

SUFFICIENT CONDITIONS TO OBTAIN A STEADY STATE-STEFAN PROBLEM WITH INTERNAL ENERGY AND DIRICHLET AND ROBIN BOUNDARY CONDITIONS.

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ABSTRACT: We consider the problem of the steady-state temperature distribution of a material submitted to an internal energy g . We assume the material is contained in a regular bounded domain $\Omega \subset \mathbb{R}^n$, $\partial\Omega = \Gamma_1 \cup \Gamma_2 \cup \Gamma_3$, with a fixed positive temperature on Γ_1 and a Robin condition for the heat flux on Γ_2 (or Newton's type law). We obtain monotonicity properties for the temperature and we state two different sufficient conditions in order it is of non constant sign in Ω , that is, we have steady state two-phase Stefan problem. We also show an example where the necessary and sufficient conditions are given.

RESUMEN: Consideramos el problema de la distribución estacionaria de temperatura de un material sometido a una energía interna g . Asumimos que el material está contenido en un dominio regular acotado $\Omega \subset \mathbb{R}^n$, $\partial\Omega = \Gamma_1 \cup \Gamma_2 \cup \Gamma_3$, con temperatura positiva fija sobre Γ_1 y tal que el flujo satisface sobre Γ_2 una condición de Robin (o ley de tipo Newton). Demostramos propiedades de monotonía para la temperatura y establecemos dos tipos diferentes de condiciones suficientes para que la temperatura no sea de signo constante en Ω , es decir, para tener un problema estacionario de Stefan a dos fases. También mostramos un ejemplo donde se dan condiciones necesarias y suficientes.

KEY WORDS : steady-state Stefan problem, variational inequalities, mixed boundary conditions.

AMS SUBJECT CLASSIFICATION : 35R35, 35J05, 35J85.

1. INTRODUCTION

We consider the problem of the steady-state temperature distribution of a body or a container with a material which is submitted to an internal energy g . We assume the body to be a regular bounded domain $\Omega \subset \mathbb{R}^n$, with $\partial\Omega = \Gamma_1 \cup \Gamma_2 \cup \Gamma_3$ and Γ_i are sufficiently regular disjoint portions of $\partial\Omega$ for $i = 1, 2, 3$ and Γ_1 and Γ_2 have positive $(n - 1)$ dimensional measure. For simplicity we state the phase change temperature of the material occupying Ω to be $0^\circ C$.

We fix the temperature on Γ_1 to be positive and the heat flux on Γ_2 to verify a Newton's type law (or a Robin condition) and to be null on Γ_3 .

In this paper we study necessary or sufficient conditions for data such that we have two phases of the material in $\overline{\Omega}$ (i.e. the temperature is of non constant sign in $\overline{\Omega}$). ([3], [5], [8], [9]).

The weak formulation of our problem is given by

$$(1) \quad \begin{cases} -\Delta u = g & \text{in } \mathcal{D}'(\Omega) \\ u|_{\Gamma_1} = B > 0 \\ -\frac{\partial u}{\partial n}\Big|_{\Gamma_2} = \alpha(u - T) \\ \frac{\partial u}{\partial n}\Big|_{\Gamma_3} = 0 \end{cases}$$

where $u = k_2\theta^+ - k_1\theta^-$, θ is the temperature, k_1 and k_2 are respectively the thermal conductivities of solid and liquid phase, $b = \frac{B}{k_2}$ and T are fixed temperatures on Γ_1 and Γ_2 respectively, α is a heat transfer coefficient on Γ_2 and n denotes the exterior normal to Γ_2 . ([4])

It is well known ([7], [10]), that the variational formulation of (1) is given by

$$(2) \quad \begin{cases} a_\alpha(u, v) = L_{\alpha T}g(v), & \forall v \in V_1 \\ u \in K_B \end{cases}$$

where

$$(3) \quad a_\alpha(u, v) = \int_{\Omega} \nabla u \nabla v \, dx + \alpha \int_{\Gamma_2} \gamma_2(u) \gamma_2(v) \, ds$$

$$(4) \quad L_{\alpha T}g(v) = \int_{\Omega} gv \, dx + \alpha \int_{\Gamma_2} T \gamma_2(v) \, ds$$

with

$$(5) \quad V = H^1(\Omega), \quad V_1 = \{v \in V : \gamma_1(v) = 0\}, \quad K_B = \{v \in V : \gamma_1(v) = B\}$$

and γ_i are the 0-order trace operator on Γ_i for $i = 1, 2$.

