0$ |  |
|  | $H_{17}\left(0^{+}\right)=\alpha_{7}$ |
|  | $H_{18}$ |
|  | $H_{19}\left(0^{+}\right)=0$ |
|  | $H_{20}\left(0^{+}\right)=0$ |
|  | $H_{23}\left(0^{+}\right)=\alpha_{9}$ |
|  | $H_{24}\left(0^{+}\right)=$ |
|  | $H_{25}$ |
|  | $H_{27}\left(0^{+}\right)=\alpha_{10}$ |
|  | $\left.0^{+}\right)=\alpha_{11}$ |
|  | $H_{29}\left(0^{+}\right)=\alpha_{12}$ |
|  | $H_{30}\left(0^{+}\right)=\alpha_{13}$ |
|  | $H_{31}\left(0^{+}\right)=\alpha_{14}$ |
|  | $H_{32}\left(0^{+}\right)=$ |
|  | $H_{33}\left(0^{+}\right)=\frac{\sigma h_{0} c_{1}}{\lambda k_{1}}$ |
|  | $H_{34}\left(0^{+}\right)=0$ |
|  | $H_{35}\left(0^{+}\right)=1$ |
|  | $H_{36}\left(0^{+}\right)=+\infty$ |
|  | $H_{37}\left(0^{+}\right)=+\infty$ |
|  | $H_{38}\left(0^{+}\right)=\frac{E c_{2}}{\lambda \sqrt{\pi}}$ |
|  | $H_{39}\left(0^{+}\right)=0$ |
|  | $H_{40}\left(0^{+}\right)=\alpha_{16}$ |
|  | $H_{41}\left(0^{+}\right)=+\infty$ |
|  | $H_{42}\left(0^{+}\right)=\alpha_{17}$ |
| $H_{43}\left(0^{+}\right)=\alpha_{18}$ |  |
| $H_{44}\left(0^{+}\right)=0$ |  |
|  | $H_{45}\left(0^{+}\right)=\frac{\sigma h_{0} c_{1}}{\lambda k_{1}}$ |
| $H_{46}\left(0^{+}\right)=+\infty$ |  |
| $H_{47}\left(0^{+}\right)=+\infty$ |  |
| $H_{48}\left(0^{+}\right)=\alpha_{19}$ |  |
| $H_{49}\left(0^{+}\right)=+\infty$ |  |
| $H_{50}\left(0^{+}\right)=+\infty$ |  |
| $H_{51}\left(0^{+}\right)=\frac{\sigma h_{0} c_{1}}{\lambda k_{1}}$ |  |
| $H_{52}\left(0^{+}\right)=\frac{\sigma h_{0} c_{1}}{\lambda k_{1}}$ |  |
| $H_{53}\left(0^{+}\right)=\alpha_{8}$ |  |
| $H_{54}\left(0^{+}\right)=\frac{\sigma h_{0} c_{1}}{\lambda k_{1}}$ |  |
| $H_{5 s}\left(0^{+}\right)=0$ |  |
| $H_{56}\left(0^{+}\right)=+\infty$ |  |
| $H_{57}\left(0^{+}\right)=+\infty$ |  |
|  | $H_{58}\left(0^{+}\right)=\frac{\sqrt{\pi}}{2}$ |

with
$\alpha_{3}=\frac{2}{\sqrt{\pi}}+\frac{\gamma \sqrt{\pi}}{2 D}$

| $H_{15}(+\infty)=-\infty$ | $H_{15}^{\prime}(x)<0$ |
| :---: | :---: |
| $H_{16}(+\infty)=1$ | $H_{16}^{\prime}(x)>0$ |
| $H_{17}(+\infty)=-\infty$ | $H_{17}^{\prime}(x)<0$ |
| $H_{18}(+\infty)=+\infty$ | $H_{i 8}^{\prime}(x)>0$ |
| $H_{19}(+\infty)=+\infty$ | $H_{19}^{\prime}(x)>0$ |
| $H_{20}(+\infty)=+\infty$ | $H_{20}^{\prime}(x)>0$ |
| $H_{23}(+\infty)=-\infty$ | $H_{23}^{\prime}(x)<0$ |
| $H_{24}(+\infty)=0$ | $H_{24}^{\prime}(x)<0$ |
| $\mathrm{H}_{25}(+\infty)=+\infty$ | $H_{25}(x)>0$ |
| $H_{27}(+\infty)=-\infty$ | $H_{27}^{\prime}(x)<0$ |
| $H_{28}(+\infty)=-\infty$ | $H_{28}^{\prime}(x)<0$ |
| $H_{29}(+\infty)=-\infty$ | $H_{29}^{\prime}(x)<0$ |
| $H_{30}(+\infty)=+\infty$ | $H_{30}^{\prime}(x)>0$ |
| $H_{31}(+\infty)=+\infty$ | $H_{31}^{\prime}(x)>0$ |
| $H_{32}(+\infty)=-1-\alpha_{14}$ | $H_{32}^{\prime}(x)<0$ |
| $H_{33}(+\infty)=-\infty$ | $H_{33}^{\prime}(x)<0$ |
| $H_{34}(+\infty)=+\infty$ | $H_{34}^{\prime}(x)>0$ |
| $H_{3,5}(+\infty)=-\infty$ | $H_{35}^{\prime}(x)<0$ |
| $H_{36}(+\infty)=-\infty$ | $H_{36}^{\prime}(x)<0$ |
| $H_{37}(+\infty)=-\infty$ | $H_{37}^{\prime}(x)<0$ |
| $H_{38}(+\infty)=+\infty$ | $H_{38}^{\prime}(x)>0$ |
| $H_{39}(+\infty)=+\infty$ | $H_{39}^{\prime}(x)>0$ |
| $H_{40}(+\infty)=+\infty$ | $H_{40}^{\prime}(x)>0$ |
| $H_{41}(+\infty)=0$ | $H_{41}^{\prime}(x)<0$ |
| $H_{42}(+\infty)=+\infty$ | $H_{42}^{\prime}(x)>0$ |
| $H_{43}(+\infty)=-1-\frac{E c_{2}}{\lambda}$ | $H_{43}^{\prime}(x)<0$ |
| $H_{44}(+\infty)=+\infty$ | $H_{44}^{\prime}(x)>0$ |
| $H_{45}(+\infty)=-\infty$ | $H_{45}^{\prime}(x)<0$ |
| $H_{46}(+\infty)=-\infty$ | $H_{46}^{\prime}(x)<0$ |
| $H_{47}(+\infty)=-\infty$ | $H_{47}^{\prime}(x)<0$ |
| $H_{48}(+\infty)=-1$ | $H_{4 x}^{\prime}(x)<0$ |
| $H_{49}(+\infty)=-\infty$ | $H_{49}^{\prime}(x)<0$ |
| $H_{50}(+\infty)=-\infty$ | $H_{501}^{\prime}(x)<0$ |
| $H_{51}(+\infty)=-\infty$ | $H_{51}^{\prime}(x)<0$ |
| $H_{52}(+\infty)=-\infty$ | $H_{52}^{\prime}(x)<0$ |
| $H_{53}(+\infty)=+\infty$ | $H_{53}^{\prime}(x)>0$ |
| $H_{54}(+\infty)=-\infty$ | $H_{54}^{\prime}(x)<0$ |
| $H_{5 s}(+\infty)=+\infty$ | $H_{S S}^{\prime}(x)>0$ |
| $H_{56}(+\infty)=-\infty$ | $H_{56}^{\prime}(x)<0$ |
| $H_{57}(+\infty)=-\infty$ | $H_{57}^{\prime}(x)<0$ |
| $H_{58}(+\infty)=0$ | $H_{S K}^{\prime}(x)<0$ |

