# Convergence of optimal control problems governed by second kind parabolic variational inequalities 

Mahdi BOUKROUCHE ${ }^{1}$, Domingo A. TARZIA ${ }^{2 \dagger}$<br>1.Lyon University, UJM F-42023, CNRS UMR 5208, Institut Camille Jordan, 23 Paul Michelon, 42023 Saint-Etienne Cedex 2, France;<br>2.Departamento de Matemática-CONICET, FCE, Austral University, Paraguay 1950, S2000FZF Rosario, Argentina


#### Abstract

We consider a family of optimal control problems where the control variable is given by a boundary condition of Neumann type. This family is governed by parabolic variational inequalities of the second kind. We prove the strong convergence of the optimal control and state systems associated to this family to a similar optimal control problem. This work solves the open problem left by the authors in IFIP TC7 CSMO2011.


Keywords: Parabolic variational inequalities of the second kind; Aubin compactness arguments; Boundary control; Convergence of optimal control problems; Tresca boundary conditions, free boundary problems

## 1 Introduction

The motivation of this paper is to prove the strong convergence of the boundary optimal controls and state systems associated with a family of second kind parabolic variational inequalities. In this paper, we solve the open question, left in [1] and we generalize our work [2], to study the Neumann boundary optimal controls governed by second kind parabolic variational inequalities.

To illustrate the problem, we consider, for example, two free boundary problems which leads to the second kind parabolic variational inequalities.

We assume that the boundary of a multidimensional regular domain $\Omega$ is given by $\partial \Omega=\Gamma_{1} \cup \Gamma_{2} \cup \Gamma_{3}$ with $\operatorname{meas}\left(\Gamma_{1}\right)>0$ and meas $\left(\Gamma_{3}\right)>0$. We consider a family of optimal control problems where the control variable is given by a boundary condition of Neumann type whose state system is governed by a free boundary problem with Tresca conditions on a portion $\Gamma_{2}$ of the boundary, with a flux $f$ on $\Gamma_{3}$ as the control variable, given by

## Problem 1

$$
\begin{aligned}
& \dot{u}-\Delta u=g \text { in } \Omega \times(0, T) \\
& \left|\frac{\partial u}{\partial n}\right|<q \Rightarrow u=0 \text { on } \Gamma_{2} \times(0, T) \\
& \left|\frac{\partial u}{\partial n}\right|=q \Rightarrow \exists k>0: u=-k \frac{\partial u}{\partial n} \text { on } \Gamma_{2} \times(0, T), \\
& u=b \text { on } \Gamma_{1} \times(0, T) \\
& -\frac{\partial u}{\partial n}=f \text { on } \Gamma_{3} \times(0, T)
\end{aligned}
$$

with the initial condition

$$
u(0)=u_{b} \text { on } \Omega
$$

and the compatibility condition on $\Gamma_{1} \times(0, T)$

$$
u_{b}=b \text { on } \Gamma_{1} \times(0, T)
$$

where $q>0$ is the Tresca friction coefficient on $\Gamma_{2}$ [3-5].

We define the spaces $\mathcal{F}=L^{2}\left((0, T) \times \Gamma_{3}\right), V=H^{1}(\Omega)$, $V_{0}=\left\{v \in V: v_{\left.\right|_{\Gamma_{1}}}=0\right\}, H=L^{2}(\Omega), \mathcal{H}=L^{2}(0, T ; H)$, $\mathcal{V}=L^{2}(0, T ; V)$ and the closed convex set $K_{b}=\{v \in V$ : $\left.v_{\left.\right|_{\Gamma_{1}}}=b\right\}$. Let

$$
\left\{\begin{array}{l}
g \in \mathcal{H}, \quad b \in L^{2}\left(0, T, H^{\frac{1}{2}}\left(\Gamma_{1}\right)\right), \quad f \in \mathcal{F}  \tag{1}\\
q \in L^{2}\left((0, T) \times \Gamma_{2}\right), \quad q>0, \quad u_{b} \in K_{b}
\end{array}\right.
$$

The variational formulation of Problem 1 leads to the following parabolic variational problem:

Problem 2 Let $g, b, q, u_{b}$ and $f$ be as in (1). Find $u=u_{f} \in \mathcal{C}(0, T, H) \cap L^{2}\left(0, T ; K_{b}\right)$ with $\dot{u} \in \mathcal{H}$, such that $u(0)=u_{b}$, and for $t \in(0, T)$,

$$
\begin{aligned}
& \langle\dot{u}, v-u\rangle+a(u, u-v)+\Phi(v)-\Phi(u) \\
& \geqslant(g, v-u)-\int_{\Gamma_{3}} f(v-u) \mathrm{d} s, \quad \forall v \in K_{b} .
\end{aligned}
$$

where $(\cdot, \cdot)$ is the scalar product in $H, a$ and $\Phi$ are defined by

$$
\begin{equation*}
a(u, v)=\int_{\Omega} \nabla u \nabla v \mathrm{~d} x, \quad \Phi(v)=\int_{\Gamma_{2}} q|v| \mathrm{d} s . \tag{2}
\end{equation*}
$$

The functional $\Phi$ comes from the Tresca condition on $\Gamma_{2}[4-5]$. We consider also the following problem where we change, in Problem 1, only the Dirichlet condition on $\Gamma_{1} \times(0, T)$ by the Newton law or a Robin boundary condition, i.e.,

## Problem 3

$\dot{u}-\Delta u=g$ in $\Omega \times(0, T)$,

$$
\begin{aligned}
&\left|\frac{\partial u}{\partial n}\right|<q \Rightarrow u=0 \text { on } \Gamma_{2} \times(0, T), \\
&\left|\frac{\partial u}{\partial n}\right|=q \Rightarrow \exists k>0: u=-k \frac{\partial u}{\partial n} \text { on } \Gamma_{2} \times(0, T), \\
&-\frac{\partial u}{\partial n}=h(u-b) \text { on } \Gamma_{1} \times(0, T), \\
&-\frac{\partial u}{\partial n}=f \text { on } \Gamma_{3} \times(0, T),
\end{aligned}
$$

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${ }^{\dagger}$ Corresponding author. E-mail: DTarzia@ austral.edu.ar. Tel.: +54-341-5223093; fax: +54-341-5223001.
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with the initial condition

$$
u(0)=u_{b} \text { on } \Omega
$$

and the condition of compatibility on $\Gamma_{1} \times(0, T)$

$$
u_{b}=b \text { on } \Gamma_{1} \times(0, T) .
$$

The variational formulation of Problem 3 leads to the following parabolic variational problem.