The bilinear form a_1 is coercive on V , i.e.

$$(6) \quad \exists \lambda_1 > 0 / \quad a_1(v, v) \geq \lambda_1 \|v\|_V^2, \quad \forall v \in V$$

and therefore a_α is also coercive on V with a constant of coerciveness given by $\lambda_\alpha = \lambda_1 \cdot \min(1, \alpha)$. ([7], [10])

Moreover, if $\alpha > 0$, $g \in L^2(\Omega)$, $B \in H^{\frac{1}{2}}(\Gamma_1)$ and $T \in H^{\frac{1}{2}}(\Gamma_2)$, problem (2) has a unique solution $u_{\alpha g T B} \in K_B$ ([4], [7], [10]). In the sequel we will study the solution operator

$$\Lambda : \mathbb{R}^+ \times L^2(\Omega) \times H^{\frac{1}{2}}(\Gamma_2) \times H^{\frac{1}{2}}(\Gamma_1) \longrightarrow H^1(\Omega)$$

such that

$$(7) \quad (\alpha, g, T, B) \longmapsto \Lambda(\alpha, g, T, B) = u_{\alpha g T B}.$$

It is well known that some difficulties due to the mixed boundary conditions do arrives to prove regularity of $u_{\alpha g T B}$ ([6]). Sufficient conditions to obtain a H^2 regularity for elliptic mixed boundary problems are given, among others, in [1] and, recently in [2], but our interest is to state hypothesis in order to have a solution of non-constant sign.

We obtain monotonicity properties of $\Lambda(\alpha, g, T, B)$ and we state two different sufficient conditions for it to be of non constant sign in Ω (see Theorems 8 and 10).

Finally we check everything in an example for which we obtain the necessary and sufficient conditions in order to have a steady-state two-phase Stefan problem in Ω .

Throughout this paper, the statement $f \geq 0$ where f belongs to any subspace of $L^2(A)$ will mean $f(x) \geq 0$ a.e. in A , and in order to simplify notation, on Γ_1 and Γ_2 , u will mean $\gamma_1(u)$ and $\gamma_2(u)$ respectively.

2. ASYMPTOTIC BEHAVIOUR OF $\Lambda(\alpha, g, T, B)$

It is physically reasonable to expect that if $\alpha \rightarrow +\infty$ then $u_{\alpha gTB}|_{\Gamma_2} \rightarrow T$.

This fact leads us to consider the function $u = u_{gTB}$, which is the unique solution of the mixed partial differential problem

$$(8) \quad \begin{cases} -\Delta u = g & \text{in } \mathcal{D}'(\Omega) \\ u|_{\Gamma_1} = B \\ u|_{\Gamma_2} = T \\ \frac{\partial u}{\partial n}|_{\Gamma_3} = 0 \end{cases}$$

whose equivalent variational formulation is given by:

$$(9) \quad \begin{cases} a(u, v) = L_g(v), & \forall v \in V_{12} \\ u \in K_{BT} \end{cases}$$

where

$$(10) \quad a(u, v) = \int_{\Omega} \nabla u \nabla v \, dx, \quad L_g(v) = \int_{\Omega} g v \, dx$$

with

$$(11) \quad V_{12} = \{v \in V / v = 0 \text{ on } \Gamma_1 \cup \Gamma_2\}$$

$$(12) \quad K_{BT} = \{v \in V / v = B \text{ on } \Gamma_1 \text{ and } v = T \text{ on } \Gamma_2\}$$

Theorem 1. *If $\alpha > 1$, $g \in L^2(\Omega)$, $B \in H^{\frac{1}{2}}(\Gamma_1)$ and $T \in H^{\frac{1}{2}}(\Gamma_2)$, $u_{\alpha gTB}$ and u_{gTB} are respectively the unique solutions of problems (2) and (9) then: 1)*

$$\|u_{\alpha gTB} - u_{gTB}\|_{H^1(\Omega)} \leq \frac{k}{\lambda_1}$$

$$2) \quad (\alpha - 1) \int_{\Gamma_2} (u_{\alpha gTB} - u_{gTB})^2 \, ds \leq \frac{k^2}{\lambda_1}$$

where k is a constant which depends on u_{gTB} and λ_1 depends only on Ω and Γ_2 .

$$3) \quad \lim_{\alpha \rightarrow +\infty} u_{\alpha gTB} = u_{gTB} \text{ strongly in } V.$$