$\alpha_{4}=\frac{2}{\sqrt{\pi}}+\frac{(1-\epsilon) \gamma \sqrt{\pi}}{2 D}$

$$
\begin{array}{ll}
\alpha_{6}=\frac{D c_{1}}{\lambda \sqrt{\pi}}\left(1-\frac{E k_{2}}{h_{\mathrm{o}} a_{2} \sqrt{\pi}}\right) & \alpha_{5}=\frac{E h_{\mathrm{o}} k_{2}}{D \lambda \rho k_{1} a_{2}}\left(\frac{h_{\mathrm{o}} a_{2} \sqrt{\pi}}{E k_{2}}-F_{1}\left(\frac{D k_{1}}{h_{\mathrm{o}} a_{2} \sqrt{\pi}} \alpha_{3}\right)\right) \\
\alpha_{7}=\frac{h_{\mathrm{o}}}{E \rho a_{1} c_{2}} & \alpha_{8}=\frac{E}{\lambda} \sqrt{\frac{c_{1} c_{2} k_{2}}{\pi k_{1}}} \\
\alpha_{9}=\frac{h_{\mathrm{o}}}{E \rho \sigma c_{2}}-\left(1+\frac{\gamma}{D}+\frac{\lambda}{E c_{2}}\left(1+\frac{(1-\epsilon) \gamma}{D}\right)\right) & \alpha_{10}=1-\frac{E k_{2}}{h_{\mathrm{o}} a_{2} \sqrt{\pi}} F_{1}\left(\frac{\sigma}{a_{2}}\left(1+\frac{\gamma}{D}\right)\right) \\
\alpha_{11}=\frac{h_{\mathrm{o}}}{\lambda \rho \sigma}-1-\frac{(1-\epsilon) \gamma}{D} & \alpha_{12}=\frac{h_{\mathrm{o}}}{\lambda \rho \sigma}-1-\frac{E k_{2}}{\lambda \rho \sigma a_{2} \sqrt{\pi}} F_{1}\left(\frac{\sigma}{a_{2}}\left(1+\frac{\gamma}{D}\right)\right) \\
\alpha_{14}=\frac{E k_{2}}{\lambda \rho \sigma a_{2} \sqrt{\pi}} & \alpha_{13}=\frac{\gamma}{D}-\frac{h_{0}}{\lambda \rho \sigma}+1+\frac{E k_{2}}{\lambda \rho \sigma a_{2} \sqrt{\pi}} F_{1}\left(\frac{\sigma}{a_{2}}\left(1+\frac{\gamma}{D}\right)\right) \\
\alpha_{16}=\frac{E}{\lambda} \sqrt{\frac{D c_{1} c_{2} k_{2}}{2 \sigma h_{\mathrm{o}} \sqrt{\pi}}} & \alpha_{15}=\frac{h_{\mathrm{o}}}{\lambda \rho \sigma}-1-\frac{E k_{2}}{\lambda \rho \sigma a_{2} \sqrt{\pi}} F_{1}\left(\frac{\sigma}{a_{2}}\right) \\
\alpha_{19}=1+\frac{(1-\epsilon) \gamma}{D} & \alpha_{18}=\frac{h_{\mathrm{o}}}{\lambda \rho \sigma}-1-\frac{E c_{2}}{\lambda} \\
\alpha_{19}=\frac{h_{\mathrm{o}}}{\lambda \rho \sigma}-1 & \alpha_{20}=1-\frac{E k_{2}}{h_{\mathrm{o}} a_{2} \sqrt{\pi}} .
\end{array}
$$