Problem 4 Let $g, b, q, u_{b}$ and $f$ be as in (1). For all $h>0$, find $u=u_{h f}$ in $\mathcal{C}(0, T, H) \cap \mathcal{V}$ with $\dot{u}$ in $\mathcal{H}$, such that $u(0)=u_{b}$, and for $t \in(0, T)$,

$$
\begin{aligned}
& \langle\dot{u}, v-u\rangle+a_{h}(u, u-v)+\Phi(v)-\Phi(u) \\
& \geqslant(g, v-u)-\int_{\Gamma_{3}} f(v-u) \mathrm{d} s+h \int_{\Gamma_{1}} b(v-u) \mathrm{d} s, \\
& \quad \forall v \in V,
\end{aligned}
$$

where $a_{h}$ is defined by

$$
a_{h}(u, v)=a(u, v)+h \int_{\Gamma_{1}} u v \mathrm{~d} s
$$

Moreover, from [6-9] we have that: $\exists \lambda_{1}>0$ such that
$\lambda_{h}\|v\|_{V}^{2} \leqslant a_{h}(v, v), \forall v \in V$, with $\lambda_{h}=\lambda_{1} \min \{1, h\}$,
that is, $a_{h}$ is also a bilinear, continuous, symmetric and coercive form $V \times V$ to $\mathbb{R}$. The existence and uniqueness of the solution to each of the above Problems 2 and 4, is well known see for example [3, 10-11].

The main goal of this paper is to prove in Section 2 the existence and uniqueness of a family of optimal control problems 5 and 6 where the control variable is given by a boundary condition of Neumann type whose state system is governed by a free boundary problem with Tresca conditions on a portion $\Gamma_{2}$ of the boundary, with a flux $f$ on $\Gamma_{3}$ as the control variable, using a regularization method to overcome the nondifferentiability of the functional $\Phi$. Then, in Section 3, we study the convergence when $h \rightarrow+\infty$ of the state systems and optimal controls associated to Problem 6 to the corresponding state system and optimal control associated to Problem 5. In order to obtain this last result we obtain an auxiliary strong convergence by using the Aubin compactness arguments (see Lemma 2). This paper completes our previous paper [2] and solves the open problem left in [1].

Remark here that our study still valid with the bilinear form $a$ in more general cases, provided that $a$ must be symmetric, coercive and continuous from $V \times V$ to $\mathbb{R}$.

## 2 Boundary optimal control problems

Let $M>0$ be a constant and we define the space

$$
\mathcal{F}_{-}=\{f \in \mathcal{F}: f \leqslant 0\} .
$$

We consider the following Neumannn boundary optimal control problems defined by [12-15].

Problem 5 Find the optimal control $f_{\mathrm{op}} \in \mathcal{F}_{-}$such that

$$
\begin{equation*}
J\left(f_{\mathrm{op}}\right)=\min _{f \in \mathcal{F}_{-}} J(f) \tag{3}
\end{equation*}
$$

where the cost functional $J: \mathcal{F}_{-} \rightarrow \mathbb{R}^{+}$is given by

$$
\begin{equation*}
J(f)=\frac{1}{2}\left\|u_{f}\right\|_{\mathcal{H}}^{2}+\frac{M}{2}\|f\|_{\mathcal{F}}^{2}(M>0) \tag{4}
\end{equation*}
$$

and $u_{f}$ is the unique solution to Problem 2 for a given $f \in \mathcal{F}_{-}$.

Problem 6 Find the optimal control $f_{\mathrm{op}_{h}} \in \mathcal{F}_{-}$such
that

$$
\begin{equation*}
J\left(f_{\mathrm{op}_{h}}\right)=\min _{f \in \mathcal{F}_{-}} J_{h}(f), \tag{5}
\end{equation*}
$$

where the cost functional $J_{h}: \mathcal{F}_{-} \rightarrow \mathbb{R}^{+}$is given by

$$
\begin{equation*}
J_{h}(f)=\frac{1}{2}\left\|u_{h f}\right\|_{\mathcal{H}}^{2}+\frac{M}{2}\|f\|_{\mathcal{F}}^{2}(M>0, h>0), \tag{6}
\end{equation*}
$$

and $u_{h f}$ is the unique solution to Problem 4 for a given $f \in \mathcal{F}_{-}$and $h>0$.

Theorem 1 Under the assumptions $g \geqslant 0$ in $\Omega \times(0, T)$, $b \geqslant 0$ on $\Gamma_{1} \times(0, T)$ and $u_{b} \geqslant 0$ in $\Omega$, we have the following properties:
a) the cost functional $J$ is strictly convex on $\mathcal{F}_{-}$; and
b) there exists a unique optimal control $f_{\text {op }} \in \mathcal{F}_{-}$solution to the Neumann boundary optimal control Problem 5.

Proof We give some sketch of the proof, following [2], we generalize for parabolic variational inequalities of the second kind, given in Problem 2, the estimates obtained for convex combination between $u_{4}(\mu)=u_{\mu f_{1}+(1-\mu) f_{2}}$, and $u_{3}(\mu)=\mu u_{f_{1}}+(1-\mu) u_{f_{2}}$, for any two element $f_{1}$ and $f_{2}$ in $\mathcal{F}$. The main difficulty, to prove this result comes from the fact that the functional $\Phi$ is not differentiable. To overcome this difficulty, we use the regularization method and consider for $\varepsilon>0$ the following approach of $\Phi$ defined by

$$
\begin{equation*}
\Phi_{\varepsilon}(v)=\int_{\Gamma_{2}} q \sqrt{\varepsilon^{2}+|v|^{2}} \mathrm{~d} s, \quad \forall v \in V \tag{7}
\end{equation*}
$$

which is Gateaux differentiable, with

$$
\left\langle\Phi_{\varepsilon}^{\prime}(w), v\right\rangle=\int_{\Gamma_{2}} \frac{q w v}{\sqrt{\varepsilon^{2}+|w|^{2}}} \mathrm{~d} s, \quad \forall(w, v) \in V^{2}
$$

We define $u^{\varepsilon}$ as the unique solution to the corresponding parabolic variational inequality for all $\varepsilon>0$. We obtain that for all $\mu \in[0,1]$ we have $u_{4}^{\varepsilon}(\mu) \leqslant u_{3}^{\varepsilon}(\mu)$ for all $\varepsilon>0$.