Proof. We name $u_{\alpha} = \Lambda(\alpha, g, T, B)$ and $u = u_{gTB}$, then, as $u \in K_{BT} \subset K_B$, we can choose $v = u$ in (2) and (7) and we obtain for $w_{\alpha} = u_{\alpha} - u$

$$(13) \quad \lambda_1 \|w_{\alpha}\|_{H^1}^2 + (\alpha - 1) \int_{\Gamma_2} w_{\alpha}^2 \, ds \leq a_{\alpha}(w_{\alpha}, w_{\alpha}) = -a(u, w_{\alpha}) + L_g(w_{\alpha})$$

$$\leq (\|u\|_V + \|g\|_{L_2}) \|w_\alpha\|_V$$

Then, if we call $k = \|u\|_V + \|g\|_{L_2}$ we have

$$(14) \quad \forall \alpha > 1 \quad \|w_\alpha\|_V \leq \frac{k}{\lambda_1} \quad \text{and} \quad (\alpha - 1) \int_{\Gamma_2} w_\alpha^2 ds \leq \frac{k^2}{\lambda_1}.$$

These a priori estimates ensure the existence of $u^* \in V$ and a convergent subsequence of u_α , which for simplicity we will also name u_α such that $u_\alpha \rightharpoonup u^*$ weakly in V and, moreover,

$$\liminf_{\alpha \rightarrow +\infty} \int_{\Gamma_2} (u_\alpha - T)^2 ds = 0.$$

The functional $v \rightarrow \int_{\Gamma_2} v^2 ds$ is weakly lower semicontinuous in V , then

$$0 \leq \int_{\Gamma_2} (u^* - T)^2 ds \leq \liminf_{\alpha \rightarrow +\infty} \int_{\Gamma_2} (u_\alpha - T)^2 ds = 0$$

and therefore $u^* \in K_B$. The zero order trace operator γ_1 on Γ_1 is continuous, then $u^* \in K_{BT}$, and $u = u^*$ by uniqueness of the solution of the variational equation (9). Thus, we have proved $\lim_{\alpha \rightarrow +\infty} u_{\alpha gTB} = u_{gTB}$ weakly in V . Besides, $w_\alpha \rightharpoonup 0$ weakly in V and (13) implies $w_\alpha \rightarrow 0$ strongly in V , i.e. $\lim_{\alpha \rightarrow +\infty} u_{\alpha gTB} = u_{gTB}$ strongly in V . \square

3. MONOTONICITY PROPERTIES OF $\Lambda(\alpha, g, T, B)$

We will prove the operator $\Lambda(\alpha, g, T, B)$ satisfies some monotonicity properties and order bound.

Proposition 1. *If $u_i = \Lambda(\alpha, g_i, T_i, B_i)$, $i = 1, 2$, $g_1 \leq g_2$ in $L^2(\Omega)$, $T_1 \leq T_2$ in $H^{\frac{1}{2}}(\Gamma_2)$, $B_1 \leq B_2$ in $H^{\frac{1}{2}}(\Gamma_1)$, then $u_1 \leq u_2$ a.e. in Ω .*

Proof. i) We first consider the case $g_1 = g_2 = g$ and $T_1 = T_2 = T$. $B_1 \leq B_2 \implies w = (u_1 - u_2)^+ \in V_1$ and $0 \leq a_\alpha(u_1 - u_2, w) = a(w, w) + \alpha \int_{\Gamma_2} w^2 ds = 0$.

Then, $w = 0$, that is $u_1 \leq u_2$ a.e. in Ω .

ii) Now we state $B_1 = B_2 = B$, $g_1 \leq g_2$ and $T_1 \leq T_2$. As $w = (u_1 - u_2)^+ \in V_1$, $0 \leq a_\alpha(u_1 - u_2, w) = a_\alpha(w, w) = \int_{\Omega} (g_1 - g_2) w dx + \alpha \int_{\Gamma_2} (T_1 - T_2) w ds \leq 0$ and $u_1 \leq u_2$ a.e. in Ω . The general case easily follows from (i) and (ii). \square

Proposition 2. *Let $T = T_0$ be constant on Γ_2 , $\alpha \in \mathbb{R}^+$.*

a) *If $g \leq 0$ in $L^2(\Omega)$, $B \leq T_0$ in $H^{\frac{1}{2}}(\Gamma_1)$, then $u_{\alpha g T_0 B} = \Lambda(\alpha, g, T_0, B) \leq T_0$ a.e. in Ω .*

b) *If $g \geq 0$ in $L^2(\Omega)$, $B \geq T_0$ in $H^{\frac{1}{2}}(\Gamma_1)$, then $u_{\alpha g T B} \geq T_0$ a.e. in Ω .*

c) *If $g = 0$, $B = T_0$, then $u_{\alpha g T B} = B = T_0$.*

Proof. a) $a_\alpha(u_{\alpha g T_0 B} - T_0, v) = \int_{\Omega} g v dx, \quad \forall v \in V_1$.