## APPENDIX B

The restrictions used in the text are the following:
(R1) $h_{\mathrm{o}}>\frac{E k_{2}}{a_{2} \sqrt{\pi}}$
(R2) $\frac{D k_{1}}{h_{\mathrm{o}} a_{1} \sqrt{\pi}}<f\left(x_{2}\right)$, where $x_{2}$ is the unique positive zero of $H_{2}$
(R3) $\frac{D k_{1}}{h_{0} a_{1} \sqrt{\pi}}<f\left(x_{3}\right)$, where $x_{3}$ is the unique positive zero of $H_{3}$
(R4) $\frac{D k_{1}}{h_{0} a_{1} \sqrt{\pi}}<f\left(x_{4}\right)$, where $x_{4}$ is the unique positive zero of $H_{4}$
(R5) $\frac{D k_{1}}{h_{\mathrm{o}} a_{\mathrm{j}} \sqrt{\pi}}>f\left(x_{5}\right)$, where $x_{5}$ is the unique positive zero of $H_{5}$
(R6) $\frac{D k_{1}}{h_{\mathrm{o}} a_{1} \sqrt{\pi}}<f\left(x_{7}\right)$, where $x_{7}$ is the unique positive zero of $H_{7}$
(R7) $h_{0}>\frac{D k_{1}}{a_{2} \sqrt{\pi}}\left(\frac{2}{\sqrt{\pi}}+\frac{\gamma \sqrt{\pi}}{2 D}\right) \frac{1}{\eta}$, where $\eta$ is the unique positive solution of the equation $H_{11}(x)=H_{12}(x), x>0$
(R8) $\frac{D k_{1}}{h_{\mathrm{o}} a_{1} \sqrt{\pi}}<f\left(x_{17}\right)$, where $x_{17}$ is the unique positive zero of $H_{17}$
(R9) $\frac{h_{\mathrm{o}}}{E \rho c_{2}}>1+\frac{\gamma}{D}+\frac{\lambda}{E c_{2}}\left(1+\frac{(1-\epsilon) \gamma}{D}\right)$
(R10) $h_{\circ}>\frac{D k_{1}}{2 \sigma}$
(R11) $\frac{D \rho \sigma c_{1}}{h_{0} \sqrt{\pi}}<H_{20}\left(x_{23}\right)$, where $x_{23}$ is the unique positive zero of $H_{23}$
(R12) $\frac{D k_{1}}{\sigma h_{0} \sqrt{\pi}}>H_{24}\left(x_{23}\right)$, where $x_{23}$ is the unique positive zero of $H_{23}$
(R13) $h_{0}>\frac{E k_{2}}{a_{2} \sqrt{\pi}} F_{1}\left(\frac{\sigma}{a_{2}}\left(1+\frac{\gamma}{D}\right)\right)$
(R14) $\frac{h_{o}}{\lambda \rho \sigma}>1+\frac{(1-\epsilon) \gamma}{D}$
(R15) $\frac{h_{0}}{\lambda \rho \sigma}>1+\frac{E k_{2}}{\lambda \rho \sigma a_{2} \sqrt{\pi}} F_{1}\left(\frac{\sigma}{a_{2}}\left(1+\frac{\gamma}{D}\right)\right)$
(R16) $\frac{h_{o}}{\lambda \rho \sigma}>1+\frac{\gamma}{D}+\frac{E k_{2}}{\lambda \rho \sigma a_{2} \sqrt{\pi}} F_{1}\left(\frac{\sigma}{a_{2}}\left(1+\frac{\gamma}{D}\right)\right)$
(R17) $\frac{h_{\mathrm{o}}}{\lambda \rho \sigma}>1+\frac{E k_{2}}{\lambda \rho \sigma a_{2} \sqrt{\pi}} F_{1}\left(\frac{\sigma}{a_{2}}\right)$
(R18) $\frac{D \rho \sigma c_{1}}{h_{\mathrm{o}} \sqrt{\pi}}<H_{20}\left(x_{27}\right)$, where $x_{27}$ is the unique positive zero of $H_{27}$
(R19) $\frac{D \rho \sigma c_{1}}{h_{0} \sqrt{\pi}}<H_{20}\left(x_{28}\right)$, where $x_{28}$ is the unique positive zero of $H_{28}$
(R20) $H_{20}\left(x_{29}\right)>\frac{D \rho \sigma c_{1}}{h_{0} \sqrt{\pi}}>H_{20}\left(x_{30}\right)$, where $x_{29}$ is the unique positive zero of $H_{29}, x_{30}$ is the unique positive zero of $H_{30}$
(R21) $\frac{D \rho \sigma c_{1}}{h_{\mathrm{o}} \sqrt{\pi}}<H_{20}\left(x_{32}\right)$, where $x_{32}$ is the unique positive zero of $H_{32}$
(R22) $\frac{D k_{1}}{\sigma h_{\mathrm{o}} \sqrt{\pi}}>H_{24}\left(x_{33}\right)$, where $x_{33}$ is the unique positive zero of $H_{33}$
(R23) $h_{\mathrm{o}}=\frac{D k_{1}}{a_{1} f\left(\frac{\sigma}{a_{1}}\right) \sqrt{\pi}}$
(R24) $H_{35}\left(\frac{\sigma}{a_{1}}\right)>0$ or $\frac{\sigma}{a_{1}}<x_{35}$, where $x_{35}$ is the unique positive zero of $H_{35}$
(R25) $H_{36}\left(\frac{\sigma}{a_{1}}\right)>0$ or $\frac{\sigma}{a_{1}}<x_{36}$, where $x_{36}$ is the unique positive zero of $H_{36}$
(R26) $\frac{D k_{1}}{E \rho a_{1} a_{2} c_{2}} F_{2}\left(\frac{\sigma}{a_{1}}\right) \leq 1$
(R27) $\frac{\lambda \rho \sigma a_{2} \sqrt{\pi}}{E k_{2}}\left(\frac{D k_{1}}{\lambda \rho \sigma a_{1} \sqrt{\pi}} F_{2}\left(\frac{\sigma}{a_{1}}\right)-1\right)>1$
(R28) $F_{1}\left(\frac{\sigma}{a_{2}}\right)<\frac{\lambda \rho \sigma a_{2} \sqrt{\pi}}{E k_{2}}\left(\frac{D k_{1}}{\lambda \rho \sigma a_{1} \sqrt{\pi}} F_{2}\left(\frac{\sigma}{a_{1}}\right)-1\right)$
(R29) $H_{57}\left(\frac{\sigma}{a_{1}}\right)>0$ or $\frac{\sigma}{a_{1}}<x_{57}$, where $x_{57}$ is the unique positive zero of $H_{57}$
(R30) $\frac{h_{0}}{\lambda \rho \sigma}>1+\frac{(1-\epsilon) \gamma}{D}+\frac{E k_{2}}{\lambda \rho \sigma a_{2} \sqrt{\pi}} F_{1}\left(\frac{\sigma}{a_{2}}\left(1+\frac{\gamma}{D}\right)\right)$
(R31) $\frac{D k_{1}}{\sigma h_{\mathrm{o}} \sqrt{\pi}}>H_{24}\left(x_{27}\right)$, where $x_{27}$ is the unique positive zero of $H_{27}$
(R32) $h_{\mathrm{o}}>\rho \sigma\left(\lambda+E c_{2}\right)$
(R33) $H_{43}\left(\frac{\sigma}{a_{1}}\right)>0$ or $\frac{\sigma}{a_{1}}<x_{43}$, where $x_{43}$ is the unique positive zero of $H_{43}$ when (R32) is verified
(R34) $\frac{D k_{1}}{\sigma h_{0} \sqrt{\pi}}>H_{24}\left(x_{28}\right)$, where $x_{28}$ is the unique positive zero of $H_{28}$
(R35) $\frac{D k_{1}}{\sigma h_{\mathrm{o}} \sqrt{\pi}}>H_{24}\left(x_{29}\right)$, where $x_{29}$ is the unique positive zero of $H_{29}$
(R36) $\frac{D k_{1}}{\sigma h_{\mathrm{o}} \sqrt{\pi}}<H_{24}\left(x_{30}\right)$, where $x_{30}$ is the unique positive zero of $H_{31}$
(R37) $\frac{D k_{1}}{\sigma h_{\mathrm{o}} \sqrt{\pi}}>H_{24}\left(x_{32}\right)$, where $x_{32}$ is the unique positive zero of $H_{32}$
(R38) $\frac{D k_{1}}{\sigma h_{\mathrm{o}} \sqrt{\pi}}>H_{24}\left(x_{45}\right)$, where $x_{45}$ is the unique positive zero of $H_{45}$
(R39) $H_{47}\left(\frac{\sigma}{a_{1}}\right)>0$ or $\frac{\sigma}{a_{1}}<x_{47}$, where $x_{47}$ is the unique positive zero of $H_{47}$
(R40) $H_{55}\left(\frac{\sigma}{a_{1}}\right)<\frac{D c_{1}}{\lambda \sqrt{\pi}}$
(R41) $\frac{D k_{1}}{\sigma h_{0} \sqrt{\pi}}>H_{24}\left(x_{54}\right)$, where $x_{54}$ is the unique positive zero of $H_{54}$
(R42) $\frac{D k_{1}}{\sigma h_{6} \sqrt{\pi}}>H_{24}\left(x_{49}\right)$, where $x_{49}$ is the unique positive zero of $H_{49}$
(R43) $H_{50}\left(\frac{\sigma}{a_{1}}\right)>0$ or $\frac{\sigma}{a_{1}}<x_{50}$, where $x_{50}$ is the unique positive zero of $H_{50}$
(R44) $\frac{D k_{1} a_{2}}{E a_{1} k_{2}} F_{2}\left(\frac{\sigma}{a_{1}}\right)>1$
(R45) $F_{1}\left(\frac{\sigma}{a_{2}}\right)<\frac{D k_{1} a_{2}}{E a_{1} k_{2}} F_{2}\left(\frac{\sigma}{a_{1}}\right)$
(R46) $\frac{D k_{1}}{\sigma h_{0} \sqrt{\pi}}>H_{24}\left(x_{51}\right)$, where $x_{51}$ is the unique positive zero of $H_{51}$
(R47) $\frac{D k_{1}}{\sigma h_{\mathrm{o}} \sqrt{\pi}}<H_{24}\left(x_{52}\right)$, where $x_{52}$ is the unique positive zero of $H_{52}$
(R48) $H_{56}\left(\frac{\sigma}{a_{1}}\right)>0$ or $\frac{\sigma}{a_{1}}<x_{56}$, where $x_{56}$ is the unique positive zero of $H_{56}$.