When $\varepsilon \rightarrow 0$ we have that for $i=1, \ldots, 4$,

$$
\begin{equation*}
u_{i}^{\varepsilon} \rightarrow u_{i} \text { strongly in } \mathcal{V} \cap L^{\infty}(0, T ; H) \tag{8}
\end{equation*}
$$

As $f \in \mathcal{F}_{-}, g \geqslant 0$ in $\Omega \times(0, T), b \geqslant 0$ in $\Gamma_{1} \times(0, T)$ and $u_{b} \geqslant 0$ in $\Omega$, we obtain by the weak maximum principle that for all $\mu \in[0,1]$ we have $0 \leqslant u_{4}(\mu)$, and so following [2], we obtain

$$
\begin{equation*}
0 \leqslant u_{4}(\mu) \leqslant u_{3}(\mu) \text { in } \Omega \times[0, T], \quad \forall \mu \in[0,1] . \tag{9}
\end{equation*}
$$

Then, for all $\mu \in[0,1]$, and for all $f_{1}, f_{2}$ in $\mathcal{F}_{-}$, and by using $f_{3}(\mu)=\mu f_{1}+(1-\mu) f_{2}$, we obtain

$$
\begin{align*}
& \mu J\left(f_{1}\right)+(1-\mu) J\left(f_{2}\right)-J\left(f_{3}(\mu)\right) \\
& =\frac{1}{2}\left(\left\|u_{3}(\mu)\right\|_{\mathcal{H}}^{2}-\left\|u_{4}(\mu)\right\|_{\mathcal{H}}^{2}\right)+\frac{1}{2} \mu(1-\mu)\left\|u_{f_{1}}-u_{f_{2}}\right\|_{\mathcal{H}}^{2} \\
& \quad+\frac{M}{2} \mu(1-\mu)\left\|f_{1}-f_{2}\right\|_{\mathcal{F}}^{2} . \tag{10}
\end{align*}
$$

Then, $J$ is strictly convex functional on $\mathcal{F}_{-}$, and therefore, there exists a unique optimal $f_{\text {op }} \in \mathcal{F}_{-}$solution to the Neumann boundary optimal control Problem 5. This completes the proof.

Theorem 2 Under the assumptions $g \geqslant 0$ in $\Omega \times(0, T)$, $b \geqslant 0$ in $\Gamma_{1} \times(0, T)$ and $u_{b} \geqslant 0$ in $\Omega$, we have the following properties:
a) the cost functional $J_{h}$ are strictly convex on $\mathcal{F}_{-}$, for all $h>0$; and
b) there exists a unique optimal control $f_{h_{\mathrm{op}}} \in \mathcal{F}_{-}$solution to the Neumann boundary optimal control problem 6, for all $h>0$.

Proof We follow a similar method to the one developed
in Theorem 1 for all $h>0$.

## 3 Convergence when $h \rightarrow+\infty$

In this section, we study the convergence of the Neumann optimal control problem 6 to the optimal control problem 5 when $h \rightarrow \infty$. For a given $f \in \mathcal{F}$, we have first the following result which generalizes $[2,6-7,16]$.

Lemma 1 Let $u_{h f}$ be the unique solution to Problem 4 and $u_{f}$ the unique solution to Problem 2, and then,
$u_{h f} \rightarrow u_{f} \in \mathcal{V}$ strongly as $h \rightarrow+\infty, \forall f \in \mathcal{F}$.
Proof Following [2], we take $v=u_{f}(t)$ in the variational inequality of Problem 4 where $u=u_{h f}$, and recalling that $u_{f}(t)=b$ on $\Gamma_{1} \times[0, T]$, taking $\phi_{h}(t)=u_{h f}(t)-u_{f}(t)$ we obtain for $h>1$, that $\left\|u_{h f}\right\|_{\mathcal{V}}$ is bounded for all $h>1$ and for all $f \in \mathcal{F}$. Then, there exists $\eta \in \mathcal{V}$ such that (when $h \rightarrow+\infty)$

$$
u_{h f} \rightharpoonup \eta \text { weakly in } \mathcal{V},
$$

and

$$
u_{h f} \rightarrow b \text { strongly on } L^{2}\left((0, T) \times \Gamma_{1}\right),
$$

and so $\eta(0)=u_{b}$.
Let $\varphi$ be in $L^{2}\left(0, T, H_{0}^{1}(\Omega)\right)$ and take the variational inequality of Problem 4 where $u=u_{h f}, v=u_{h f}(t) \pm \varphi(t)$, and then, as $\left\|u_{h f}\right\|_{\mathcal{V}}$ is bounded for all $h>1$, we deduce that $\left\|\dot{u}_{h f}\right\|_{L^{2}\left(0, T, H^{-1}(\Omega)\right)}$ is also bounded for all $h>1$. Then, we conclude that

$$
\begin{equation*}
u_{h f} \rightharpoonup \eta \text { in } \mathcal{V} \text { weak, and in } L^{\infty}(0, T, H) \text { weak star, } \tag{11}
\end{equation*}
$$

and $\dot{u}_{h f} \rightharpoonup \dot{\eta}$ in $L^{2}\left(0, T, H^{-1}(\Omega)\right)$ weak.
From the variational inequality of Problem 4, taking $v \in$ $K$ so $v=b$ on $\Gamma_{1}$, we obtain a.e. $t \in[0, T]$

$$
\begin{aligned}
& \left\langle\dot{u}_{h f}, v-u_{h f}\right\rangle+a\left(u_{h f}, v-u_{h f}\right)-h \int_{\Gamma_{1}}\left|u_{h f}-b\right|^{2} \mathrm{~d} s \\
& \geqslant \Phi\left(u_{h f}\right)-\Phi(v)+\left(g, v-u_{h f}\right)-\int_{\Gamma_{3}} f\left(v-u_{h f}\right) \mathrm{d} s
\end{aligned}
$$

for all $v \in K$, and then, as $h>0$ we have a.e. $t \in[0, T]$,

$$
\begin{align*}
& \left\langle\dot{u}_{h f}, v-u_{h f}\right\rangle+a\left(u_{h f}, v-u_{h f}\right) \\
& \geqslant \Phi\left(u_{h f}\right)-\Phi(v)+\left(g, v-u_{h f}\right)-\int_{\Gamma_{3}} f\left(v-u_{h f}\right) \mathrm{d} s \\
& \quad \forall v \in K . \tag{12}
\end{align*}
$$