$B \leq T_0 \implies w = (u_{\alpha g T_0 B} - T_0)^+ \in V_1$ and then $0 \leq a_\alpha(w, w) = \int_{\Omega} g w dx \leq 0$.

Thus $u_{\alpha g T_0 B} \leq T_0$.

b) and c) are obtained similarly. \square

An immediate consequence of this property and Theorem 1 is the following:

Proposition 3. *Let $T = T_0$ be constant on Γ_2 , $u_{\alpha_i} = \Lambda(\alpha_i, g, T_0, B)$, $i = 1, 2$ and $1 < \alpha_1 \leq \alpha_2$.*

a) If $g \leq 0$ in $L^2(\Omega)$, $B \leq T_0$ in $H^{\frac{1}{2}}(\Gamma_1)$, then $u_{\alpha_1} \leq u_{\alpha_2}$ and $u_{\alpha g T_0 B} \uparrow u_{g T_0 B}$ strongly in V .

b) If $g \geq 0$ in $L^2(\Omega)$, $B \geq T_0$ in $H^{\frac{1}{2}}(\Gamma_1)$, then $u_{\alpha_2} \leq u_{\alpha_1}$ and $u_{\alpha g T_0 B} \downarrow u_{g T_0 B}$ strongly in V .

Proof. We consider $w = (u_1 - u_2)^+ \in V_1$. Then, as $u_1 \leq T_0$ by Prop. 3, we have

$$0 \leq a(w, w) = a(u_1 - u_2, w) = \int_{\Gamma_2} [\alpha_1 (T_0 - u_1) - \alpha_2 (T_0 - u_2)] w \, ds =$$

$$(\alpha_1 - \alpha_2) \int_{\Gamma_2} (T_0 - u_1) w \, ds - \alpha_2 \int_{\Gamma_2} w^2 \, ds \leq 0.$$

Then $u_{\alpha_1} \leq u_{\alpha_2}$.

b) is obtained in a similar way. \square

The next result easily follows from propositions 2 and 3.

Proposition 4. *Let $T = T_0$ be constant on Γ_2 and $g_1 \leq 0 \leq g_2$ in $L^2(\Omega)$, $T_1 \leq T_0 \leq T_2$ in $H^{\frac{1}{2}}(\Gamma_2)$, $B_1 \leq T_0 \leq B_2$ in $H^{\frac{1}{2}}(\Gamma_1)$, then $u_{\alpha_1 g_1 T_1 B_1} \leq T_0 \leq u_{\alpha_2 g_2 T_2 B_2}$ a.e. in Ω , $\forall \alpha_1, \alpha_2 \in \mathbb{R}^+$.*

4. NECESSARY AND SUFFICIENT CONDITIONS FOR THE EXISTENCE OF TWO PHASES IN Ω

From now on we will consider without loss of generality $B \geq 0$ as in the Introduction.

It is physically reasonable to expect that, in the case $g \geq 0$, $T \geq 0$ and $B \geq 0$ the solution $u_{\alpha g T B}$ will be non negative, but this result is a trivial consequence of proposition 3, taking into account that $u_{\alpha 0 0 0} = 0$. Thus, denying the above statement we have the following:

Proposition 5. *If $B \geq 0$, a necessary condition for the existence of two phases in Ω ($u_{\alpha g T B}$ is of non constant sign in Ω) is that $g \leq 0$ or $T \leq 0$, non simultaneously null.*

In order to obtain sufficient conditions we will first analyze the case of constant data.