## APPENDIX C

Table 1

| Case | Unknown coefficients | Restrictions | Solution |
| :---: | :---: | :---: | :---: |
| 1 | $c_{2}, \sigma, \omega$ | (R2) | $\sigma=a_{1} \xi_{1}, \quad c_{2}=\frac{c_{1} k_{2}}{k_{1}} \frac{B^{2}}{W^{2}\left(\xi_{1}\right)}, \quad \omega=a_{1} w\left(\xi_{1}\right)$ <br> where $\xi_{1}$ is the unique positive solution of the equation $f(x)=\frac{D k_{1}}{h_{0} a_{1} \sqrt{\pi}}, \quad x>0$ <br> and $B$ is the only positive solution of the equation $H_{1}(x)=\frac{\lambda k_{1} \sqrt{\pi}}{E c_{1} k_{2}} W\left(\xi_{1}\right) H_{2}\left(\xi_{1}\right), \quad x>0 .$ |
| 2 | $\lambda, \sigma, \omega$ | (R1) <br> (R3) | $\sigma=a_{1} \xi_{1}, \quad \lambda=\frac{h_{\mathrm{o}}}{\rho a_{1}} \frac{H_{3}\left(\xi_{1}\right)}{G\left(\xi_{1}\right)}, \quad \omega=a_{1} W\left(\xi_{1}\right)$ |

where $\xi_{1}$ is given as in case 1 .

| 3 | $\epsilon, \sigma, \omega$ | $\begin{aligned} & \text { (R1) } \\ & \text { (R4) } \\ & \text { (R5) } \end{aligned}$ | $\sigma=a_{1} \xi_{1}, \quad \epsilon=\frac{2 D}{\gamma \sqrt{\pi}} F_{2}\left(\xi_{1}\right) H_{5}\left(\xi_{1}\right), \quad \omega=a_{1} W\left(\xi_{1}\right)$ <br> where $\xi_{1}$ is given as in case 1 . |
| :---: | :---: | :---: | :---: |
| 4 | $\gamma, \sigma, \omega$ | $\begin{aligned} & \text { (R1) } \\ & \text { (R6) } \end{aligned}$ | $\sigma=a_{1} \xi_{1}, \quad \gamma=\frac{2 D}{\sqrt{\pi}}\left(\frac{a_{2}}{a_{1}} B-\xi_{1}\right) F_{2}\left(\xi_{1}\right), \quad \omega=a_{1} W\left(\xi_{1}\right)$ <br> where $\xi_{1}$ is given as in case 1 and $B$ is the only positive solution of the equation $H_{6}(x)=\frac{h_{0}}{\lambda \rho a_{1}} \exp \left(-\xi_{1}^{2}\right)-\epsilon \xi_{1}, \quad x>\frac{a_{1}}{a_{2}} \xi_{1} .$ |
| 5 | $c_{1}, \sigma, \omega$ | (R7) | $\sigma=\frac{D k_{1}}{h_{\mathrm{o}} \sqrt{\pi}} \frac{1}{H_{24}\left(\xi_{1}\right)}, \quad c_{1}=\frac{\pi h_{\mathrm{o}}^{2}}{D^{2} \rho k_{1}} f^{2}\left(\xi_{1}\right), \quad \omega=\frac{D k_{1}}{h_{\mathrm{o}} \sqrt{\pi}} H_{8}\left(\xi_{1}\right)$ <br> where $\xi_{1}$ is the unique positive solution of the equation $H_{9}(x)=H_{10}(x), \quad x>0 .$ |
| 6 | $k_{1}, \sigma, \omega$ | (R1) | $\sigma=\frac{h_{0} \sqrt{\pi}}{D \rho c_{1}} H_{20}\left(\xi_{1}\right), \quad k_{1}=\frac{\pi h_{o}^{2}}{D^{2} \rho c_{1}} f^{2}\left(\xi_{1}\right), \quad \omega=\frac{h_{\mathrm{o}} \sqrt{\pi}}{D \rho c_{1}} H_{13}\left(\xi_{1}\right)$ <br> where $\xi_{1}$ is the unique positive solution of the equation $H_{14}(x)=H_{15}(x), \quad x>0 .$ |
| 7 | $k_{2}, \sigma, \omega$ | (R8) | $\sigma=a_{1} \xi_{1}, \quad k_{2}=\frac{k_{1} c_{2}}{c_{1}} \frac{W^{2}\left(\xi_{1}\right)}{B^{2}}, \quad \omega=a_{1} W\left(\xi_{1}\right)$ <br> where $\xi_{1}$ is given as in case 1 and $B$ is the only positive solution of the equation $\frac{1}{H_{16}(x)}=\frac{\lambda}{E c_{2}} \frac{H_{2}\left(\xi_{1}\right)}{W\left(\xi_{1}\right)}, \quad x>0 .$ |
| 8 | $\rho, \sigma, \omega$ | - | $\sigma=\frac{\lambda k_{1}}{h_{0} c_{1}} \xi_{1} H_{18}\left(\xi_{1}\right), \quad \rho=\frac{h_{0}^{2} c_{1}}{\lambda^{2} k_{1}} \frac{1}{H_{18}^{2}\left(\xi_{1}\right)}, \quad \omega=\frac{\lambda k_{1}}{h_{0} c_{1}} w\left(\xi_{1}\right) H_{18}\left(\xi_{1}\right)$ <br> where $\xi_{1}$ is the unique positive solution of the equation $H_{19}(x)=\frac{D c_{1}}{\lambda \sqrt{\pi}}, \quad x>0 .$ |