Therefore, using (11) and passing to the limit, when $h \rightarrow$ $+\infty$, we obtain

$$
\begin{aligned}
& \langle\dot{\eta}, v-\eta\rangle+a(\eta, v-\eta)+\Phi(v)-\Phi(\eta) \\
& \geqslant(g, v-\eta)-\int_{\Gamma_{3}} f(v-\eta) \mathrm{d} s, \forall v \in K \text { a.e. } t \in[0, T]
\end{aligned}
$$

and $\eta(0)=u_{b}$. Using the uniqueness of the solution to Problem 2, we obtain $\eta=u_{f}$.

To prove the strong convergence, we take $v=u_{f}(t)$ in the variational inequality of Problem 4,

$$
\begin{aligned}
& \left\langle\dot{u}_{h f}, u_{f}-u_{h f}\right\rangle+a_{h}\left(u_{h f}, u_{f}-u_{h f}\right)+\Phi\left(u_{f}\right)-\Phi\left(u_{h f}\right) \\
& \geqslant\left(g, u_{f}-u_{h f}\right)+h \int_{\Gamma_{1}} b\left(u_{f}-u_{h f}\right) \mathrm{d} s \\
& \quad-\int_{\Gamma_{3}} f\left(u_{f}-u_{h f}\right) \mathrm{d} s
\end{aligned}
$$

a.e. $t \in[0, T]$, and thus as $u_{f}=u_{b}$ on $\Gamma_{1} \times[0, T]$, we put $\phi_{h}=u_{h f}-u_{f}$, and so a.e. $t \in[0, T]$ we have

$$
\begin{aligned}
& \left\langle\dot{\phi_{h}}, \phi_{h}\right\rangle+a\left(\phi_{h}, \phi_{h}\right)+h \int_{\Gamma_{1}}\left|\phi_{h}\right|^{2} \mathrm{~d} s+\Phi\left(u_{h f}\right)-\Phi\left(u_{f}\right) \\
& \leqslant\left\langle\dot{u}_{f}, \phi_{h}\right\rangle+a\left(u_{f}, \phi_{h}\right)+\left(g, \phi_{h}\right)-\int_{\Gamma_{3}} f \phi_{h} \mathrm{~d} s,
\end{aligned}
$$

and so

$$
\begin{aligned}
& \frac{1}{2}\left\|\phi_{h}\right\|_{L^{\infty}(0, T, H)}^{2}+\lambda_{h}\left\|\phi_{h}\right\|_{\mathcal{V}}^{2}+\Phi\left(u_{h f}\right)-\Phi\left(u_{f}\right) \\
& \leqslant-\int_{0}^{T}\left\langle\dot{u}_{f}(t), \phi_{h}(t)\right\rangle \mathrm{d} t-\int_{0}^{T} a\left(u_{f}(t), \phi_{h}(t)\right) \mathrm{d} t \\
& \quad+\int_{0}^{T}\left(g(t), \phi_{h}(t)\right) \mathrm{d} t-\int_{0}^{T} \int_{\Gamma_{3}} f \phi_{h} \mathrm{~d} s \mathrm{~d} t .
\end{aligned}
$$

Using the weak semicontinuity of $\Phi$ and the weak convergence (11), the right side of the above inequality tends to zero when $h \rightarrow+\infty$, and then, we deduce the strong convergence of $\phi_{h}=u_{h f}-u_{f}$ to 0 in $\mathcal{V} \cap L^{\infty}(0, T, H)$, for all $f \in \mathcal{F}_{-}$and the proof holds.

We prove now the following lemma by using the Aubin compactness arguments. Lemma 2 is very important and necessary which allow us to conclude this paper. Indeed this result is needed to pass to the limit exactly in the last term of the inequality (22) in the proof of the main Theorem 3.
Lemma 2 Let $u_{h f_{\mathrm{op}_{h}}}$ the state system defined by the unique solution to Problem 4, where the flux $f$ is replaced by $f_{\mathrm{op}_{h}}$. Then, for $h \rightarrow+\infty$, we have

$$
\begin{equation*}
u_{h f_{\mathrm{op}_{h}}} \rightarrow u_{f} \text { in } L^{2}((0, T) \times \partial \Omega) \tag{13}
\end{equation*}
$$

where $u_{f}$ is the the state system defined by the unique solution to Problem 2 with the flux $f$ on $\Gamma_{3}$.

Proof Let consider the variational inequality of Problem 4 with $u=u_{h f_{\mathrm{op}_{h}}}$ and $f=f_{\mathrm{op}_{h}}$, i.e.,

$$
\begin{align*}
& \left\langle u_{h f_{\mathrm{op}_{h}}}, v-u_{h f_{\mathrm{op}_{h}}}\right\rangle+a_{h}\left(u_{h f_{\mathrm{op}_{h}}}, v-u_{h f_{\mathrm{op}_{h}}}\right)+\Phi(v) \\
& -\Phi\left(u_{h f_{\mathrm{op}_{h}}}\right) \\
& \geqslant\left(g, v-u_{h f_{\mathrm{op}_{h}}}\right)-\int_{\Gamma_{3}} f_{\mathrm{op}_{h}}\left(v-u_{h f_{\mathrm{op}_{h}}}\right) \mathrm{d} s \\
& \quad+h \int_{\Gamma_{1}} b\left(v-u_{h f_{\mathrm{op}_{h}}}\right) \mathrm{d} s, \forall v \in V \tag{14}
\end{align*}
$$

and let $\varphi \in L^{2}\left(0, T ; H_{0}^{1}(\Omega)\right)$, and set $v=u_{h f_{\mathrm{op}_{h}}}(t) \pm \varphi(t)$ in (14), we obtain

$$
\left\langle\dot{u}_{h f_{\mathrm{op}_{h}}}, \varphi\right\rangle=(g, \varphi)-a\left(u_{h f_{\mathrm{op}_{h}}}, \varphi\right) .
$$