Lemma 1. *Let g, T, B be constant, then*

$$u_{\alpha g T B} = \Lambda(\alpha, g, T, B) = T + (B - T) U_{1\alpha} + g U_{2\alpha}$$

where $U_{1\alpha} = \Lambda(\alpha, 0, 0, 1)$ and $U_{2\alpha} = \Lambda(\alpha, 1, 0, 0)$ verify:

i) $0 \leq U_{1\alpha} \leq 1$ in $L^2(\Omega)$.

ii) $U_{1\alpha} \downarrow u_{001}$ strongly in V when $\alpha \rightarrow +\infty$.

iii) $\forall 0 < \varepsilon < 1$, $\exists \Gamma_{\alpha,\varepsilon}^1 \subset \Gamma_2$ with $|\Gamma_{\alpha,\varepsilon}^1| = \int_{\Gamma_{\alpha,\varepsilon}^1} ds > 0$ such that

$U_{1\alpha} \leq \varepsilon$ a.e. on $\Gamma_{\alpha,\varepsilon}^1 \forall \alpha \geq \frac{C_1}{\varepsilon^2}$ where $C_1 = 1 + \frac{\|u_{001}\|_V^2}{\lambda_1 |\Gamma_2|}$ depends only on Ω , Γ_1 and Γ_2 .

iv) $0 \leq U_{2\alpha}$ in $L^2(\Omega)$.

v) $U_{2\alpha} \downarrow u_{100}$ in $L^2(\Omega)$ when $\alpha \rightarrow \infty$.

vi) $\forall \varepsilon > 0$, $\exists \Gamma_{\alpha,\varepsilon}^2 \subset \Gamma_2$ with $|\Gamma_{\alpha,\varepsilon}^2| = \int_{\Gamma_{\alpha,\varepsilon}^2} ds > 0$ such that

$U_{2\alpha} \leq \varepsilon$ a.e. on $\Gamma_{\alpha,\varepsilon}^2 \forall \alpha \geq \frac{C_2}{\varepsilon^2}$ where $C_2 = 1 + \frac{(\|u_{001}\|_V + |\Omega|)^2}{\lambda_1 |\Gamma_2|}$ depends only on Ω , Γ_1 and Γ_2 .

vii) $U_{2\alpha} \in L^\infty(\Omega)$.

viii) $\forall \alpha > 1$, $\exists \Omega_1 \subset \Omega$ with $|\Omega_1| = \int_{\Omega_1} dx > 0$ such that

$U_{2\alpha} \geq C_3$ a.e. in Ω_1 where $C_3 = \frac{\lambda_1 \|u_{100}\|_V^2}{|\Omega|}$ depends only on Ω and Γ_2 .

Proof. It is well known from the linearity of Problem I that, if g , T and B are constant

$$u_{\alpha gTB} = B\Lambda(\alpha, 0, 0, 1) + T\Lambda(\alpha, 0, 1, 0) + g\Lambda(\alpha, 1, 0, 0)$$

and moreover,

$$\Lambda(\alpha, 0, 1, 0) = 1 - \Lambda(\alpha, 0, 0, 1).$$

Then, we have

$$u_{\alpha gTB} = T + (B - T)U_{1\alpha} + gU_{2\alpha}$$

where $U_{1\alpha} = \Lambda(\alpha, 0, 0, 1)$ and $U_{2\alpha} = \Lambda(\alpha, 1, 0, 0)$ are solutions of (2).

(i), (ii), (iv) and (v) are obviously consequences of the maximum principle or of propositions 2-4.

Besides, from Theorem 1

$$(\alpha - 1) \int_{\Gamma_2} (U_{1\alpha} - u_{001})^2 ds = (\alpha - 1) \int_{\Gamma_2} U_{1\alpha}^2 ds \leq \frac{k^2}{\lambda_1}.$$

Then, for any fixed $0 < \varepsilon < 1$, if $\alpha \geq \frac{C_1}{\varepsilon^2} > 1 + \frac{k^2}{\lambda_1 \varepsilon^2 |\Gamma_2|}$ where $C_1 = 1 + \frac{\|u_{001}\|_V^2}{\lambda_1 |\Gamma_2|}$.

It results

$$\int_{\Gamma_2} U_{1\alpha}^2 ds \leq \varepsilon^2 |\Gamma_2|.$$

$\therefore \exists \Gamma_{\alpha,\varepsilon}^1$ with $|\Gamma_{\alpha,\varepsilon}^1| > 0$ such that $U_{1\alpha} \leq \varepsilon$ a.e. on $\Gamma_{\alpha,\varepsilon}^1$.

A similar result (vi) holds for $U_{2\alpha}$, where the constant C_1 is replaced by $C_2 = 1 + \frac{(\|u_{001}\|_V + |\Omega|)^2}{\lambda_1 |\Gamma_2|}$.

In order to prove vii) we will seek for $F \in L^\infty(\bar{\Omega})$ such that $U_{2\alpha} \leq F$.

We propose a polynomial

$$F(x) = -\frac{1}{2}(x_1)^2 + \sum_{i=1}^n p_i x_i + p_0 + q \in L^\infty(\bar{\Omega})$$

with $-\frac{1}{2}(x_1)^2 + \sum_{i=1}^n p_i x_i + p_0 \geq 0, \forall x \in \bar{\Omega}$.