## APPENDIX D

Table 2
$\left.\begin{array}{ccccc}\hline \text { Case } & \begin{array}{c}\text { Unknown } \\ \text { coefficients }\end{array} & \text { Restrictions } & \text { Solution } \\ \hline 9 & \omega, k_{1}, k_{2} & (\mathrm{R} 9) & \omega=\sigma H_{25}\left(\xi_{1}\right), & k_{1}=\rho \sigma^{2} c_{1} \frac{1}{\xi_{1}^{2}},\end{array} k_{2}=\rho \sigma^{2} c_{2} \frac{H_{25}^{2}\left(\xi_{1}\right)}{B^{2}}\right)$
where $\xi_{1}$ is the unqiue solution of the equation

$$
H_{20}(x)=\frac{D \rho \sigma c_{1}}{h_{0} \sqrt{\pi}}, \quad x>0
$$

and $B$ is the only solution of the equation

$$
\frac{1}{H_{16}(x)}=\frac{\lambda}{E c_{2}} \frac{H_{21}\left(\xi_{1}\right)}{W\left(\xi_{1}\right)}, \quad x>0
$$

$10 \quad \omega, c_{1}, k_{2}$| $\left(\begin{array}{l}\text { (R9) } \\ \text { (R10) }\end{array}\right.$ | $\omega=\sigma H_{25}\left(\xi_{1}\right), \quad c_{1}=\frac{k_{1}}{\rho \sigma^{2}} \xi_{1}^{2}, \quad k_{2}=\rho \sigma^{2} c_{2} \frac{H_{25}^{2}\left(\xi_{1}\right)}{B^{2}}$ |
| ---: | ---: | ---: |

where $\xi_{1}$ is the unique solution of the equation

$$
H_{24}(x)=\frac{D k_{1}}{\sigma h_{\mathrm{o}} \sqrt{\pi}}, \quad x>0
$$

and $B$ is given as in case 9 .

| 11 | $\omega, \lambda, k_{1}$ | $\begin{aligned} & \text { (R13) } \\ & \text { (R18) } \end{aligned}$ | $\omega=\sigma H_{25}\left(\xi_{1}\right), \quad \lambda=\frac{h_{0}}{\rho \sigma} \frac{H_{20}\left(\xi_{1}\right)}{G\left(\xi_{1}\right)}, \quad k_{1}=\rho \sigma^{2} c_{1} \frac{1}{\xi_{1}^{2}}$ <br> where $\xi_{1}$ is given as in case 9 . |
| :---: | :---: | :---: | :---: |
| 12 | $\omega, k_{1}, c_{2}$ | $\begin{aligned} & \text { (R14) } \\ & \text { (R19) } \end{aligned}$ | $\omega=\sigma H_{25}\left(\xi_{1}\right), \quad k_{1}=\rho \sigma^{2} c_{1} \frac{1}{\xi_{1}^{2}}, \quad c_{2}=\frac{k_{2}}{\rho \sigma^{2}} \frac{B^{2}}{H_{25}^{2}\left(\xi_{1}\right)}$ <br> where $\xi_{1}$ is given as in case 9 and $B$ is the only solution of the equation $H_{1}(x)=\frac{\lambda \rho \sigma^{2} \sqrt{\pi}}{E k_{2}} H_{25}\left(\xi_{1}\right) H_{28}\left(\xi_{1}\right), \quad x>0 .$ |
| 13 | $\omega, \epsilon, k_{1}$ | $\begin{aligned} & \text { (R16) } \\ & \text { (R20) } \end{aligned}$ | $\omega=\sigma H_{25}\left(\xi_{1}\right), \quad k_{1}=\rho \sigma^{2} c_{1} \frac{1}{\xi_{1}^{2}}, \quad \epsilon=1-\frac{2 D}{\gamma \sqrt{\pi}} H_{29}\left(\xi_{1}\right) H_{58}\left(\xi_{1}\right)$ |

where $\xi_{1}$ is given as in case 9 .