By integration in times for $t \in(0, T)$, we obtain

$$
\int_{0}^{T}\left\langle\dot{u}_{h f_{\mathrm{op}_{h}}}, \varphi\right\rangle \mathrm{d} t=\int_{0}^{T}(g, \varphi) \mathrm{d} t-\int_{0}^{T} a\left(u_{h f_{\mathrm{op} h}}, \varphi\right) \mathrm{d} t
$$

and thus, for $A=\left(c\|g\|_{\mathcal{H}}+\left\|u_{h f_{\text {op }_{h}}}\right\|_{\mathcal{V}}\right)$, we obtain

$$
\left|\int_{0}^{T}<\dot{u}_{h f_{\mathrm{op}_{h}}}, \varphi>\mathrm{d} t\right| \leqslant A\|\varphi\|_{L^{2}\left(0, T ; H_{0}^{1}(\Omega)\right)}
$$

where $c$ comes from the Poincaré inequality, and as in Lemma 1 we can obtain that $u_{h f_{\text {op }}}$ is bounded in $\mathcal{V}$, and so there exists a positive constant $C$ such that

$$
\begin{equation*}
\left\|\dot{u}_{h f_{\mathrm{op}_{h}}}\right\|_{L^{2}\left(0, T ; H^{-1}(\Omega)\right)} \leqslant C . \tag{15}
\end{equation*}
$$

Using now the Aubin compactness arguments, see for example [17] with the three Banach spaces $V, H^{\frac{2}{3}}(\Omega)$ and $H^{-1}(\Omega)$, and then,

$$
u_{h f_{\mathrm{op}_{h}}} \rightarrow u_{f} L^{2}\left(0, T ; H^{\frac{2}{3}}(\Omega)\right) .
$$

As the trace operator $\gamma_{0}$ is continuous from $H^{\frac{2}{3}}(\Omega)$ to $L^{2}(\partial \Omega)$, and then, the result follows. This completes the proof.

We give now, without need to use the notion of adjoint states [14,18], the convergence result which generalizes the result obtained in [19] for a parabolic variational equalities (see also [18, 20-23]). Other optimal control problems gouverned by variational inequalities are given in [24-26].
Theorem 3 Let $u_{h f_{\mathrm{op}_{h}}} \in \mathcal{V}, f_{\mathrm{op}_{h}} \in \mathcal{F}_{-}$and $u_{f_{\mathrm{op}}} \in \mathcal{V}$,
$f_{\text {op }} \in \mathcal{F}_{-}$be respectively the state systems and the optimal controls defined in Problems 4 and 2. Then,

$$
\begin{align*}
& \lim _{h \rightarrow+\infty}\left\|u_{h f_{\mathrm{op}_{h}}}-u_{f_{\mathrm{op}}}\right\| \mathcal{V} \\
& =\lim _{h \rightarrow+\infty}\left\|u_{h f_{\mathrm{op}_{h}}}-u_{f_{\mathrm{op}}}\right\|_{L^{\infty}(0, T, H)} \\
& =\lim _{h \rightarrow+\infty}\left\|u_{h f_{\mathrm{op}_{h}}}-u_{f_{\mathrm{op}}}\right\|_{L^{2}\left((0, T) \times \Gamma_{1}\right)}=0,  \tag{16}\\
& \lim _{h \rightarrow+\infty}\left\|f_{\mathrm{op}_{h}}-f_{\mathrm{op}}\right\|_{\mathcal{F}}=0 . \tag{17}
\end{align*}
$$

Proof We have first

$$
\begin{aligned}
J_{h}\left(f_{\mathrm{op}_{h}}\right) & =\frac{1}{2}\left\|u_{h f_{\mathrm{op}_{h}}}\right\|_{\mathcal{H}}^{2}+\frac{M}{2}\left\|f_{\mathrm{op}_{h}}\right\|_{\mathcal{F}}^{2} \\
& \leqslant \frac{1}{2}\left\|u_{h f}\right\|_{\mathcal{H}}^{2}+\frac{M}{2}\|f\|_{\mathcal{F}}^{2}
\end{aligned}
$$

for all $f \in \mathcal{F}_{-}$, and then, for $f=0 \in \mathcal{F}_{-}$we obtain

$$
\begin{equation*}
J_{h}\left(f_{\mathrm{op}_{h}}\right)=\frac{1}{2}\left\|u_{h f_{\mathrm{op}_{h}}}\right\|_{\mathcal{H}}^{2}+\frac{M}{2}\left\|f_{\mathrm{op}_{h}}\right\|_{\mathcal{F}}^{2} \leqslant \frac{1}{2}\left\|u_{h 0}\right\|_{\mathcal{H}}^{2}, \tag{18}
\end{equation*}
$$

where $u_{h 0} \in \mathcal{V}$ is the solution to the following parabolic variational inequality:

$$
\begin{aligned}
& \left\langle\dot{u}_{h 0}, v-u_{h 0}\right\rangle+a_{h}\left(u_{h 0}, v-u_{h 0}\right)+\Phi(v)-\Phi\left(u_{h 0}\right) \\
& \geqslant \int_{\Omega} g\left(v-u_{h 0}\right) \mathrm{d} x+h \int_{\Gamma_{1}} b\left(v-u_{h 0}\right) \mathrm{d} s, \text { a.e. } t \in[0, T]
\end{aligned}
$$ for all $v \in V$ and $u_{h 0}(0)=u_{b}$.