Then, $\forall q > 0$ it results $F(x) \geq q > 0$ in $\bar{\Omega}$.

Moreover,

$$\left(\frac{\partial F}{\partial n} + \alpha F \right) \Big|_{\Gamma_2} \geq (\nabla F \times n) \Big|_{\Gamma_2} + \alpha q \quad \forall \alpha > 0$$

and

$$|\nabla F \times n| \leq \|\nabla F\| = \left[(-x_1 + p_1)^2 + \sum_{i \neq 1} p_i^2 \right]^{\frac{1}{2}} \leq A \quad \text{on } \partial\Omega.$$

If we choose $q > A$, it results $\left(\frac{\partial F}{\partial n} + \alpha F \right) \Big|_{\Gamma_2} \geq -A + \alpha q > 0, \forall \alpha > 1$.

Then, we obtain from Proposition 2

$$F(x) = \Lambda \left(\alpha, 1, \frac{1}{\alpha} \left(\frac{\partial F}{\partial n} + \alpha F \right) \Big|_{\Gamma_2}, F|_{\Gamma_1} \right) \geq \Lambda(\alpha, 1, 0, 0) = U_{2\alpha} \quad \text{a.e. in } \bar{\Omega}.$$

Therefore $U_{2\alpha} \in L^\infty(\bar{\Omega})$.

To prove viii) we recall that, from (v) we have $\forall \alpha > 1$

$$\int_{\Omega} U_{2\alpha} dx \geq \int_{\Omega} u_{100} dx = a(u_{100}, u_{100}) = a_1(u_{100}, u_{100}) \geq \lambda_1 \|u_{100}\|_V^2.$$

Therefore $\exists \Omega_1 \subset \Omega$ with $|\Omega_1| = \int_{\Omega_1} dx > 0$ such that

$$U_{2\alpha} \geq C_3 = \frac{\lambda_1 \|u_{100}\|_V^2}{|\Omega|} \quad \text{a.e. in } \bar{\Omega}_1.$$

□

Theorem 2. Let g, T, B be constant and $B > 0$. If $g < -\frac{\max(T, B)}{c_3}$, then $u_{\alpha g T B}$ is of non constant sign in $\Omega \forall \alpha > 1$.

Proof. By lemma 7 we have

$$u_{\alpha g T B} = T + (B - T) U_{1\alpha} + g U_{2\alpha}$$

and $\exists \Omega_1 \subset \Omega$ with $|\Omega_1| > 0$ such that $U_{2\alpha} \geq c_3$ a.e. in Ω_1 .

For $g < 0$ we obtain

i) If $T \leq B$, $u_{\alpha g T B} \leq B + g U_{2\alpha} \leq B + g C_3$ a.e. in $\bar{\Omega}_1$.

ii) If $T > B$, $u_{\alpha gTB} \leq T + gU_{2\alpha} \leq T + gC_3$ a.e. in $\bar{\Omega}_1$.

Then, if $\max(T, B) < -gC_3$ it results $u_{\alpha gTB}|_{\bar{\Omega}_1} < 0$ and as we have $u_{\alpha gTB}|_{\Gamma_1} = B > 0$ it follows that u is of non constant sign in $\bar{\Omega}$. \square

An immediate consequence of this theorem and the monotone property 2 is the following:

Corollary 1. Let $\alpha > 0$, $g \in L^2(\Omega)$, $B \in H^{\frac{1}{2}}(\Gamma_1)$ and $T \in H^{\frac{1}{2}}(\Gamma_2)$ be such that $g \leq g_0$ a.e. in Ω , $0 < B \leq B_0$ a.e. on Γ_1 and $T \leq B_0$ a.e. on Γ_2 , where the constants g_0 and B_0 satisfy $g_0 < -\frac{B_0}{C_3}$ then $u_{\alpha gTB}$ is of non constant sign in $\bar{\Omega}$.

Theorem 3. Let $g, T, B > 0$ be constant. Then, for all fixed $0 < \varepsilon < 1$ and $\alpha \geq \frac{C_1}{\varepsilon^2}$, $u_{\alpha gTB}$ changes sign in $\bar{\Omega}$ if one of the following conditions are satisfied:

i) $g \leq 0$ and $T < -\frac{\varepsilon B}{1 - \varepsilon}$

or

ii) $g > 0$ and $T < -\frac{\varepsilon B + gH}{1 - \varepsilon}$, with $H = \operatorname{ess\,sup}_{\bar{\Omega}} |F(x)|$ where $F(x)$ is the

polynomial defined in lemma 7 vii).