| 14 | $\omega, \gamma, k_{1}$ | $\begin{aligned} & \text { (R17) } \\ & \text { (R21) } \end{aligned}$ | $\omega=B a_{2}, \quad \gamma=\frac{2 D}{\sqrt{\pi}}\left(\frac{a_{2}}{\sigma} B-1\right) H_{58}\left(\xi_{1}\right), \quad k_{1}=\rho \sigma^{2} c_{1} \frac{1}{\xi_{1}^{2}}$ <br> where $\xi_{1}$ is given as in case 9 and $B$ is the only solution of the equation $H_{31}(x)=\frac{h_{\mathrm{o}}}{\lambda \rho \sigma} \exp \left(-\xi_{1}^{2}\right)-\epsilon, \quad x>\frac{\sigma}{a_{2}} .$ |
| :---: | :---: | :---: | :---: |
| 15 | $\omega, \rho, k_{2}$ | $\begin{aligned} & \text { (R10) } \\ & (\mathrm{R} 22) \end{aligned}$ | $\omega=\sigma H_{25}\left(\xi_{1}\right), \quad \rho=\frac{k_{1}}{\sigma^{2} c_{1}} \xi_{1}^{2}, \quad k_{2}=\frac{k_{1} c_{2}}{c_{1}} \frac{W^{2}\left(\xi_{1}\right)}{B^{2}}$ <br> where $\xi_{1}$ is given as in case 10 and $B$ is the only solution of the equation $\frac{1}{H_{16}(x)}=\frac{\lambda}{E c_{2}} \frac{H_{45}\left(\xi_{1}\right)}{\xi_{1} W\left(\xi_{1}\right)}, \quad x>0 .$ |
| 16 | $\omega, \lambda, k_{2}$ | $\begin{aligned} & \text { (R23) } \\ & \text { (R24) } \end{aligned}$ | $\begin{aligned} \omega=a_{1} W\left(\frac{\sigma}{a_{1}}\right), \quad k_{2}= & \frac{k_{1} c_{2}}{c_{1}} \frac{W^{2}\left(\frac{\sigma}{a_{1}}\right)}{B^{2}}, \quad \lambda=\frac{1}{G\left(\frac{\sigma}{a_{1}}\right)}\left(\frac{D c_{1}}{\sqrt{\pi}} F_{2}\left(\frac{\sigma}{a_{1}}\right)-E c_{2} \frac{W\left(\frac{\sigma}{a_{1}}\right)}{H_{16}(B)}\right) \\ & \text { for any } B>H_{16}^{-1}\left(\frac{E c_{2} \sqrt{\pi}}{D c_{1}} \frac{W\left(\frac{\sigma}{a_{1}}\right)}{F_{2}\left(\frac{\sigma}{a_{1}}\right)}\right) . \end{aligned}$ |

Table 2. contd.

| Case | Unknown coefficients | Restrictions | Solution |
| :---: | :---: | :---: | :---: |
| 17 | $\omega, \epsilon, k_{2}$ | $\begin{aligned} & \text { (R23) } \\ & \text { (R25) } \end{aligned}$ | $\omega=a_{1} W\left(\frac{\sigma}{a_{1}}\right), \quad k_{2}=\frac{k_{1} c_{2}}{c_{1}} \frac{W^{2}\left(\frac{\sigma}{a_{1}}\right)}{B^{2}},$ |
|  |  |  | $\epsilon=\frac{2 D}{\gamma \sqrt{\pi}} W\left(\frac{\sigma}{a_{1}}\right) F_{2}\left(\frac{\sigma}{a_{1}}\right)-\frac{2 D}{\lambda \gamma^{\sqrt{\pi}}} F_{2}\left(\frac{\sigma}{a_{1}}\right)\left(\frac{D c_{1}}{\sqrt{\pi}} F_{2}\left(\frac{\sigma}{a_{1}}\right)-E c_{2} \frac{W\left(\frac{\sigma}{a_{1}}\right)}{H_{16}(B)}\right)$ <br> for any $H_{16}^{-1}\left(\frac{1}{A}\right)<B<H_{16}^{-1}\left(\frac{1}{C}\right)$ if $\frac{\sigma}{a_{1}}<x_{37}$ |
|  |  |  | or for any $B>H_{16}^{-1}\left(\frac{1}{A}\right)$ if $x_{37} \leq \frac{\sigma}{a_{1}}<x_{36}$ |
|  |  |  | where $A=\frac{\lambda}{E c_{2} W\left(\frac{\sigma}{a_{1}}\right)}\left(H_{41}\left(\frac{\sigma}{a_{1}}\right)-\frac{\sigma}{a_{1}}\right)$ |
|  |  |  | $C=\frac{\lambda}{E c_{2} W\left(\frac{\sigma}{a_{1}}\right)} H_{46}\left(\frac{\sigma}{a_{1}}\right) .$ |


| $18, ~$(R23) <br> (R26) <br> (R27) <br> (R28) | $\omega=\omega(\gamma)=\sigma\left(1+\frac{\gamma \sqrt{\pi}}{2 D} \frac{1}{H_{58}\left(\frac{\sigma}{a_{1}}\right)}\right)$, |
| :--- | :--- | :--- |
|  | $\epsilon=\epsilon(\gamma)=1-\frac{\sigma}{\omega(\gamma)-\sigma}\left(\frac{D k_{1}}{\lambda \rho \sigma \sqrt{\pi}} H_{58}\left(\frac{\sigma}{a_{1}}\right)-1\right)-\frac{1}{\omega(\gamma)-\sigma} \frac{E k_{2}}{\lambda \rho a_{2} \sqrt{\pi}} F_{1}\left(\frac{\omega(\gamma)}{a_{2}}\right)$ |
|  | for any $0<\gamma<\frac{2 D a_{2}}{a_{1} \sqrt{\pi}} F_{2}\left(\frac{\sigma}{a_{1}}\right)\left(F_{1}^{-1}(B)-\frac{\sigma}{a_{2}}\right)$ |

where

$$
B=\frac{\lambda \rho \sigma_{2} \sqrt{\pi}}{E k_{2}}\left(\frac{D k_{1}}{\lambda \rho \sigma a_{1} \sqrt{\pi}} F_{2}\left(\frac{\sigma}{a_{1}}\right)-1\right) .
$$

$19 \omega,$| $(\mathrm{R} 23)$ |
| :---: |
| $(\mathrm{R} 29)$ |$\omega, \quad \omega=c_{1}, k_{2} W\left(\frac{\sigma}{a_{1}}\right), \quad c_{2}=\frac{\lambda}{E} \frac{H_{16}(B) H_{57}\left(\frac{\sigma}{\sigma_{1}}\right)}{W\left(\frac{\sigma}{a_{1}}\right)}, \quad k_{2}=\frac{\lambda k_{1} \sqrt{\pi}}{E c_{1}} \frac{W\left(\frac{\sigma}{a_{1}}\right) H_{57}\left(\frac{\sigma}{a_{1}}\right)}{H_{1}(B)}$

for any $B>0$.
$20 \quad \omega, \rho, k_{1} \quad-\quad \omega=\sigma H_{25}\left(\xi_{1}\right), \quad \rho=\frac{h_{0} \sqrt{\pi}}{D \sigma c_{1}} H_{20}\left(\xi_{1}\right), \quad k_{1}=\frac{\sigma h_{0} \sqrt{\pi}}{D} H_{24}\left(\xi_{1}\right)$
where $\xi_{1}$ is the unique solution of the equation

|  |
| :---: |
| $21 \quad \omega, c_{40}(x)=H_{41}(x), \quad x>0$. |

where $\xi_{1}$ is the unique solution of the equation

$$
\frac{h_{0}}{\lambda \rho \sigma} H_{27}(x)=H_{42}(x), \quad x>0
$$

$22 \omega, \lambda, c_{1}$| $(\mathrm{R} 10)$ |
| :--- | :--- |
| $(\mathrm{R} 13)$ |
| $(\mathrm{R} 31)$ |$\quad \omega=\sigma H_{25}\left(\xi_{1}\right), \quad \lambda=\frac{h_{0}}{\rho \sigma} \frac{H_{26}\left(\xi_{1}\right)}{G\left(\xi_{1}\right)}, \quad c_{1}=\frac{k_{1}}{\rho \sigma^{2}} \xi_{1}^{2}$

where $\xi_{1}$ is given as in case 10.