Taking $v=u_{b} \in K_{b}$ we obtain that $\left\|u_{h 0}-u_{b}\right\|_{\mathcal{V}}$ is bounded independently of $h$, and then, $\left\|u_{h 0}\right\|_{\mathcal{H}}$ is bounded independently of $h$. Therefore, we deduce with (18) that $\left\|u_{h f_{\text {op }_{h}}}\right\|_{\mathcal{H}}$ and $\left\|f_{\mathrm{op}_{h}}\right\|_{\mathcal{F}}$ are also bounded independently of $h$. Therefore, there exist $\tilde{f} \in \mathcal{F}_{-}$and $\eta$ in $\mathcal{H}$ such that
$f_{\mathrm{op}_{h}} \rightharpoonup \tilde{f}$ in $\mathcal{F}_{-}$and $u_{h f_{\mathrm{op}_{h}}} \rightharpoonup \eta$ in $\mathcal{H}$ (weakly).
Taking now $v=u_{f_{\text {op }}}(t) \in K_{b}$ in Problem (4), for $t \in[0, T]$, with $u=u_{h f_{\mathrm{op}_{h}}}$ and $f=f_{\mathrm{op}_{h}}$, we obtain

$$
\begin{aligned}
& \left\langle\dot{u}_{h f_{\mathrm{op}_{h}}}, u_{f_{\mathrm{op}_{\mathrm{p}}}}-u_{h f_{\mathrm{op}_{h}}}\right\rangle+a_{1}\left(u_{h f_{\mathrm{op}_{h}}}, u_{\left.f_{\mathrm{opp}}-u_{h f_{\mathrm{op}_{h}}}\right)}\right. \\
& +(h-1) \int_{\Gamma_{1}} u_{h f_{\mathrm{op}_{h}}}\left(u_{f_{\mathrm{op}}}-u_{h f_{\mathrm{op}_{h}}}\right) \mathrm{d} s+\Phi\left(u_{f_{\mathrm{op}}}\right) \\
& -\Phi\left(u_{h f_{\mathrm{op}_{h}}}\right)
\end{aligned}
$$

$$
\begin{aligned}
& -\int_{\Gamma_{3}} f_{\mathrm{op}_{h}}\left(u_{\left.f_{\mathrm{opp}}-u_{h f_{\mathrm{op}_{h}}}\right) \mathrm{d} s, \text { a.e. } t \in[0, T] . ~ . ~ . ~ . ~}^{\text {. }}\right. \text {. }
\end{aligned}
$$

As $u_{f_{\mathrm{op}}}=b$ on $\Gamma_{1} \times[0, T]$, taking $\phi_{h}=u_{f_{\mathrm{op}}}-u_{h f_{\mathrm{op}_{h}}}$, we obtain

$$
\begin{aligned}
& \frac{1}{2}\left\|\phi_{h}\right\|_{L^{\infty}(0, T ; H)}^{2}+\lambda_{1}\left\|\phi_{h}\right\|_{\mathcal{V}}^{2} \\
& +(h-1) \int_{0}^{T} \int_{\Gamma_{1}}\left|\phi_{h}(t)\right|^{2} \mathrm{~d} s \mathrm{~d} t \\
& \leqslant \int_{0}^{T} \int_{\Gamma_{3}} f_{\mathrm{op}_{h}} \phi_{h} \mathrm{~d} s \mathrm{~d} t-\int_{0}^{T}\left(g(t), \phi_{h}(t)\right) \mathrm{d} t \\
& \quad+\int_{0}^{T} \int_{\Gamma_{2}} q\left|\phi_{h}(t)\right| \mathrm{d} s \mathrm{~d} t+\int_{0}^{T}\left\langle\dot{u}_{f_{\mathrm{op}}}(t) \phi_{h}(t)\right\rangle \mathrm{d} t \\
& \quad+\int_{0}^{T} a\left(u_{f_{\mathrm{op}}}(t), \phi_{h}(t)\right) \mathrm{d} t .
\end{aligned}
$$

As $f_{\mathrm{op}_{h}}$ is bounded in $\mathcal{F}_{-}$, from (15) $\dot{u}_{f_{\mathrm{op}}}$ is bounded in $L^{2}\left(0, T ; H^{-1}(\Omega)\right)$, and $u_{h f_{\mathrm{op}_{h}}}$ is also bounded in $\mathcal{V}$, all independently on $h$, and so there exists a positive constant $C$ which does not depend on $h$ such that

$$
\begin{aligned}
& \left\|\phi_{h}\right\| \mathcal{V}=\left\|u_{h f_{\mathrm{op}_{h}}}-u_{f_{\mathrm{op}}}\right\| \mathcal{V} \leqslant C,\left\|\phi_{h}\right\|_{L^{\infty}(0, T, H)} \leqslant C \\
& (h-1) \int_{0}^{T} \int_{\Gamma_{1}}\left|u_{h f_{\mathrm{op}_{h}}}-b\right|^{2} \mathrm{~d} s \mathrm{~d} t \leqslant C
\end{aligned}
$$

and then, $\eta \in \mathcal{V}$ and

$$
\begin{align*}
& u_{h f_{\mathrm{op}_{h}}} \rightharpoonup \eta \text { in } \mathcal{V} \text { and in } L^{\infty}(0, T, H) \text { weak star, }  \tag{20}\\
& u_{h f_{\mathrm{op}_{h}}} \rightarrow b \text { in } L^{2}\left((0, T) \times \Gamma_{1}\right) \text { strong } \tag{21}
\end{align*}
$$

and so $\eta(t) \in K_{b}$ for all $t \in[0, T]$.
Now, taking $v \in K$ in Problem 4 where $u=u_{h f_{\mathrm{op}_{h}}}$ and $f=f_{\text {op }_{h}}$ so

$$
\begin{aligned}
& \left\langle\dot{u}_{h f_{\mathrm{op}_{h}}}, v-u_{h f_{\mathrm{op}_{h}}}\right\rangle+a_{h}\left(u_{h f_{\mathrm{op}_{h}}}, v-u_{h f_{\mathrm{op}_{h}}}\right)+\Phi(v) \\
& -\Phi\left(u_{h f_{\mathrm{op}_{h}}}\right) \\
& \geqslant\left(f_{\mathrm{op}_{h}}, v-u_{h f_{\mathrm{op}_{h}}}\right)+h \int_{\Gamma_{1}} b\left(v-u_{h f_{\mathrm{op}_{h}}}\right) \mathrm{d} s \\
& \quad-\int_{\Gamma_{3}} f_{\mathrm{op}_{h}}\left(v-u_{h f_{\mathrm{op}_{h}}}\right) \mathrm{d} s, \text { a.e. } t \in[0, T]
\end{aligned}
$$