Proof.

$$u_{\alpha gTB} = T + (B - T)U_{1\alpha} + gU_{2\alpha}$$

We consider $T < 0$ and from lemma 7 iii), if $0 < \varepsilon < 1$ and $\alpha \geq \frac{C_1}{\varepsilon^2}$ there exists $\Gamma_{\alpha, \varepsilon}^1$ such that

$$u_{\alpha gTB} \leq T + (B - T)\varepsilon + gU_{2\alpha} \quad \text{a.e. on } \Gamma_{\alpha, \varepsilon}^1.$$

Then, we have

a) If $g \leq 0$

$$u_{\alpha gTB} \leq (1 - \varepsilon)T + \varepsilon B \quad \text{a.e. on } \Gamma_{\alpha, \varepsilon}^1$$

and therefore, for $T < -\frac{\varepsilon B}{1 - \varepsilon}$ it results $u_{\alpha gTB}|_{\Gamma_{\alpha, \varepsilon}^1} < 0$.

b) If $g > 0$, we have proved in lemma 7 vii) that

$$|U_{2\alpha}|_{L^\infty(\bar{\Omega})} \leq |F(x)|_{L^\infty(\Omega)} = H$$

where H depends only on $\bar{\Omega}$, then

$$u_{\alpha gTB} \leq (1 - \varepsilon)T + \varepsilon B + gH < 0 \quad \text{a.e. on } \Gamma_{\alpha, \varepsilon}^1 \iff T < -\frac{\varepsilon B + gH}{1 - \varepsilon}.$$

\square

Again, an immediate consequence of this theorem and the monotone property 2 is the following

Corollary 2. For any fixed $0 < \varepsilon < 1$ and $\alpha \geq \frac{C_1}{\varepsilon^2}$, let $g \in L^2(\Omega)$, $B \in H^{\frac{1}{2}}(\Gamma_1)$ and $T \in H^{\frac{1}{2}}(\Gamma_2)$ be such that

i) $g \leq g_0 \leq 0$ a.e. in Ω , $0 < B \leq B_0$ a.e. on Γ_1 and $T \leq T_0 < -\frac{\varepsilon B}{1 - \varepsilon}$.

or

ii) $0 \leq g \leq g_0$ a.e. in Ω , $0 < B \leq B_0$ a.e. on Γ_1 and $T \leq T_0 < -\frac{\varepsilon B + H}{1 - \varepsilon}$
 then $u_{\alpha gTB}$ is of non constant sign in $\bar{\Omega}$.

5. EXAMPLE

We consider $\Omega = (0, x_0) \times (0, y_0)$, $\Gamma_1 = \{0\} \times [0, y_0]$, $\Gamma_2 = \{x_0\} \times [0, y_0]$ and $\Gamma_3 = (0, x_0) \times \{0\} \cup (0, x_0) \times \{y_0\}$.

Let g , T and $B > 0$ be constant, then

$$(15) \quad U_{1\alpha} = 1 - \frac{\alpha}{1 + \alpha x_0} x, \quad U_{2\alpha} = -\frac{x^2}{2} + \frac{x_0(2 + \alpha x_0)}{2(1 + \alpha x_0)} x$$

and

$$(16) \quad u_{\alpha gTB} = -\frac{g}{2} x^2 + \frac{(T - B)\alpha + \frac{gx_0}{2}(2 + \alpha x_0)}{(1 + \alpha x_0)} x + B.$$

In this case we can obtain a necessary and sufficient condition for $u_{\alpha gTB}$ to be of non constant sign in Ω . From (16) we have

$$(17) \quad \begin{aligned} u_{\alpha gTB} &= -\frac{g}{2}(x - x_{v_\alpha})^2 + y_{v_\alpha} \\ x_{v_\alpha} &= \frac{(T - B)\alpha + \frac{gx_0}{2}(2 + \alpha x_0)}{g(1 + \alpha x_0)} x + B \\ y_{v_\alpha} &= B + \frac{gx_{v_\alpha}^2}{2}. \end{aligned}$$

Then $u_{\alpha gTB}$ is of non constant sign in Ω if and only if

$$(18) \quad g \geq 0 \text{ and } u_{\alpha gTB}(x_0) < 0$$

or

$$(19) \quad g < 0, y_{v_\alpha} < 0, \text{ and } 0 < x_{v_\alpha} \leq x_0 \text{ or } g \geq 0 \text{ and } u_{\alpha gTB}(x_0) < 0.$$