Table 2. contd.

| Case | Unknown <br> coefficients | Restrictions | Solution |
| :---: | :---: | :---: | :---: |
| 23 | $\omega, \gamma, k_{2}$ | (R23) | $\omega=a_{1} W\left(\frac{\sigma}{a_{1}}, \gamma\right), \quad k_{2}=\frac{k_{1} c_{2}}{c_{1}} \frac{W^{2}\left(\frac{\sigma}{a_{1}}, \gamma\right)}{B^{2}}$ |
|  | (R43) |  |  |
|  |  | for any $0<\gamma<\frac{2 D F_{2}\left(\frac{\sigma}{a_{1}}\right)}{\left(1-\epsilon+\frac{E c_{2}}{\lambda}\right) \sqrt{\pi}} H_{50}\left(\frac{\sigma}{a_{1}}\right)$ |  |

where $B=B(\gamma)$ is the unique solution of the equation

| $\frac{1}{H_{16}(x)}=\frac{\lambda}{E c_{2}} \frac{H_{57}\left(\frac{\sigma}{a_{1}}, \gamma\right)}{W\left(\frac{\sigma}{a_{1}}, \gamma\right)}, \quad x>0$. |  |
| :---: | :---: | :---: |
| $24 \quad \omega, \epsilon, c_{1} \quad$$(\mathrm{R} 16)$ <br> $(\mathrm{R} 35)$ | $\omega=\sigma H_{25}\left(\xi_{1}\right), \quad c_{1}=\frac{k_{1}}{\rho \sigma^{2}} \xi_{1}^{2}, \quad \epsilon=1-\frac{2 D}{\gamma \sqrt{\pi}} H_{29}\left(\xi_{1}\right) H_{58}\left(\xi_{1}\right)$ |

where $\xi_{1}$ is given as in case 10 .

| 25 | $\omega, c_{1}, c_{2}$ | $(\mathrm{R} 10)$ <br> $(\mathrm{R} 14)$ <br> (R34) |
| :--- | :--- | :--- |$\quad \omega=\sigma H_{25}\left(\xi_{1}\right), \quad c_{1}=\frac{k_{1}}{\rho \sigma^{2}} \xi_{1}^{2}, \quad c_{2}=\frac{k_{2}}{\rho \sigma^{2}} \frac{B^{2}}{H_{25}^{2}\left(\xi_{1}\right)}$

where $\xi_{t}$ is given as in case 10 and $B$ is given as in case 12.
$26 \quad \omega, \rho, c_{1} \quad(\mathrm{R} 10) \quad \omega=\sigma H_{25}\left(\xi_{1}\right), \quad \rho=\frac{k_{2}}{\sigma^{2} c_{2}} \frac{B^{2}}{H_{25}^{2}\left(\xi_{1}\right)}, \quad c_{1}=\frac{k_{1} c_{2}}{k_{2}} \frac{W^{2}\left(\xi_{1}\right)}{B^{2}}$
where $\xi_{1}$ is given as in case 10 and $B$ is the only solution of the equation

where $\xi_{1}$ is given as in case 10 and $B$ is given as in case 14.

| $28 \omega, \lambda, c_{2}(\mathrm{R} 23)$ | $\omega=a_{1} W\left(\frac{\sigma}{a_{1}}\right), \quad c_{2}=\frac{c_{1} k_{2}}{k_{1}} \frac{B^{2}}{W^{2}\left(\frac{\sigma}{a_{1}}\right)}$, |  |
| :---: | :---: | :---: |
| $\lambda=\frac{1}{G\left(\frac{\sigma}{a_{1}}\right)}\left(\frac{D c_{1}}{\sqrt{\pi}} F_{2}\left(\frac{\sigma}{a_{1}}\right)-\frac{E c_{1} k_{2}}{k_{1} \sqrt{\pi}} \frac{H_{1}(B)}{W\left(\frac{\sigma}{a_{1}}\right)}\right)$ |  |  |
| 29 | $\omega, \rho, c_{2}$ | (R10) |
|  | (R38) | $\omega=\sigma H_{25}\left(\xi_{1}\right), \quad \rho=\frac{k_{1}}{\sigma^{2} c_{1}} \xi_{1}^{2}, \quad c_{2}=\frac{c_{1} k_{2}}{k_{1}} \frac{B^{2}}{W^{2}\left(\xi_{1}\right)}$ |

where $\xi_{1}$ is given as in case 10 and $B$ is the only solution of the equation

$$
H_{1}(x)=\frac{\lambda k_{1} \sqrt{\pi}}{E c_{1} k_{2}} H_{25}\left(\xi_{1}\right) H_{45}\left(\xi_{1}\right), \quad x>0 .
$$

Table 2. contd.