as $v \in K_{b}$ so $v=b$ on $\Gamma_{1}$, and thus, we have

$$
\begin{aligned}
& \left\langle\dot{u}_{h f_{\mathrm{op}_{h}}}, u_{h f_{\mathrm{op}_{h}}}-v\right\rangle+a\left(u_{h f_{\mathrm{op}_{h}}}, u_{h f_{\mathrm{op}_{h}}}-v\right) \\
& +h \int_{\Gamma_{1}}\left|u_{h f_{\mathrm{op}_{h}}}-b\right|^{2} \mathrm{~d} s+\Phi\left(u_{h f_{\mathrm{op}_{h}}}\right)-\Phi(v) \\
& -\left(g, v-u_{h f_{\mathrm{op}_{h}}}\right) \\
& \leqslant \int_{\Gamma_{3}} f_{\mathrm{op}_{h}}\left(v-u_{h f_{\mathrm{op}_{h}}}\right) \mathrm{d} s \text { a.e. } t \in[0, T] .
\end{aligned}
$$

Thus,

$$
\begin{align*}
& \left\langle\dot{u}_{h f_{\mathrm{op}_{h}}}, u_{h f_{\mathrm{op}_{h}}}-v\right\rangle+a\left(u_{h f_{\mathrm{op}_{h}}}, u_{h f_{\mathrm{op}_{h}}}-v\right) \\
& +\Phi\left(u_{h f_{\mathrm{op}_{h}}}\right)-\Phi(v) \\
& \leqslant-\left(g, v-u_{h f_{\mathrm{op}_{h}}}\right) \\
& \quad-\int_{\Gamma_{3}} f_{\mathrm{op}_{h}}\left(v-u_{h f_{\mathrm{op}_{h}}}\right) \mathrm{d} s \text { a.e. } t \in[0, T] . \tag{22}
\end{align*}
$$

Using Lemma 2, (19) and (20), we deduce that [3,27]
$\langle\dot{\eta}, v-\eta\rangle+a(\eta, v-\eta)+\Phi(v)-\Phi(\eta)$
$\left.\geqslant(f, v-\eta)-\int_{\Gamma_{3}} \tilde{f}(v-\eta)\right) \mathrm{d} s, \quad \forall v \in K$, a.e. $t \in[0, T]$, so also by the uniqueness of the solution to Problem 2, we obtain

$$
\begin{equation*}
u_{\tilde{f}}=\eta \tag{23}
\end{equation*}
$$

We prove that $\tilde{f}=f_{\text {op }}$. Indeed, we have

$$
\begin{aligned}
J(\tilde{f}) & =\frac{1}{2}\|\eta\|_{\mathcal{H}}^{2}+\frac{M}{2}\|\tilde{f}\|_{\mathcal{F}}^{2} \\
& \leqslant \liminf _{h \rightarrow+\infty}\left\{\frac{1}{2}\left\|u_{h f_{\mathrm{op}_{h}}}\right\|_{\mathcal{H}}^{2}+\frac{M}{2}\left\|f_{\mathrm{op}_{h}}\right\|_{\mathcal{F}}^{2}\right\} \\
& =\liminf _{h \rightarrow+\infty} J_{h}\left(f_{\mathrm{op}_{h}}\right) \\
& \leqslant \liminf _{h \rightarrow+\infty} J_{h}(f)=\liminf _{h \rightarrow+\infty}\left\{\frac{1}{2}\left\|u_{h f}\right\|_{\mathcal{H}}^{2}+\frac{M}{2}\|f\|_{\mathcal{F}}^{2}\right\},
\end{aligned}
$$

and so using now the strong convergence $u_{h f} \rightarrow u_{f}$ as $h \rightarrow+\infty, \forall f \in \mathcal{F}_{-}$(see Lemma 1), we obtain

$$
\begin{align*}
J(\tilde{f}) \leqslant \liminf _{h \rightarrow+\infty} J_{h}\left(f_{\mathrm{op}_{h}}\right) & \leqslant \frac{1}{2}\left\|u_{f}\right\|_{\mathcal{H}}^{2}+\frac{M}{2}\|f\|_{\mathcal{F}}^{2} \\
& =J(f), \quad \forall f \in \mathcal{F}_{-}, \tag{24}
\end{align*}
$$

and then, by the uniqueness of the optimal control problem 2, we obtain

$$
\begin{equation*}
\tilde{f}=f_{\mathrm{op}} \tag{25}
\end{equation*}
$$

Now, we prove the strong convergence of $u_{h f_{\mathrm{op}_{h}}}$ to $\eta=$ $u_{f}$ in $\mathcal{V} \cap L^{\infty}(0, T ; H) \cap L^{2}\left(0, T ; L^{2}\left(\Gamma_{1}\right)\right)$, indeed taking $v=\eta$ in Problem 4 where $u=u_{h f_{\mathrm{op}_{h}}}$ and $f=f_{\mathrm{op}_{h}}$, as $\eta(t) \in K$ for $t \in[0, T]$, and so $\eta=b$ on $\Gamma_{1}$, we obtain

$$
\frac{1}{2}\left\|u_{h f_{\mathrm{op}_{h}}}-\eta\right\|_{L^{\infty}(0, T ; H)}^{2}+\lambda_{1}\left\|u_{h f_{\mathrm{op}_{h}}}-\eta\right\|_{\mathcal{V}}^{2}
$$

$$
\begin{aligned}
& +\int_{0}^{T}\left\{\Phi\left(u_{h f_{\mathrm{op}_{h}}}\right)-\Phi(\eta)\right\} \mathrm{d} t+\tilde{h}\left\|u_{h f_{\mathrm{op}_{h}}}-\eta\right\|_{L^{2}\left((0, T) \times \Gamma_{1}\right)}^{2} \\
& \leqslant \int_{0}^{T}\left(g, u_{h f_{\mathrm{op}_{h}}}(t)-\eta(t)\right) \mathrm{d} t-\int_{0}^{T}\left\langle\dot{\eta}, u_{h f_{\mathrm{op}_{h}}}-\eta\right\rangle \mathrm{d} t \\
& \quad+\int_{0}^{T} a\left(\eta(t), \eta(t)-u_{h f_{\mathrm{op}_{h}}}(t)\right) \\
& \left.\quad-\int_{\Gamma_{3}} f_{\mathrm{op}_{h}}\left(u_{h f_{\mathrm{op}_{h}}}-\eta\right)\right) \mathrm{d} s \mathrm{~d} t,
\end{aligned}
$$