We have, for $g \geq 0$

$$(20) \quad u_{\alpha gTB}(x_0) < 0 \Leftrightarrow T < -\frac{x_0}{2\alpha} g - \frac{B}{\alpha x_0}.$$

If $g < 0$, we obtain

$$(21) \quad y_{v_\alpha} < 0 \Leftrightarrow T > B - \frac{gx_0}{2\alpha}(2 + \alpha x_0) + (1 + \alpha x_0) \sqrt{(-2Bg)}$$

$$(22) \quad 0 < x_{v_\alpha} \leq x_0 \Leftrightarrow B + \frac{gx_0^2}{2} \leq T < B - \frac{gx_0}{2\alpha}(2 + \alpha x_0)$$

$$(23) \quad u_{\alpha gTB}(x_0) < 0 \Leftrightarrow T < -\frac{x_0}{2\alpha} g - \frac{B}{\alpha x_0}.$$

Taking into account (18) – (23), the necessary and sufficient conditions for $u_{\alpha gTB}$ to be of non constant sign in Ω become, $\forall \alpha > 0$:

$$(24) \quad \begin{aligned} & T < T_{1\alpha}(g) && \text{if } g \geq 0 \\ & \begin{cases} T < T_{1\alpha}(g) \text{ or} \\ T_2(g) \leq T \\ T_{3\alpha}(g) > T \end{cases} && \text{if } -\frac{2B}{x_0^2} < g < 0 \\ & T < T_{3\alpha}(g) && \text{if } g \leq -\frac{2B}{x_0^2}. \end{aligned}$$

where $T_{1\alpha}(g) = -\frac{x_0}{2\alpha}g - \frac{B}{\alpha x_0}$, $T_2(g) = B + \frac{gx_0^2}{2}$, and $T_{3\alpha}(g) = B - \frac{gx_0}{2\alpha}(2 + \alpha x_0)$.

This conditions can also be expressed as:

$$(25) \quad \begin{aligned} g &< g_{1\alpha}(T) = -\frac{2\alpha}{x_0}T - \frac{2B}{x_0^2} && \text{if } T < 0 \\ g &\leq g_2(T) = \frac{2(T-B)}{x_0^2} && \text{if } 0 \leq T < B \\ g &< g_{3\alpha}(T) = -\frac{2(T-B)}{x_0(2+\alpha x_0)} && \text{if } T \geq B. \end{aligned}$$

Then we obtain the following property:

Theorem 4. *If g, T and $B > 0$ are constant and $\Omega = (0, x_0) \times (0, y_0)$, $\Gamma_1 = \{0\} \times [0, y_0]$, $\Gamma_2 = \{x_0\} \times [0, y_0]$ and $\Gamma_3 = (0, x_0) \times \{0\} \cup (0, x_0) \times \{y_0\}$, then the necessary and sufficient conditions in order the temperature $u_{\alpha gTB}$ is of non-constant sign in Ω (we have a steady-state two-phase Stefan problem in Ω) is given by (24) or equivalently by (25).*

We compare (24) or (25) with the sufficient conditions that we have obtained in Theorems 8 and 10.

From Lemma 7, (15) and (16), we can take $C_3 = \frac{x_0^2}{a}$, with $a > 8$, $C_1 = \frac{1}{4x_0}$, $H = \frac{x_0^2}{2}$ and we have:

* From Theorem 8, $\forall \alpha > 0$

$$(26) \quad g < g_4(T) = \begin{cases} -\frac{aB}{x_0^2} & \text{if } T \leq B \\ -\frac{aT}{x_0^2} & \text{if } T > B. \end{cases}$$

* From Theorem 8, $\forall 0 < \varepsilon < 1$,

$$(27) \quad \begin{aligned} \text{i) } T &< T_4(\varepsilon) = -\frac{B}{\frac{1}{\varepsilon} - 1} && \text{if } \alpha \geq \frac{1}{4x_0\varepsilon^2}, g \leq 0 \\ \text{ii) } T &< T_5(\varepsilon) = -\frac{B}{\frac{1}{\varepsilon} - 1} - \frac{x_0^2}{2(1-\varepsilon)}g && \text{if } \alpha \geq \frac{1}{4x_0\varepsilon^2}, g > 0 \end{aligned}$$

It is easy to verify that (26) \implies (25) and (27) \implies (24).

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