| Case | Unknown coefficients | Restrictions | Solution |
| :---: | :---: | :---: | :---: |
| 30 | $\omega, \boldsymbol{\epsilon}, c_{2}$ | $\begin{aligned} & \text { (R23) } \\ & \text { (R39) } \end{aligned}$ | $\begin{gathered} \omega=a_{1} W\left(\frac{\sigma}{a_{1}}\right), \quad c_{2}=\frac{c_{1} k_{2}}{k_{1}} \frac{B^{2}}{W^{2}\left(\frac{\sigma}{a_{1}}\right)}, \\ \epsilon=\frac{2 D}{\gamma \sqrt{\pi}} W\left(\frac{\sigma}{a_{1}}\right) F_{2}\left(\frac{\sigma}{a_{1}}\right)-\frac{2 D}{\lambda \gamma \sqrt{\pi}} F_{2}\left(\frac{\sigma}{a_{1}}\right)\left(\frac{D c_{1}}{\sqrt{\pi}} F_{2}\left(\frac{\sigma}{a_{1}}\right)-\frac{E c_{1} k_{2}}{k_{1} \sqrt{\pi}} \frac{H_{1}(B)}{W\left(\frac{\sigma}{a_{1}}\right)}\right) \end{gathered}$ <br> for any $H_{1}^{-1}(A)<B<H_{1}^{-1}(C)$ if $\frac{\sigma}{a_{1}}<x_{46}$ <br> or for any $0<B<H_{1}^{-1}(C)$ if $x_{46} \leq \frac{\sigma}{a_{1}}<x_{47}$ <br> where $A=\frac{\lambda k_{1} \sqrt{\pi}}{E c_{1} k_{2}} W\left(\frac{\sigma}{a_{1}}\right) H_{46}\left(\frac{\sigma}{a_{1}}\right)$ $C=\frac{\lambda k_{1} \sqrt{\pi}}{E c_{1} k_{2}} W\left(\frac{\sigma}{a_{1}}\right) H_{47}\left(\frac{\sigma}{a_{1}}\right) .$ |
| 31 | $\omega, \gamma, c_{2}$ | $\begin{aligned} & \text { (R23) } \\ & \text { (R40) } \end{aligned}$ | $\omega=a_{1} W\left(\frac{\sigma}{a_{1}}, \gamma\right), \quad c_{2}=\frac{k_{2}}{\rho a_{1}^{2}} \frac{B^{2}}{W^{2}\left(\frac{\sigma}{a_{1}}, \gamma\right)},$ <br> for any $0<\gamma<\frac{2 D H_{58}\left(\frac{\sigma}{a_{1}}\right)}{(1-\epsilon) \sqrt{\pi}}\left(\frac{D c_{1}}{\lambda \sqrt{\pi}} \frac{1}{H_{55}\left(\frac{\sigma}{a_{1}}\right)}-1\right)$ <br> where $B=B(\gamma)$ is the unique solution of the equation $H_{1}(x)=\frac{\lambda k_{1} \sqrt{\pi}}{E c_{1} k_{2}} W\left(\frac{\sigma}{a_{1}}, \gamma\right) H_{57}\left(\frac{\sigma}{a_{1}}, \gamma\right), \quad x>0 .$ |
| 32 | $\omega, \lambda, \rho$ | $\begin{aligned} & \text { (R10) } \\ & \text { (R42) } \end{aligned}$ | $\omega=\sigma H_{25}\left(\xi_{1}\right), \quad \rho=\frac{k_{1}}{\sigma^{2} c_{1}} \xi_{1}^{2}, \quad \lambda=\frac{\sigma h_{0} c_{1}}{k_{1}} \frac{H_{49}\left(\xi_{1}\right)}{G\left(\xi_{1}\right)}$ <br> where $\xi_{1}$ is given as in case 10 . |
| 33 | $\omega, \lambda, \epsilon$ | $\begin{aligned} & \text { (R23) } \\ & \text { (R48) } \end{aligned}$ | $\begin{aligned} & \qquad \omega=a_{1} W\left(\frac{\sigma}{a_{1}}\right), \quad \epsilon=1-\frac{2 D}{\gamma \sqrt{\pi}} H_{58}\left(\frac{\sigma}{a_{1}}\right)\left(\frac{D c_{1}}{\lambda \sqrt{\pi}} H_{56}\left(\frac{\sigma}{a_{1}}\right)-1\right) \\ & \text { for any } \frac{\frac{D c_{1}}{\sqrt{\pi}} H_{56}\left(\frac{\sigma}{a_{1}}\right)}{1+\frac{\gamma \sqrt{\pi}}{2 D} \frac{1}{H_{56}\left(\frac{\sigma}{a_{1}}\right)}}<\lambda<\frac{D c_{1}}{\sqrt{\pi}} H_{56}\left(\frac{\sigma}{a_{1}}\right) \end{aligned}$ |
| 34 | $\omega, \lambda, \gamma$ | (R23) <br> (R44) <br> (R45) | $\begin{gathered} \omega=a_{1} W\left(\frac{\sigma}{a_{1}}, \gamma\right), \\ \lambda=\frac{D c_{1}}{G\left(\frac{\sigma}{a_{1}}, \gamma\right), \sqrt{\pi}}\left(F_{2}\left(\frac{\sigma}{a_{1}}\right)-\frac{E}{D} \sqrt{\frac{c_{2} k_{1}}{c_{1} k_{1}}} F_{1}\left(\frac{\sigma}{a_{1}} H_{25}\left(\frac{\sigma}{a_{1}}\right)\right)\right) \\ \text { for any } 0<\gamma<\frac{2 D H_{5 s}\left(\frac{\sigma}{a_{1}}\right)}{(1-\epsilon)^{\sqrt{\pi}}}\left(\frac{a_{2}}{\sigma} F_{1}^{-1}\left(\frac{D k_{1} a_{2}}{E a_{1} k_{2}} F_{2}\left(\frac{\sigma}{a_{1}}\right)\right)-1\right) . \end{gathered}$ |
| 35 | $\omega, \boldsymbol{\epsilon}, \boldsymbol{\rho}$ | $\begin{aligned} & \text { (R46) } \\ & \text { (R47) } \end{aligned}$ | $\omega=\sigma H_{25}\left(\xi_{1}\right), \quad \rho=\frac{k_{1}}{\sigma^{2} c_{1}} \xi_{1}^{2}, \quad \epsilon=1-\frac{2 D}{\gamma \sqrt{\pi}} \frac{H_{51}\left(\xi_{1}\right)}{H_{55}\left(\xi_{1}\right)}$ |

where $\xi_{1}$ is given as in case 10 .

$36 \omega, \gamma, \rho \quad$| (R10) |
| :--- |
| (R41) |

$$
\begin{aligned}
& \omega=\sigma \sqrt{\frac{c_{1} k_{2}}{k_{1} c_{2}} \frac{B}{\xi_{1}}}, \quad \rho=\frac{k_{1}}{\sigma^{2} c_{1}} \xi_{1}^{2}, \\
& \gamma=\frac{2 D}{\sqrt{\pi}}\left(\sqrt{\frac{c_{1} k_{2}}{k_{1} c_{2}}} B-\xi_{1}\right) F_{2}\left(\xi_{1}\right)
\end{aligned}
$$

where $\xi_{1}$ is given as in case 10 and $B$ is the only solution of the equation

$$
H_{53}(x)=\frac{\sigma h_{0} c_{1}}{\lambda k_{1}} \frac{\exp \left(-\xi_{1}^{2}\right)}{\xi_{1}}-\epsilon \xi_{1}, \quad x>\sqrt{\frac{k_{1} c_{2}}{c_{1} k_{2}}} \xi_{1}
$$