where $\tilde{h}=h-1$.
Using (20) and the weak semicontinuity of $\Phi$, we deduce that

$$
\begin{aligned}
& \lim _{h \rightarrow+\infty}\left\|u_{h f_{\mathrm{op}_{h}}}-\eta\right\|_{L^{\infty}(0, T ; H)} \\
& =\lim _{h \rightarrow+\infty}\left\|u_{h f_{\mathrm{op}_{h}}}-\eta\right\| \mathcal{V} \\
& =\left\|u_{h f_{\mathrm{op}_{h}}}-\eta\right\|_{L^{2}\left((0, T) \times \Gamma_{1}\right)}=0,
\end{aligned}
$$

and with (23) and (25) we deduce (16). Then, from (24) and (25), we can write

$$
\begin{align*}
J\left(f_{\mathrm{op}}\right) & =\frac{1}{2} \| u_{f_{\mathrm{opp}}\left\|_{\mathcal{H}}^{2}+\frac{M}{2}\right\| f_{\mathrm{op}} \|_{\mathcal{F}}^{2} \leqslant \liminf _{h \rightarrow+\infty} J_{h}\left(f_{\mathrm{op}_{h}}\right)} \\
& =\liminf _{h \rightarrow+\infty}\left\{\frac{1}{2}\left\|u_{h f_{\mathrm{op}_{h}}}\right\|_{\mathcal{H}}^{2}+\frac{M}{2}\left\|f_{\mathrm{op}_{h}}\right\|_{\mathcal{F}}^{2}\right\} \\
& \leqslant \lim _{h \rightarrow+\infty} J_{h}\left(f_{\mathrm{op}}\right)=J\left(f_{\mathrm{op}}\right), \tag{26}
\end{align*}
$$

and using the strong convergence (16), we obtain

$$
\begin{equation*}
\lim _{h \rightarrow+\infty}\left\|f_{\mathrm{op}_{h}}\right\|_{\mathcal{F}}=\left\|f_{\mathrm{op}}\right\|_{\mathcal{F}} \tag{27}
\end{equation*}
$$

Finally as

$$
\begin{equation*}
\left\|f_{\mathrm{op}_{h}}-f_{\mathrm{op}}\right\|_{\mathcal{F}}^{2}=\left\|f_{\mathrm{op}_{h}}\right\|_{\mathcal{F}}^{2}+\left\|f_{\mathrm{op}}\right\|_{\mathcal{F}}^{2}-2\left(f_{\mathrm{op}_{h}}, f_{\mathrm{op}}\right), \tag{28}
\end{equation*}
$$

and by the first part of (19) we have

$$
\lim _{h \rightarrow+\infty}\left(f_{\mathrm{op}_{h}}, f_{\mathrm{op}}\right)=\left\|f_{\mathrm{op}}\right\|_{\mathcal{F}}^{2},
$$

and so from (27) and (28) we obtain (17). This completes the proof.

Corollary 1 Let $u_{h f_{\mathrm{op}_{h}}}$ in $\mathcal{V}, f_{\mathrm{op}_{h}}$ in $\mathcal{F}_{-}, u_{f_{\mathrm{op}}}$ in $\mathcal{V}$ and $f_{\text {op }}$ in $\mathcal{F}_{-}$be respectively the state systems and the optimal controls defined in Problems 4 and 2. Then,

$$
\lim _{h \rightarrow+\infty}\left|J_{h}\left(f_{\mathrm{op}_{h}}\right)-J\left(f_{\mathrm{op}}\right)\right|=0
$$

Proof It follows from the definitions (3) and (4), and the convergences (16) and (17).

## 4 Conclusions

The main difference between this paper and our previous work [2] where the control variable was the function $g$, is that we consider here as a control variable the function $f$ given by the Neumann boundary condition on $\Gamma_{3}$. This change induce in the variational problems 2 and 4, and also in the proofs of Lemma 1 and Theorem 3, a new integral term on $\Gamma_{3}$. The main difficulty here is in Section 3, and the question is exactly how to pass to the limit for $h \rightarrow+\infty$ in the last integral term on $\Gamma_{3}$ in (22). To overcome this main difficulty we have introduced the new lemma 2 , which is the key of our problem. The idea of Lemma 1 and Theorem 3 and their proofs are indeed similar to those of our work [2] with the differences and difficulties mentioned just above.

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Mahdi BOUKROUCHE received his Ph.D. degree in Mathematics from Lille University France, in 1982. Since 1982 he has been associate professor in Annaba University Algeria. He received his Doctorat d'Etat es Sciences Mathematics in Applied Mathematics from Lyon University France, in 1993. Since 1993 he is a professor at the University of Saint-Etienne. His main research interests are on issues of free boundaries, hydrodynamic lubrication, cavitation, the problems of type: Hele-Shaw and Stefan, asymptotic
analysis, homogenization, dynamic PDE systems and turbulence. E-mail Mahdi.Boukrouche@univ-st-etienne.fr.


Domingo A. TARZIA received his B.S. degree in Mathematics and in Physics from Rosario University (Argentina) in 1972 and 1977 respectively. He received his PhD and Habilitation in Mathematics from Pierre-et-Marie-Curie University (Univ. Paris VI, France) in 1979 and 1991 respecively. He received also a Magister in Finance from Rosario University (Argentina) in 2010. He has been professor as the Rosario University (Argentina) during the period 1979 to 1990. Actually he is research director at CONICET since 1983 and works at Austral University (Rosario, Argentina) where he is the chair of the Mathematics Department since 1991. His research themes include the study of the free boundary problems for the heat equation, specially by variational inequalities and integral equations, and particularly by obteining explicit solutions; optimal control problems for PDE. He is president of ASAMACI (Asociacin Argentina de Matemtica Aplicada, Computacional e Industrial) and vice-president of AR-SIAM (Argentinean Section of SIAM). E-mail: DTarzia@austral.edu.ar.

